

Introduction to the Inner Life of the Proton

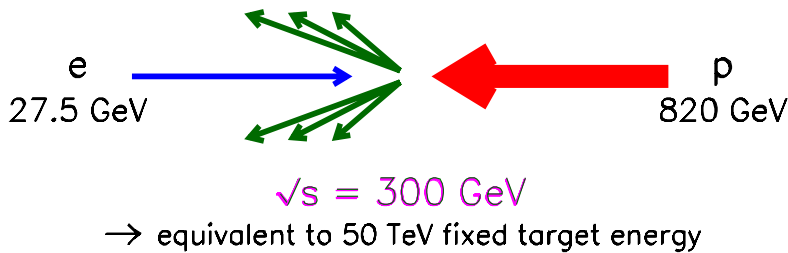
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- The **H**adron **E**lectron **R**ing-**A**nlage
- **D**eep-**I**nelastic **S**cattering
- The **Q**uark **P**arton **M**odel
- **Q**uantum **C**hromo **D**ynamics
- The Strong Coupling Constant α_s
- **P**arton **D**ensity **F**unctions

- Characteristics:

- ep -collisions with $E_e = 27.5$ GeV and $E_p = 920$ GeV
(in 1996/97: $E_p = 820$ GeV)

- boosted final state



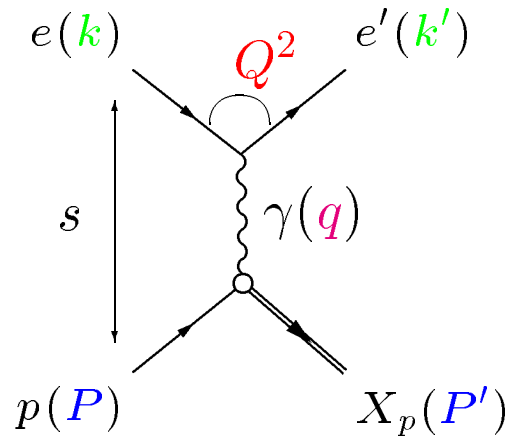
- hadron in initial state

- Possibilities:

- to understand $ep \rightarrow eX$ transition:
precise determination of the partonic structure of the proton

- test of theoretical QCD predictions

- e.g. measurement of strong coupling constant α_s ,
diffractive processes etc.



negative four-momentum transfer squared:

$$Q^2 = -q^2 = -(k - k')^2$$

parton momentum fraction (Bjorken scaling variable):

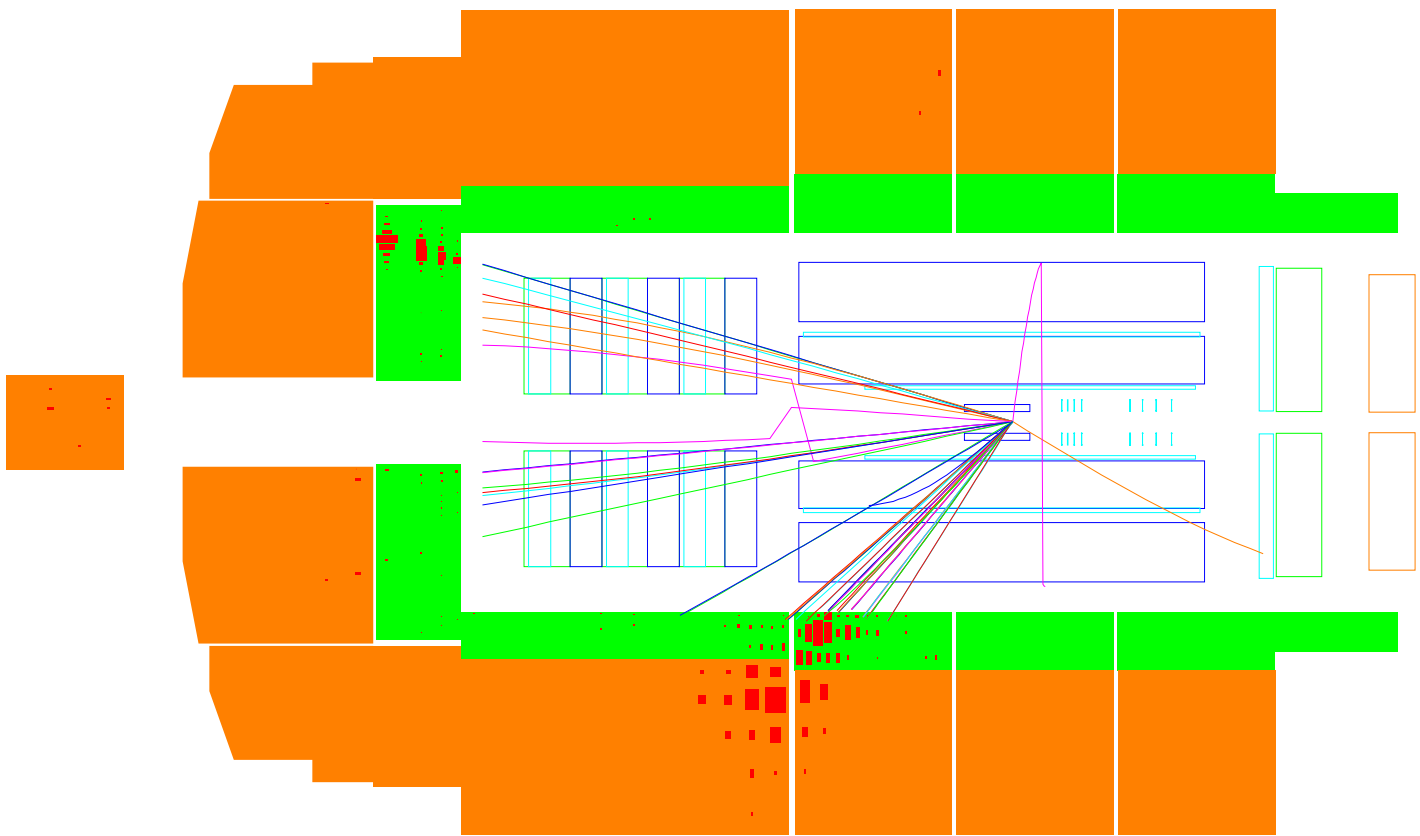
$$x = \frac{Q^2}{2P \cdot q}$$

inelasticity (fractional energy transfer in p rest frame):

$$y = \frac{q \cdot P}{k \cdot P} = 1 - \frac{E_{e'}}{E_e} \sin^2 \left(\frac{\theta_e}{2} \right)$$

$$\Rightarrow Q^2 = x \cdot y \cdot s$$

The $\Pi 1$ Detector



$$Q^2 = 34000 \text{ GeV}^2$$

$$x = 0.46$$

$$y = 0.8$$

- $Q^2 \gg m_{Proton}^2$: deep-inelastic
- differential cross-section $ep \rightarrow eX$
(neglecting Z^0 -exchange):

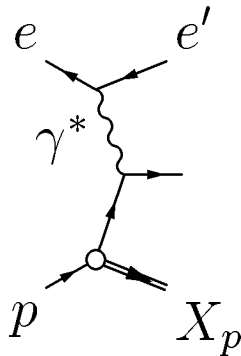
$$\frac{d^2\sigma}{dx dQ^2} = \frac{4\pi\alpha^2}{xQ^4} \cdot Y_+ \cdot [F_2(x, Q^2) - \frac{y^2}{Y_+} \cdot F_L(x, Q^2)]$$

- $\frac{y^2}{Y_+} = \frac{y^2}{1+(1-y)^2} < 1$ for all $y < 1$
- F_2 and F_L are **structure functions**, F_L corresponds to coupling of longitudinally polarized photons to the proton
- **What are structure functions?**
 - parameterization of the partonic structure of the proton
 - ⇒ **Measurement of structure functions provides information about the inner life of the proton**

For non-interacting partons (Quark-Parton Model):
 F_2 in terms of quark distributions $q_i(x)$:

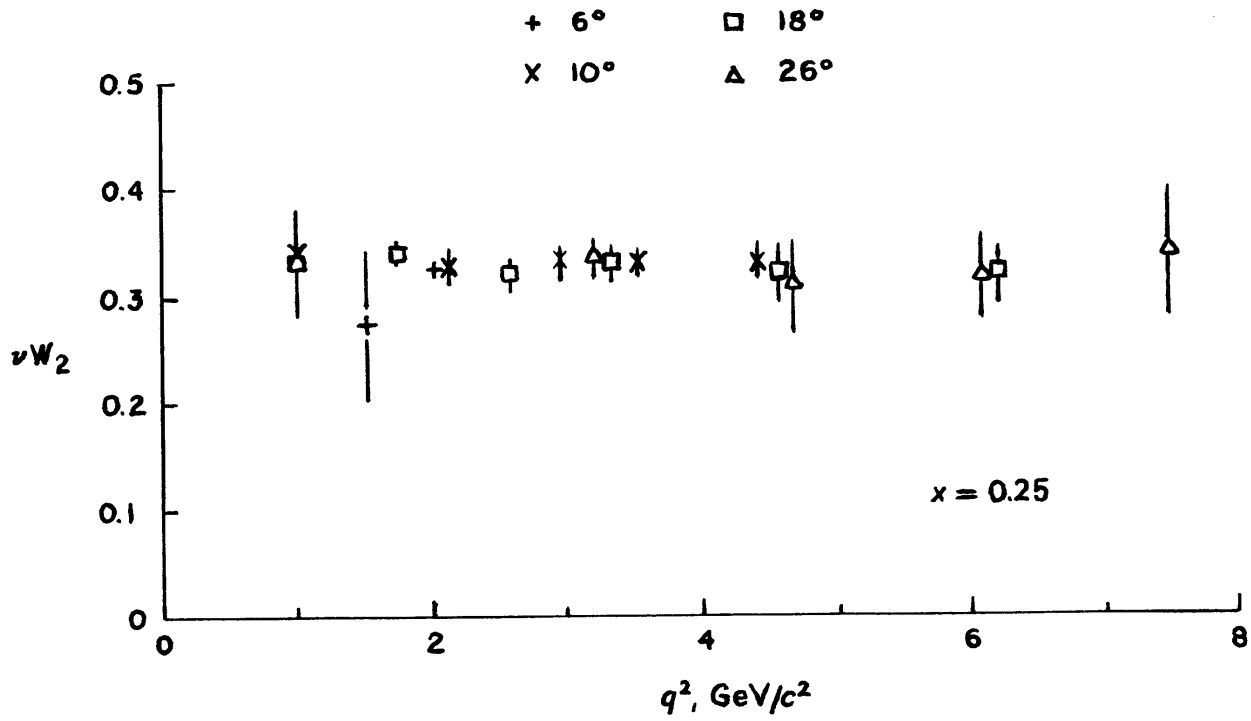
$$F_2(x) = x \sum_i e_i^2 \cdot q_i(x)$$

- virtual photon γ^* scatters off partons in the proton



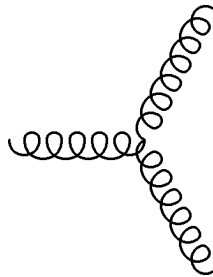
- quarks are massless spin- $\frac{1}{2}$ particles without transverse momentum k_{\perp}
 \Rightarrow no coupling of longitudinally polarized γ^* (due to spin constraints)
 $\Rightarrow F_L = 0$

- In the QPM $F_2(x)$ does not depend on Q^2
 \Rightarrow **BJORKEN SCALING**



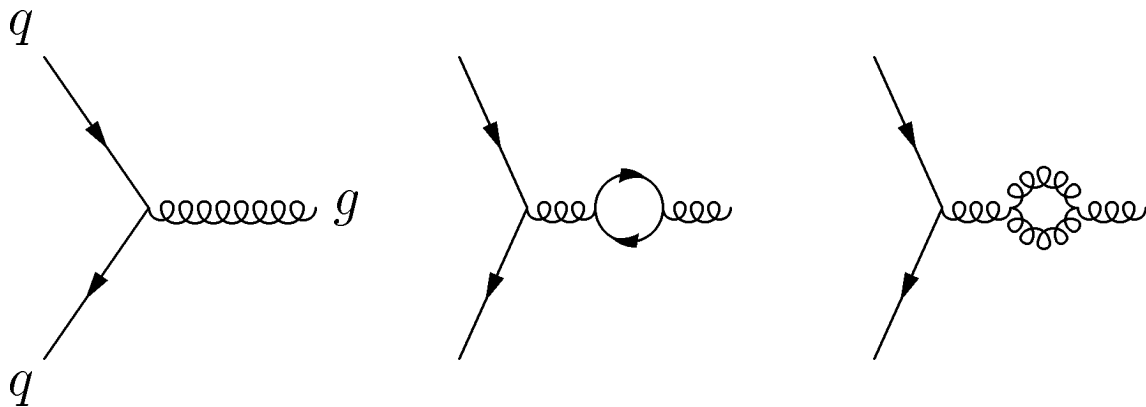
E.D. Bloom et al. Phys.Rev.Lett. (1969) 930

- QCD describes the **strong interaction** between quarks and gluons
- mathematical structure: non-abelian gauge invariant field theory
- gauge bosons in QCD are gluons
- quarks **and gluons** carry colour (\leftrightarrow QED: photons have no charge)
 \Rightarrow **selfinteracting gluons**:



- quarks can have transverse momentum due to gluon emission and absorption
 \Rightarrow coupling to longitudinally polarized γ^* possible
 \Rightarrow **$F_L \neq 0$**
- only free parameter of QCD:
strong coupling constant α_s

- QCD predictions for cross-section using **Feynman rules**:



- virtual loop corrections \Rightarrow infinite integrals
 \Rightarrow **renormalization and regularization**
- **regularization** of divergent integrals by defining upper integration boundary
- **renormalization** by redefinition of coupling constant for all other terms (introduction of renormalization scale μ_R and scheme i.e. \overline{MS})
- requirement: observables independent of μ_R :
 \Rightarrow **Renormalization Group Equation**

\Rightarrow consequence: $\alpha_s := \alpha_s(\mu_R)$

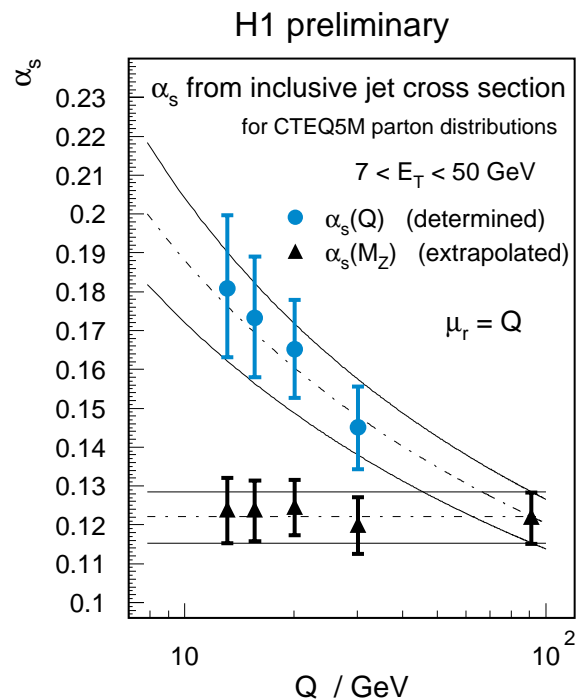
- Solution for α_s in leading order:

$$\alpha_s(Q^2) = \frac{\alpha_s(\mu_R^2)}{1 + \frac{\alpha_s(\mu_R^2)}{12\pi}(33 - 2n_f) \ln\left(\frac{Q^2}{\mu_R^2}\right)}$$

n_f = number of quark flavours

⇒ **Running coupling**

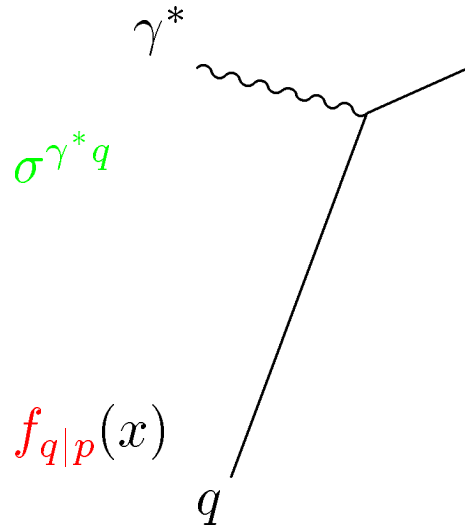
- $Q^2 < \mu_R^2 \Rightarrow \alpha_s(Q^2) > \alpha_s(\mu_R^2)$
 $Q^2 > \mu_R^2 \Rightarrow \alpha_s(Q^2) < \alpha_s(\mu_R^2)$
- standard measurement: $\alpha_s(M_Z)$ as reference scale



- **short distance scale** ($Q^2 \gtrsim 10 \text{ GeV}$) :
 - γ^* interacts with partons
 - \Rightarrow **asymptotic freedom**
 - \Rightarrow calculable in perturbative QCD

- **long distance scale** ($Q^2 \lesssim 0.1 \text{ GeV}$):
 - partons appear in bound states
 - \Rightarrow **confinement**
 - \Rightarrow perturbative QCD not applicable

Deep-inelastic γ^*p scattering =
incoherent sum of elastic γ^*q scatterings



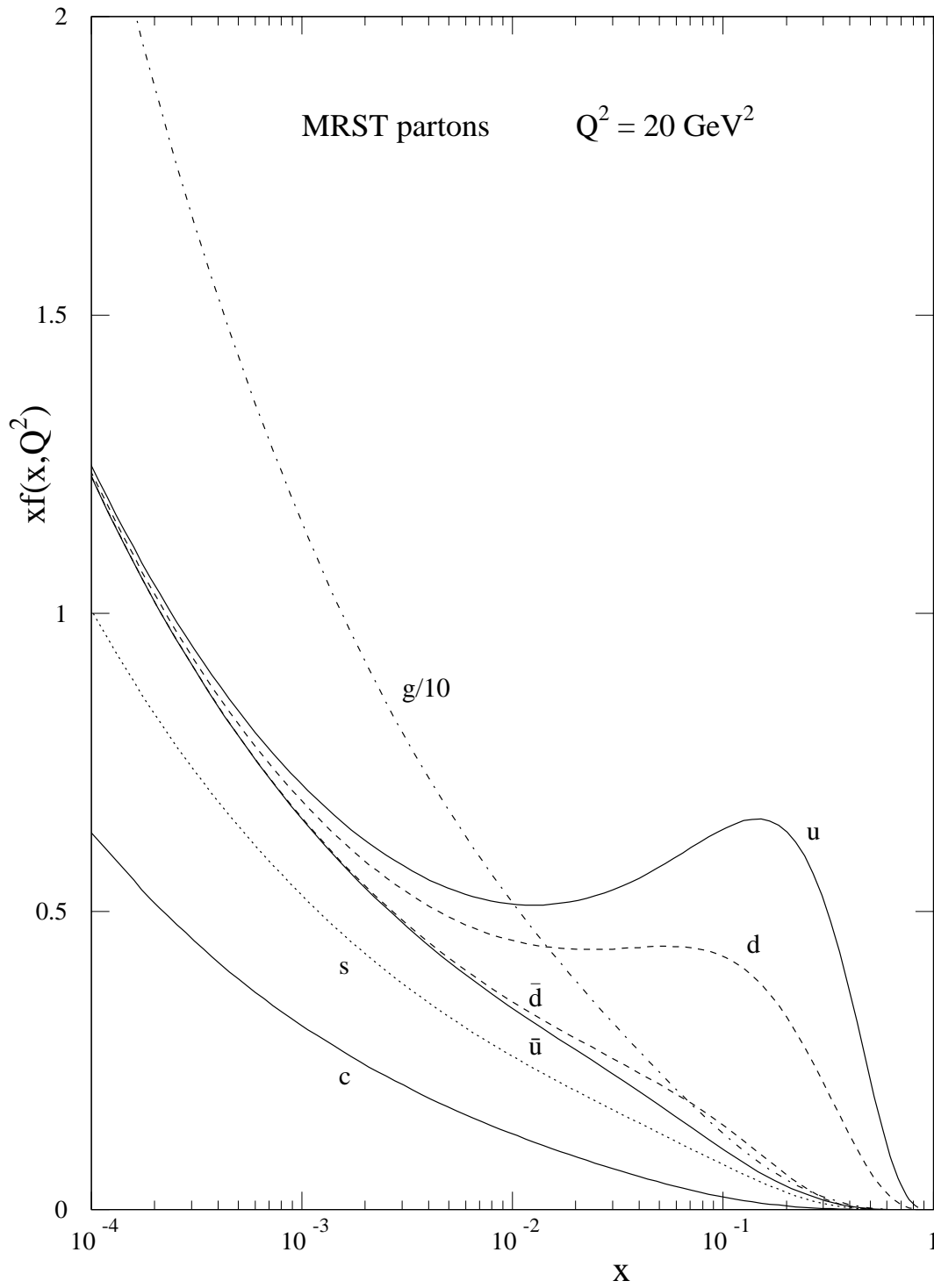
σ^{γ^*q} = electron-quark cross-section

$f_{q|p}(x) \equiv q(x) =$ Parton density function (**PDF**)

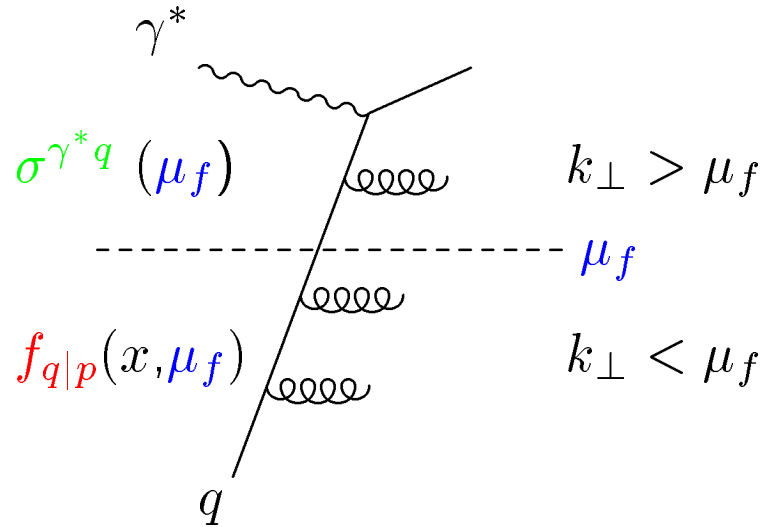
Probability density to find quark q in the proton carrying a fraction x of the proton momentum

$$\sigma^{\gamma^*p} = \sum_{q, \bar{q}} \sigma^{\gamma^*q} \cdot f_{q|p}(x)$$

Factorization of **perturbative** and **non-perturbative** process



- PDFs specify a **universal** property of the proton (\rightarrow process independent)
- Higher order corrections:



- integration over all transverse momenta k_\perp
 \Rightarrow divergences for $k_\perp \rightarrow 0$ (**infrared divergences**)
- curing divergences: redefinition of PDFs (introduction of factorization scale μ_f and scheme i.e. \overline{MS})
- for $k_\perp < \mu_f$: parton emissions are included in PDFs $f_{q|p}$

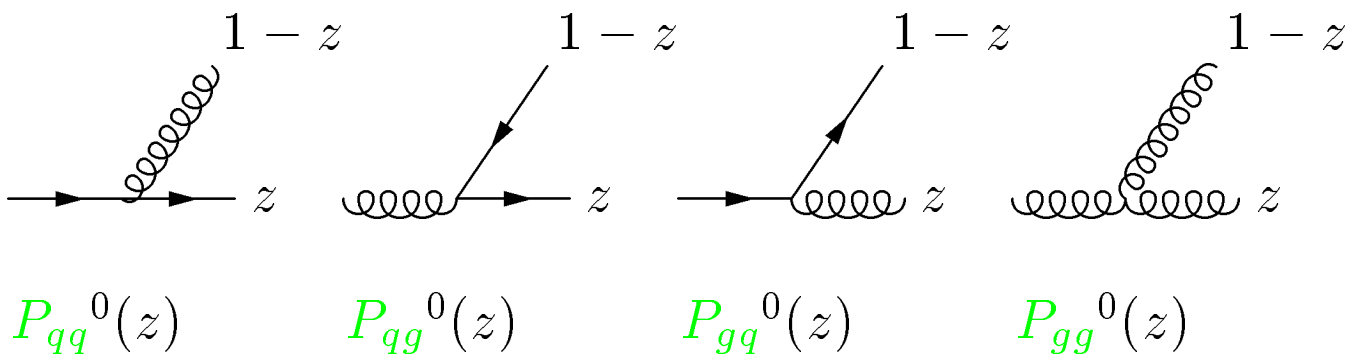
$$\sigma^{\gamma^* p} = \sum_{q, \bar{q}} \sigma^{\gamma^* q}(\mu_f^2) \cdot f_{q|p}(\mu_f^2)$$

Dokshitzer-Gribov-Lipatov-Altarelli-Parisi-equation

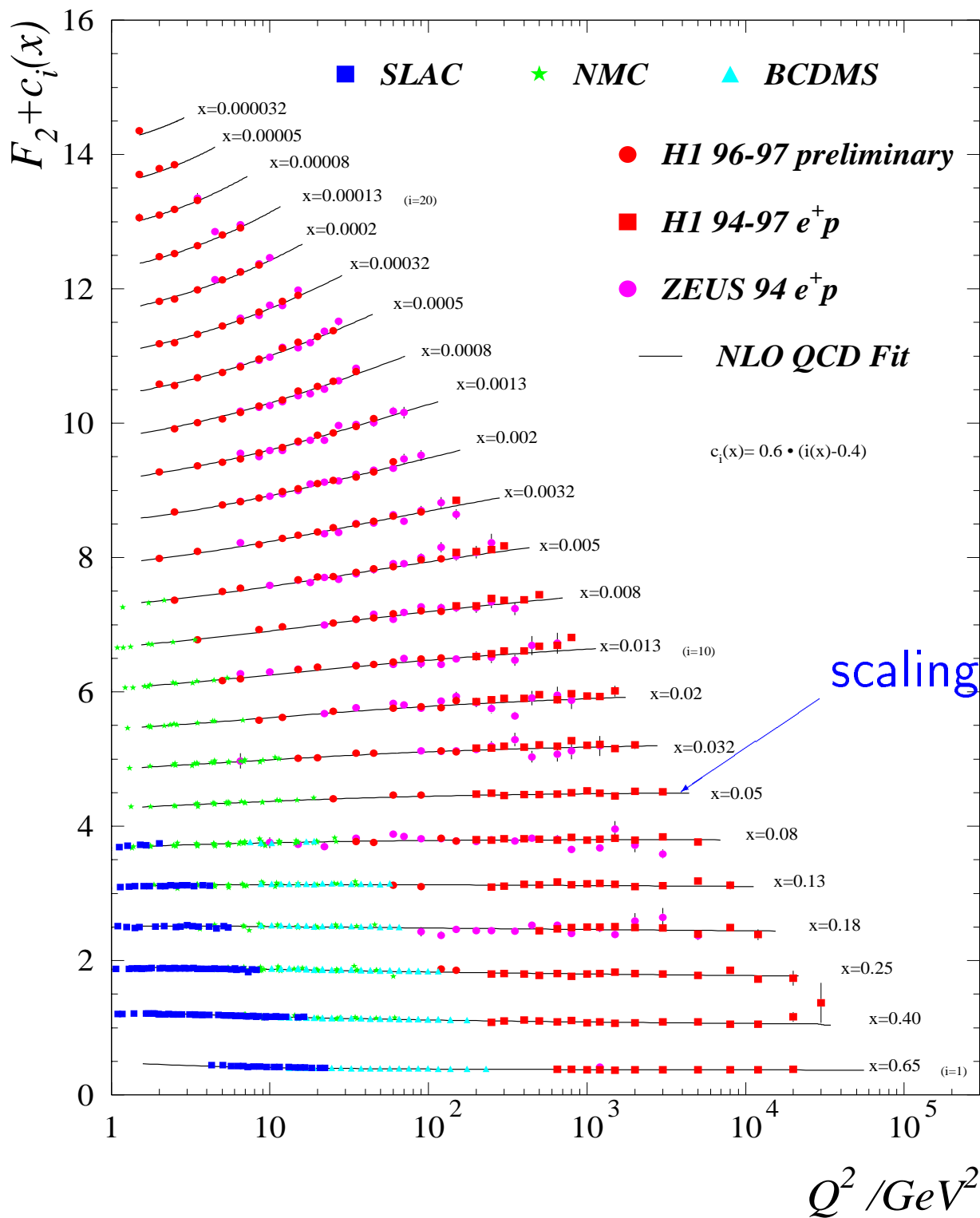
- analogous to renormalization: no dependency of observables on μ_f is required
 \Rightarrow analogue to renormalization group equation:

Dokshitzer-Gribov-Lipatov-Altarelli-Parisi-equation

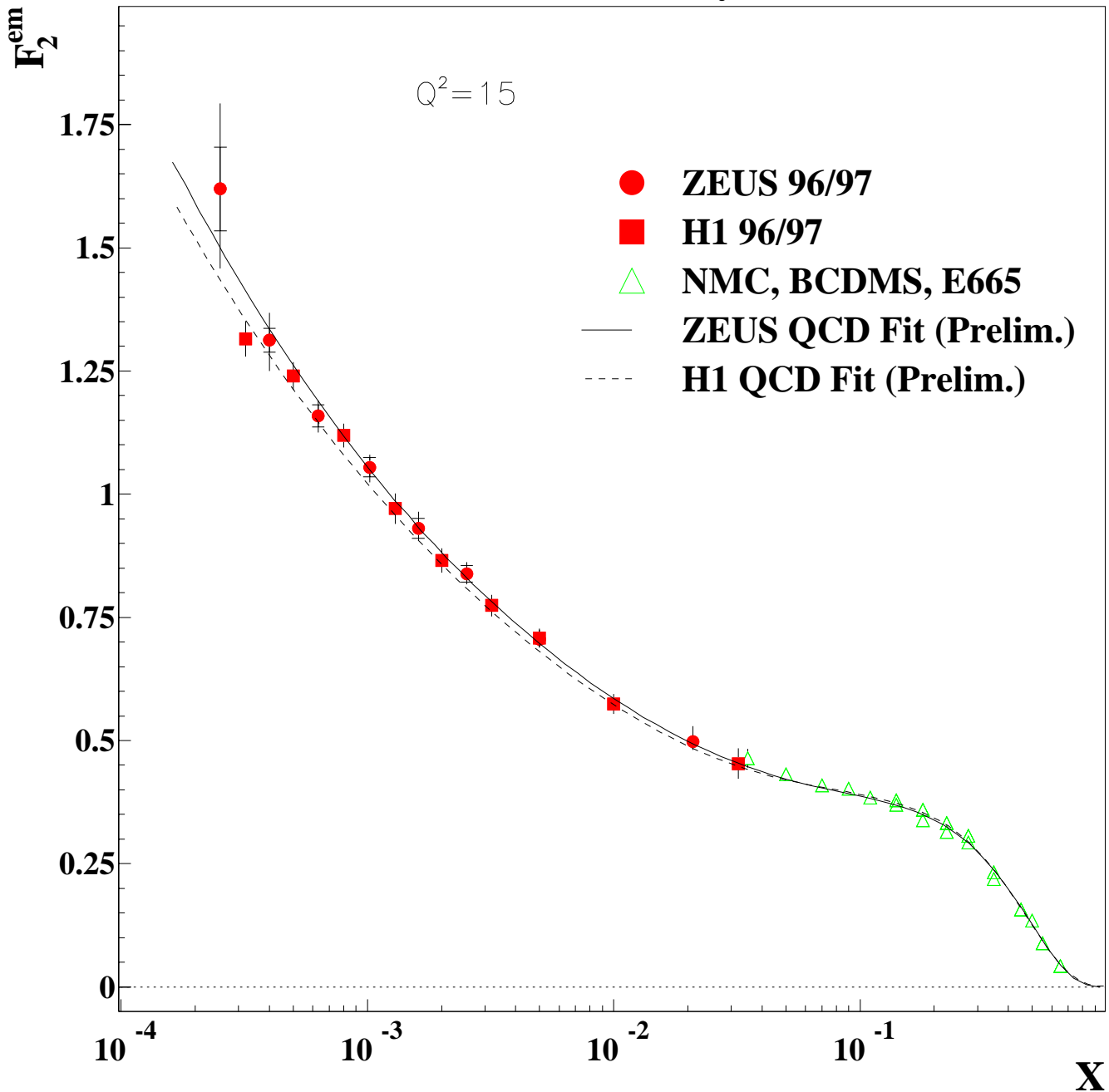
- system of coupled differential equations describing the μ_f dependence of the PDFs
- they include **splitting functions** $P_{ab}^0(z)$: probability density of finding parton a in parton b :



$$\frac{\partial F_2(x, Q^2)}{\partial \ln Q^2} = \frac{\alpha_s(Q^2)}{2\pi} \int_x^1 dz \left[P_{qq}(z) \cdot F_2\left(\frac{x}{z}, Q^2\right) + P_{qg}(z) \cdot 2n_f \cdot \frac{5}{18} \cdot \frac{x}{z} \cdot g\left(\frac{x}{z}, Q^2\right) \right]$$



ZEUS+H1 Preliminary 1996-97



- fixed target experiments measured high x
 \Rightarrow valence quarks
- strong rise of F_2 with decreasing x
 \Rightarrow sea quarks

Summary

- DIS processes at HERA provide information about the partonic structure of the proton
- Quark Parton Model: no parton-parton interaction in the proton → Bjorken Scaling
- QCD: parton interactions lead to logarithmic scaling violation mathematically described by DGLAP-equations
- α_s is a running coupling constant