

Development of a Liquid Argon Purity Monitor for ICARUS

Doktorandenseminar 10.10.01

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Overview of Purity Monitors

The e^- in Argon behaves like free $e^- \rightarrow$ its energy is not enough to ionize Ar!

But

Impurities in Ar (O_2) capture $e^- \rightarrow$ free e^- are lost!

If capture rate of the free e^- is constant

$$N(t_{\text{drift}}) = N(0) e^{-\frac{t_{\text{drift}}}{\tau}}$$

Where $N(t)$ is the number of electrons

t_{drift} is the drift time

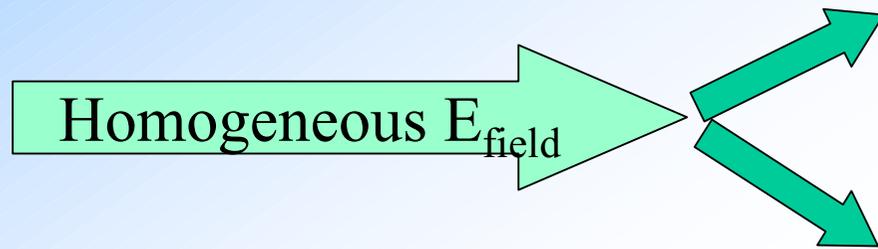
τ is the mean life-time of electrons in LAr

Our Purity Monitors measure τ

Overview of Purity Monitors

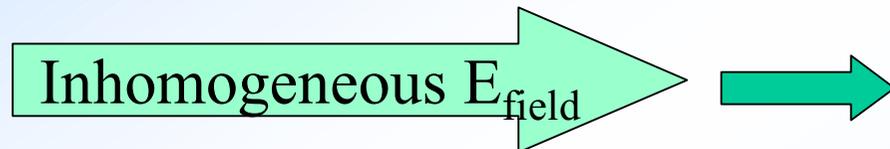
Electric Drift Field

How are e^- produced?



Photoeffect UV \rightarrow Rare Earth

α , β Sources ionize LAr



α , β , fission fragments ionize LAr
(our Purity Monitor)

Overview of Purity Monitors

Homogeneous

Electric Drift Field

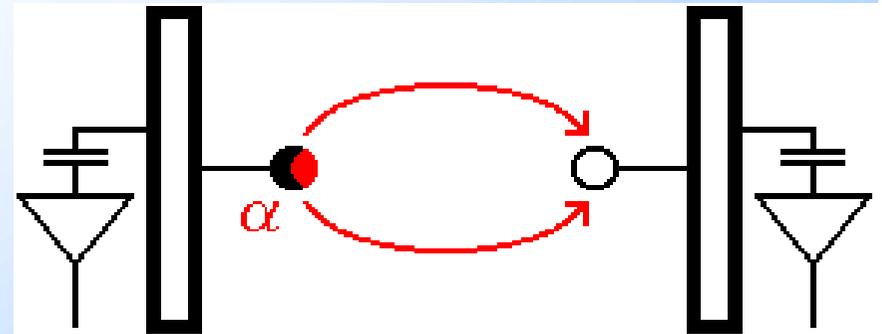
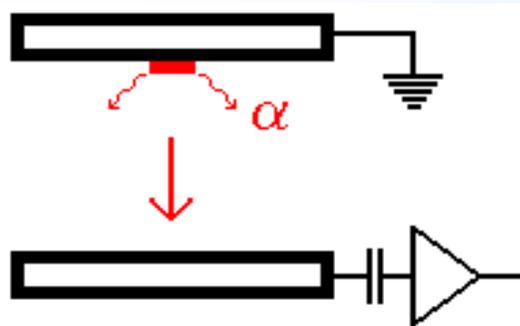
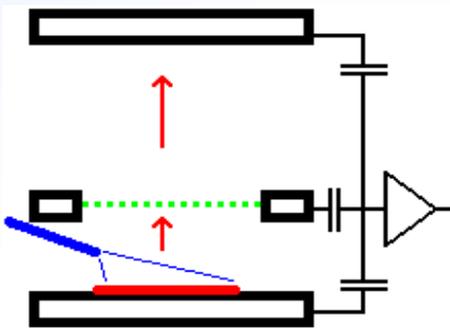
Inhomogeneous

How the free e^- are produced

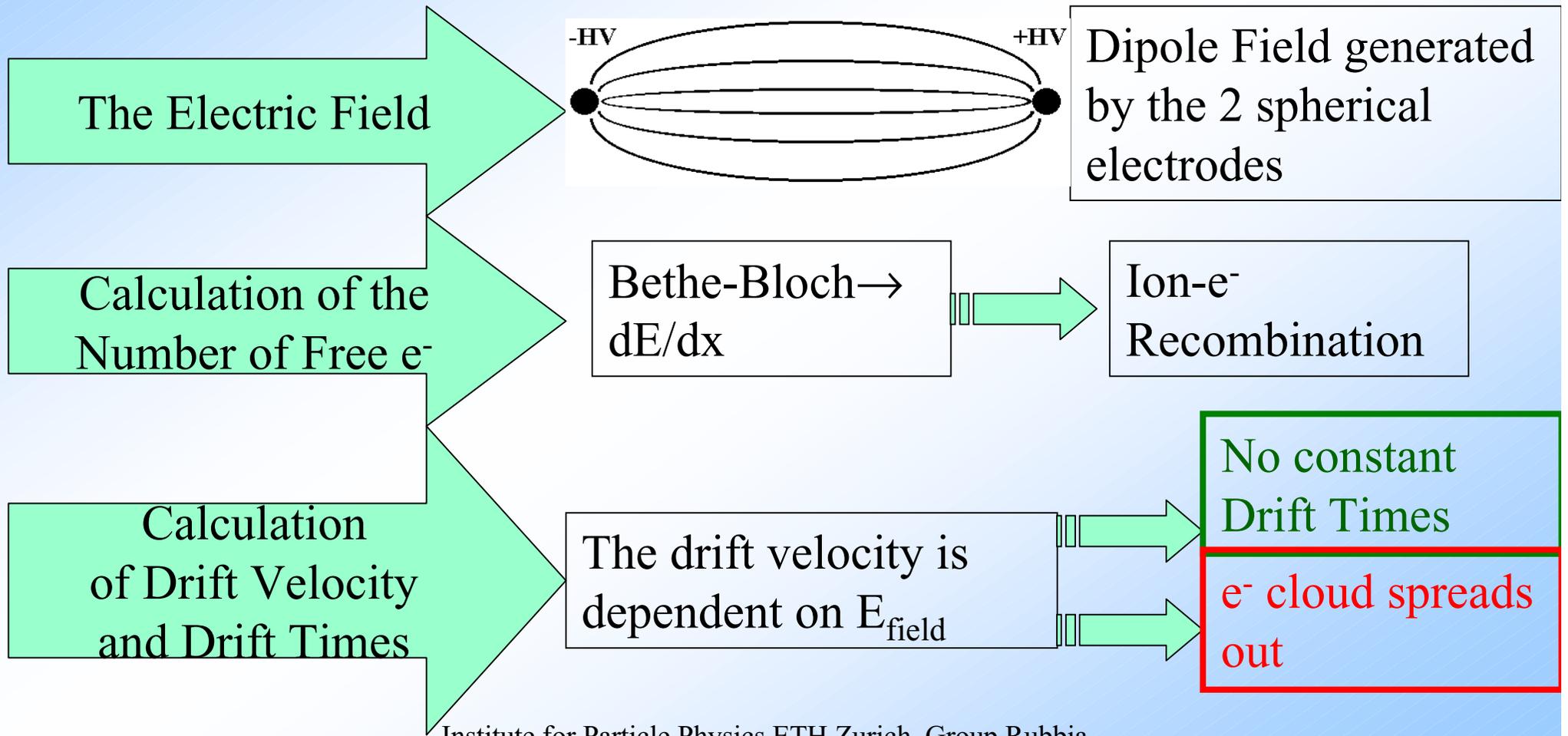
Photoeffect
UV \rightarrow Rare Earth

α , β Sources
ionize LAr

α , β , fission fragments ionize
LAr (**our Purity Monitor**)

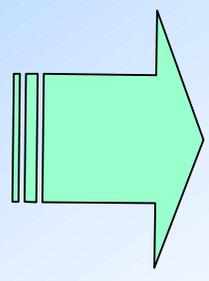
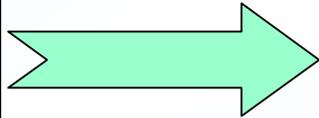


Design Study for the Purity Monitor



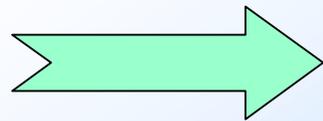
The Electric Dipole Field

Close to the source the E-field is well approximated by a single Coulomb field! \Rightarrow **High field**


$$|\vec{E}| = \frac{V \cdot R}{r^2}$$


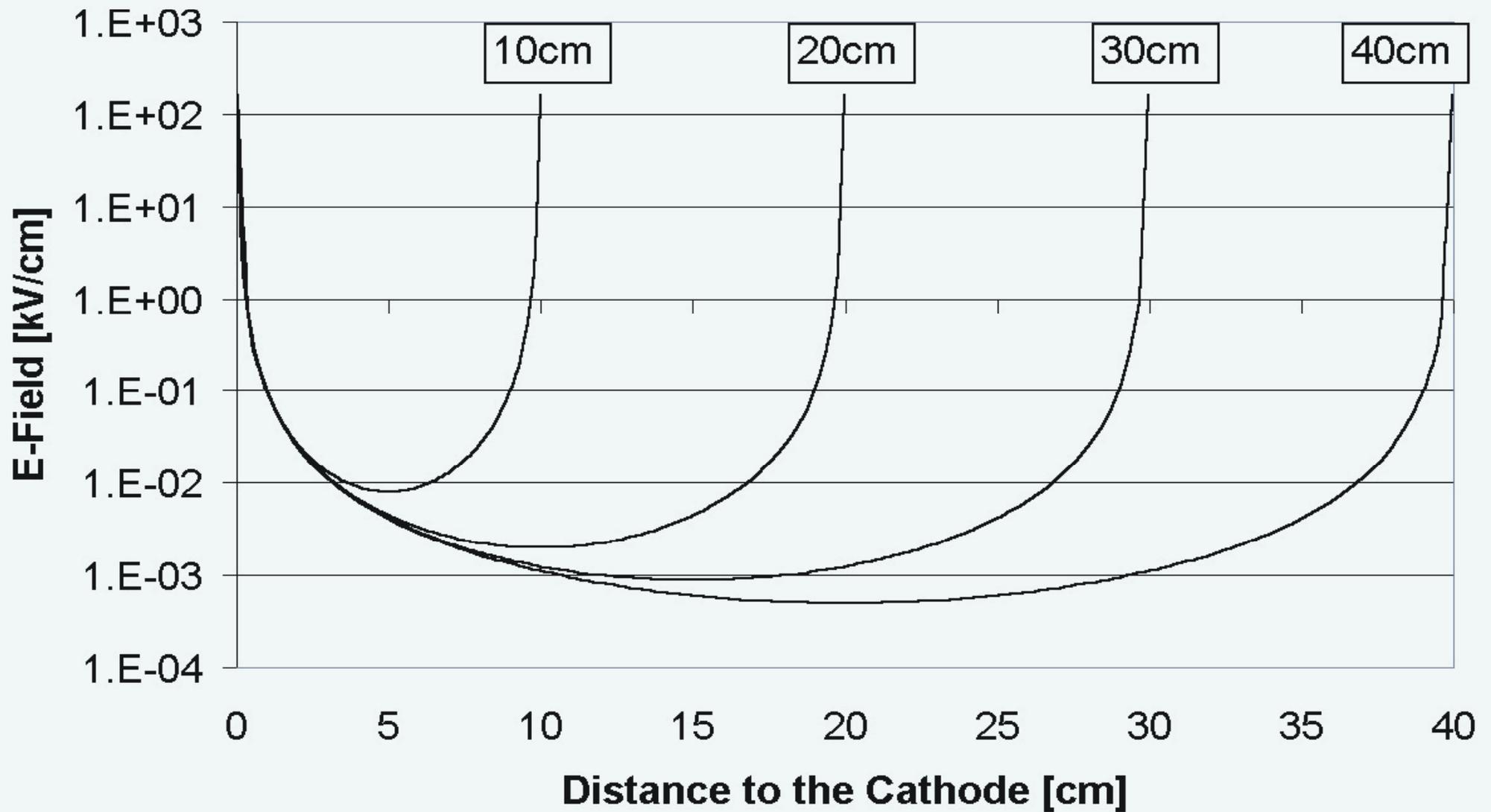
i.e: $V=2\text{kV}$, $R=0.25\text{mm}$,
 $r=R+0.05\text{mm}$
 $\Rightarrow E=56\text{kV/cm}$

Between the two electrodes the E-field is **very low!**

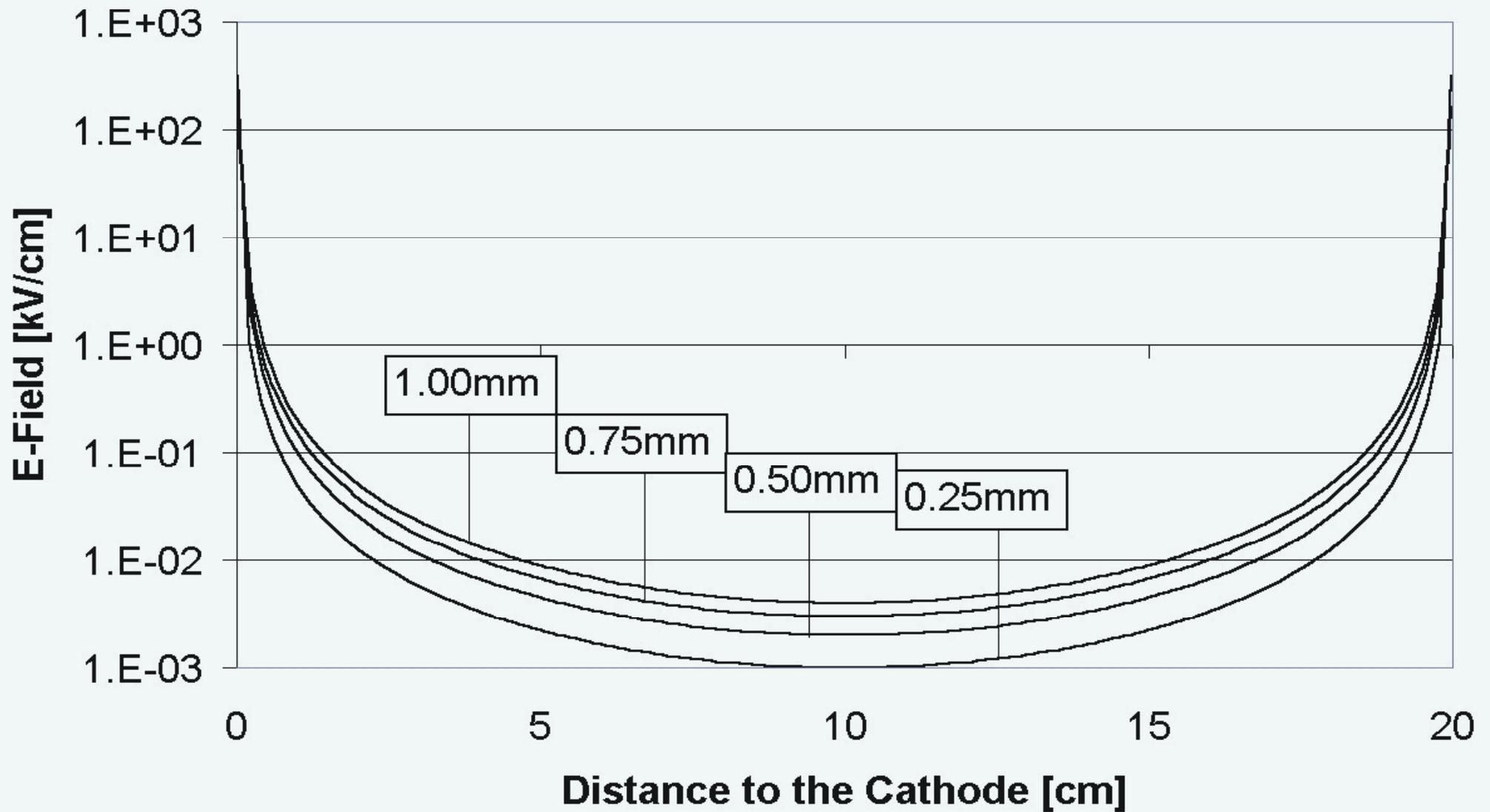


i.e: $V=\pm 2\text{kV}$, $R=0.25\text{mm}$,
 $d=100\text{mm} \Rightarrow E < 5\text{V/cm}$

The Electric Dipole Field



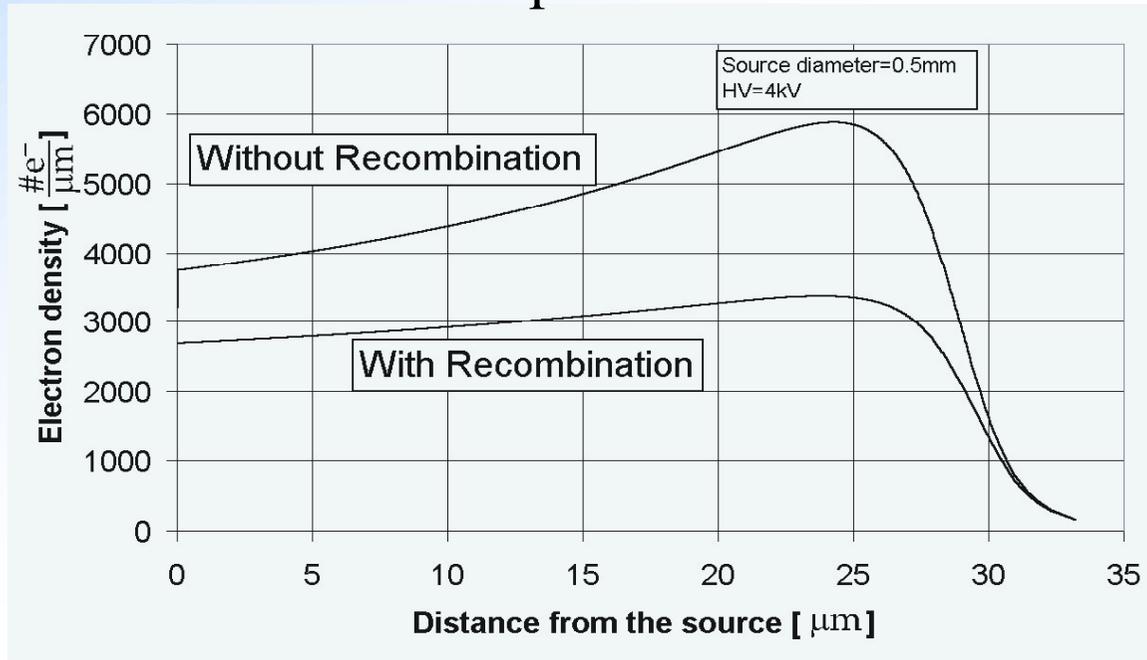
The Electric Dipole Field



Calculation of the Number of Free Electrons

Source: $\text{Pb}^{210} \rightarrow 3.72\text{MeV } \alpha\text{-particle}$

Energy deposition of α -particle in LAr



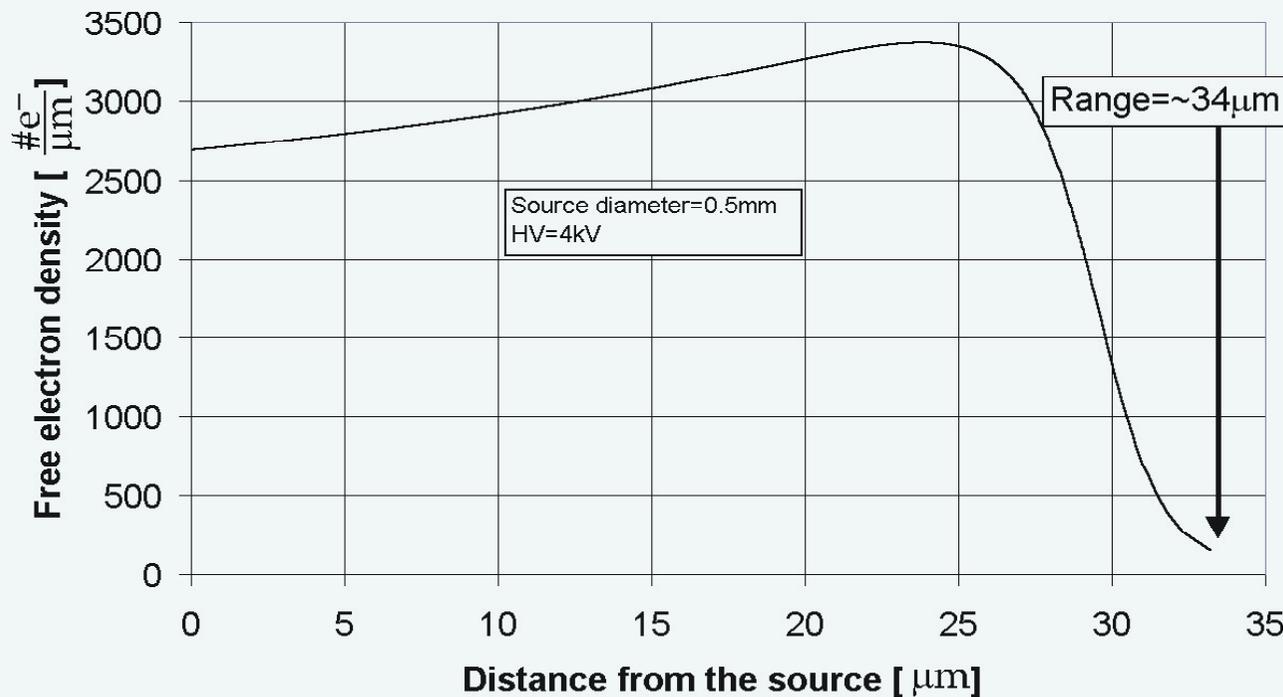
A lot of e^- -ion pairs will recombine due to the electrostatic attraction.

Recombination Models

The α -particles ionize the LAr and produce e^- -ion pairs.

Energy deposition of α -particle in LAr (Bethe-Bloch)

$$\frac{dE}{dx} = 0.1535 \cdot \rho \cdot \frac{Z}{A} \cdot \frac{z^2}{\beta^2} \cdot \left[2 \ln \left(2.511000 \frac{\gamma^4 \beta^2}{I_0} \right) - 2 \cdot \beta^2 \right]$$



The α -particles produce a very large ionization density (700 times MIP)

➡ Small dimension of the electron cloud

➡ Bigger Recombination effect!

Recombination Models

Birks Model

Semi-empirical formula for the recombination at a given E-field. It takes into account the ionization density of the particle

$$\frac{dN}{dx} = \frac{\frac{dE}{dx} \cdot \frac{1}{w}}{1 + k_1 \left(\left| \vec{E} \right| \right) \frac{dE}{dx} \cdot \frac{1}{\rho}}$$

$w=23.6\text{eV}$ is the mean energy needed to produce an e⁻-ion pair

k_1 is E-field dependent

Box Model

The e⁻-ion pairs are considered as isolated and, at the beginning, the distribution is uniform in a box of certain dimensions

$$\frac{N}{N_0} = \frac{\left| \vec{E} \right|}{C} \ln \left(1 + \frac{C}{\left| \vec{E} \right|} \right)$$

N, N_0 are the number of electrons with and without recombination

C depends on source and medium

Calibration of Models

Find C on Box Model, given $k_1(500\text{V/cm})$

Integrate the Birks formula with the known $k_1(500\text{V/cm})$ and in a homogeneous E-field (500V/cm)

$$\frac{dN}{dx} = \frac{\frac{dE}{dx} \cdot \frac{1}{w}}{1 + k_1 \left(\left| \vec{E} \right| \right) \frac{dE}{dx} \cdot \frac{1}{\rho}}$$

$w=23.6\text{eV}$ is the mean energy needed to produce an e-ion pair

k_1 is E-field dependent

Integrate the Bethe-Bloch

N_0

Calculate N/N_0

Integrate the Birks formula with a k_1 , with the Box model, find the correspondent E-field

$$\frac{N}{N_0} = \frac{\left| \vec{E} \right|}{C} \ln \left(1 + \frac{C}{\left| \vec{E} \right|} \right)$$

N, N_0 are the number of electrons with and without recombination

C depends on source and medium

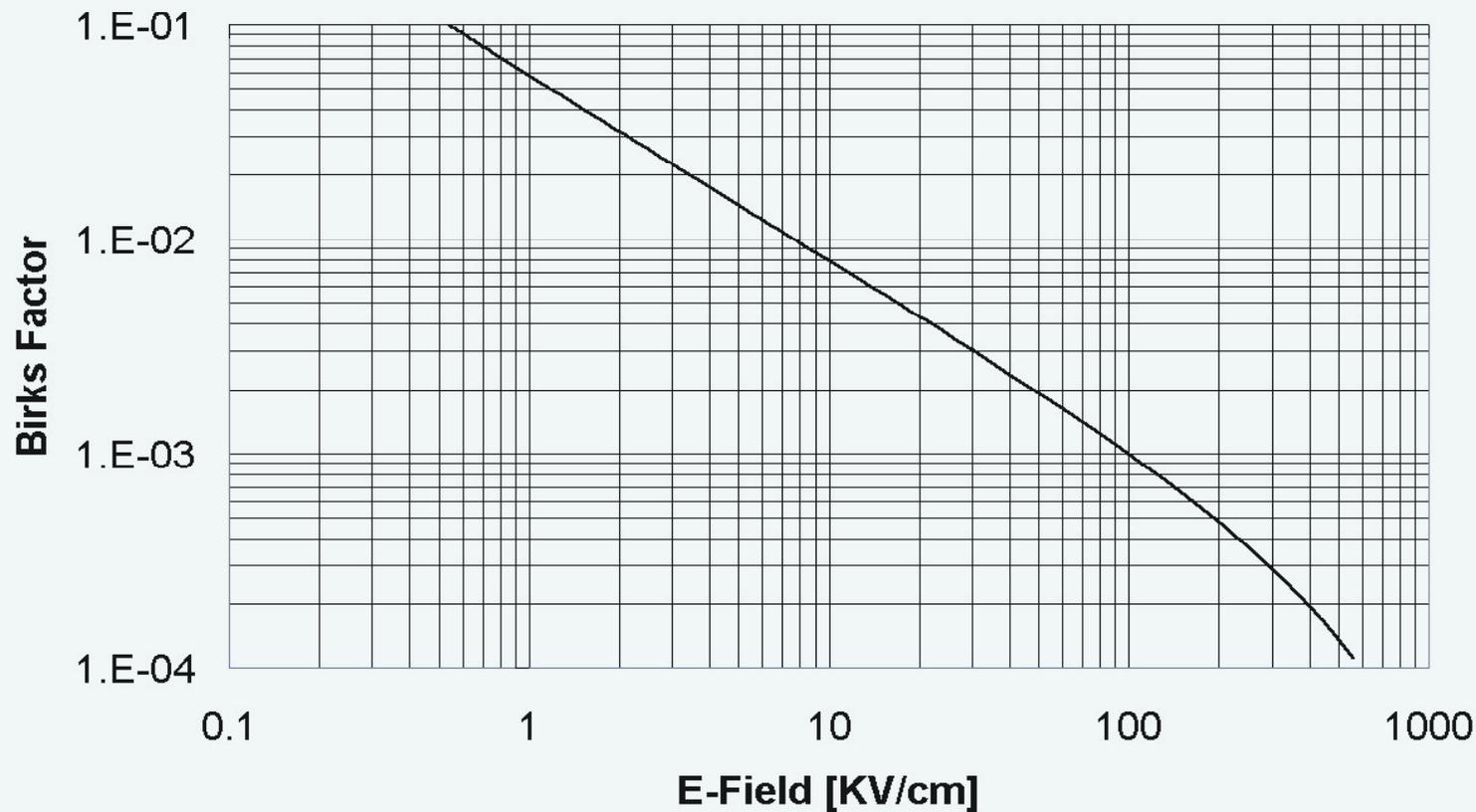
Find k_1 on Birks Model, given C

Recombination Models: Birks

Model

$$\frac{dN}{dx} = \frac{\frac{dE}{dx} \cdot \frac{1}{w}}{1 + k_1 \left(\left| \vec{E} \right| \right) \frac{dE}{dx} \cdot \frac{1}{\rho}}$$

$w=23.6\text{eV}$ is the mean energy needed to produce an e^- -ion pair



k_1 is known for LAr at 500V/cm

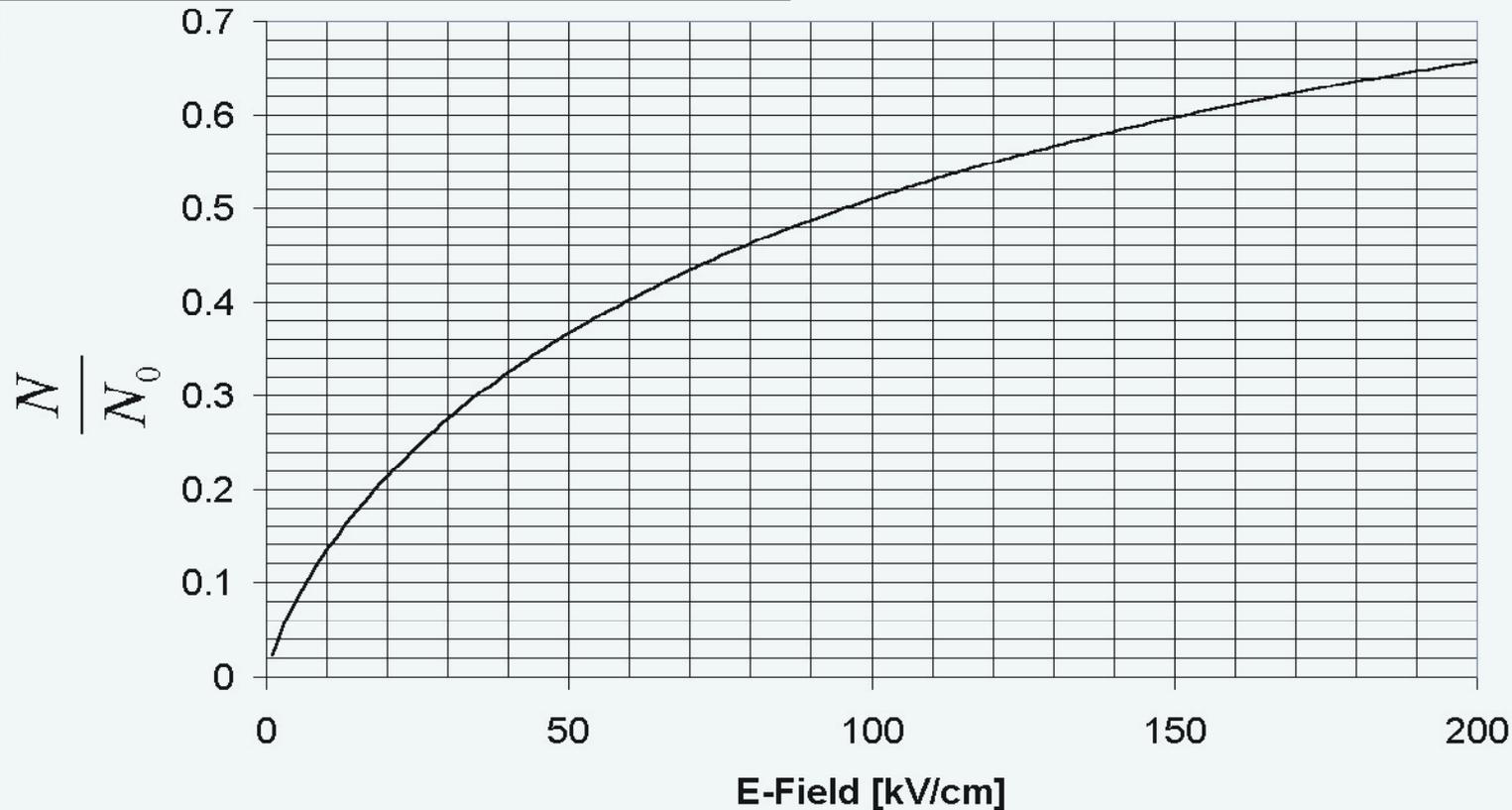
Calibration

k_1 is found with the Box Model

Recombination Models: Box Model

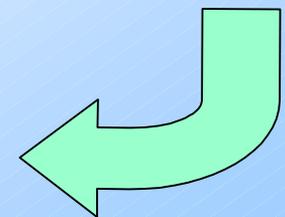
$$\frac{N}{N_0} = \frac{|\vec{E}|}{C} \ln \left(1 + \frac{C}{|\vec{E}|} \right)$$

N, N_0 are the number of electrons with and without recombination



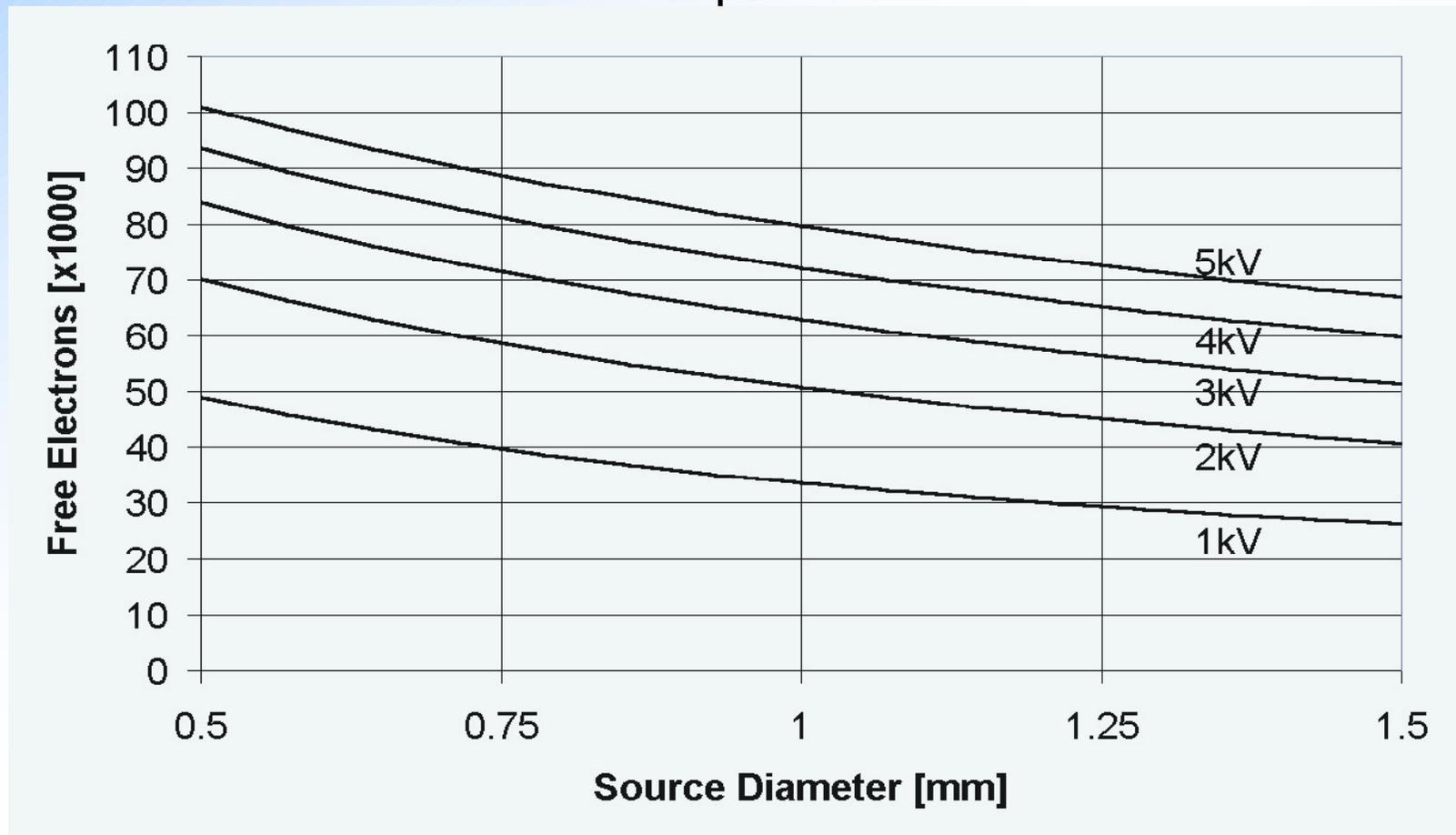
C for LAr and Pb²¹⁰
 α -particle=239.9

The Box model gives the recombination in a homogeneous E-field.



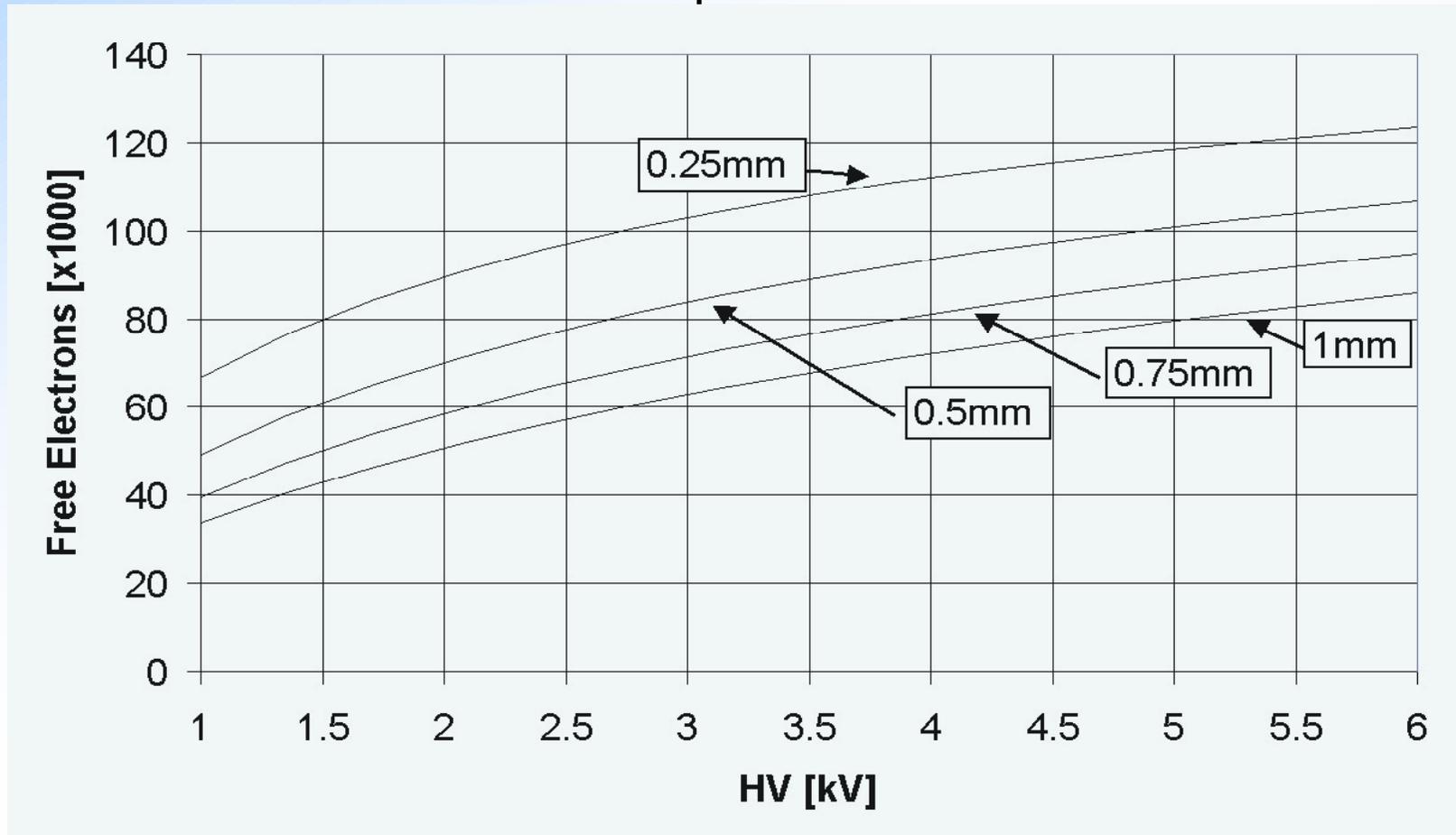
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Calculation of the Number of Free Electrons

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Calculation of Drift Velocity and Drift Times

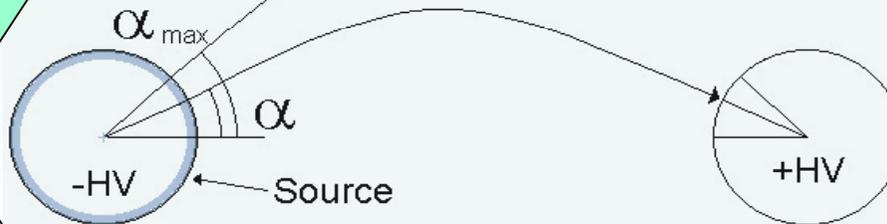
- The e^- are tracked along the E-field lines with a computer program.
- The drift time should be comparable to the mean life-time for a precise life-time measurement.

$t_{\text{drift}}(\text{HV}, d, l, \alpha)$

Field Shaping Electrodes



$$V_d \simeq P_3 \cdot \left| \vec{E} \right| \ln \left(1 + \frac{P_4}{\left| \vec{E} \right|} \right) + P_5 \cdot \left(\left| \vec{E} \right| \right)^{P_6}$$



e^- - cloud diffusion

$$\delta_{\text{diff}} = \sqrt{2 D \cdot t}$$

Where $D \simeq 5 \text{ cm}^2/\text{s}$

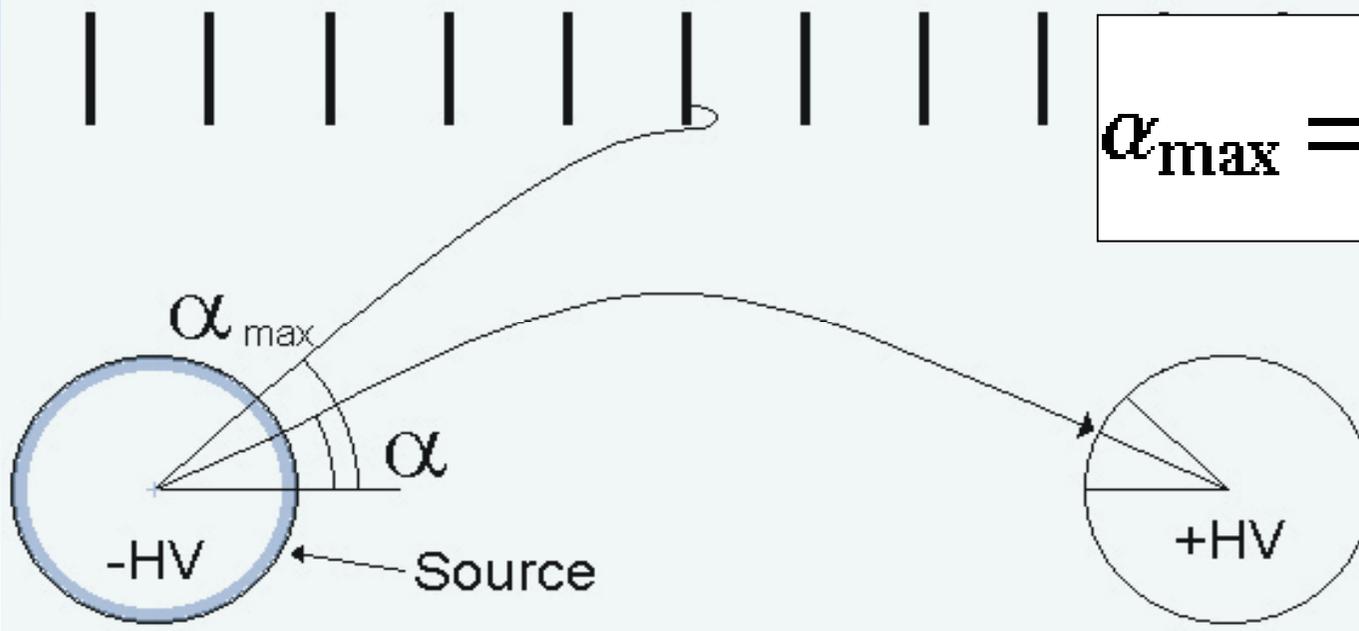
$$\alpha_{\text{max}} = 31.4^\circ \cdot \frac{D_{\text{Chamber}}}{l}$$

Calculation of Drift Velocity and Drift Times

The different drift times are given by the starting angle (α)

→The diameter of the Field Shaping Electrodes gives a limit on possible drift times

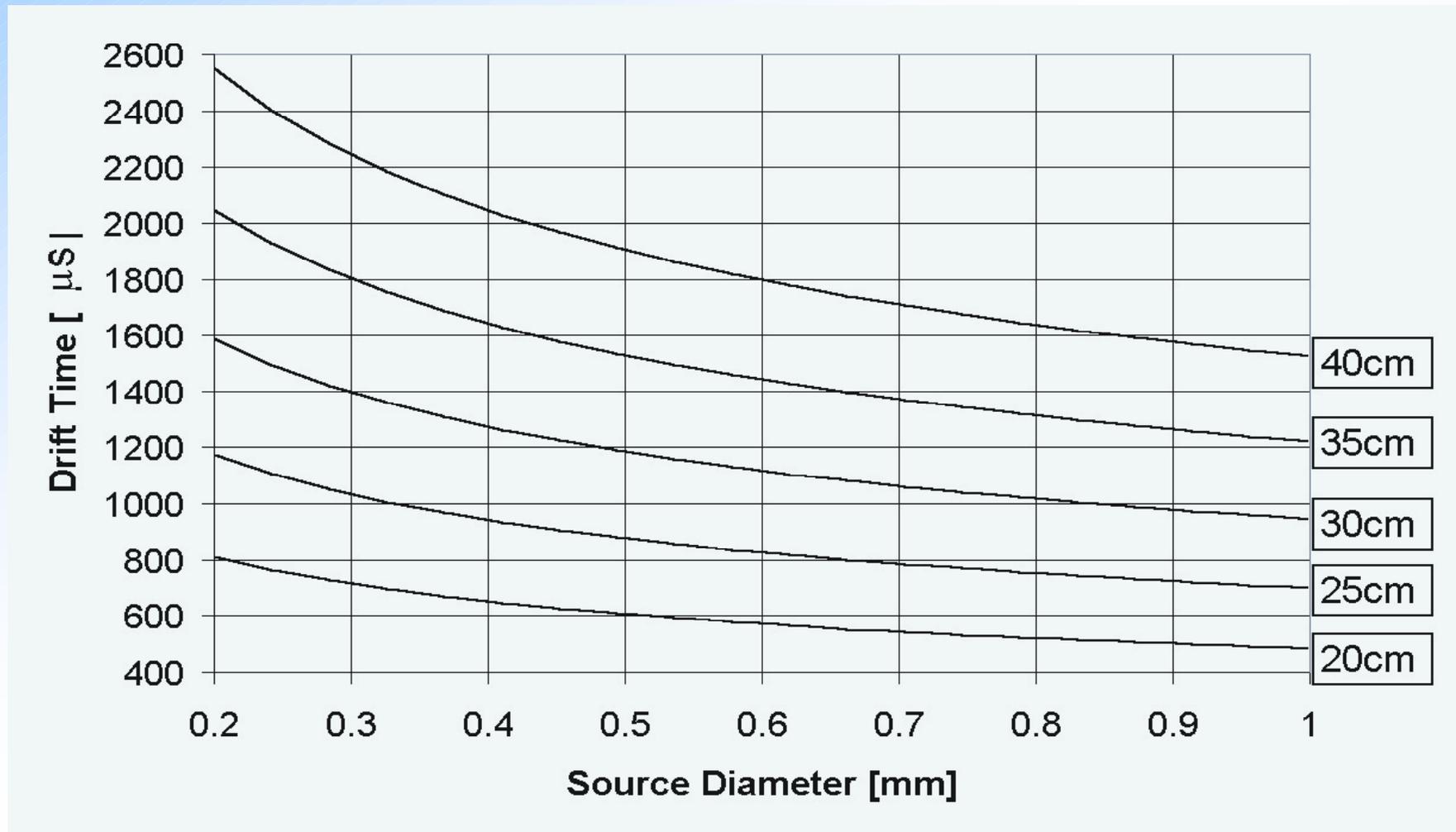
Field Shaping Electrodes



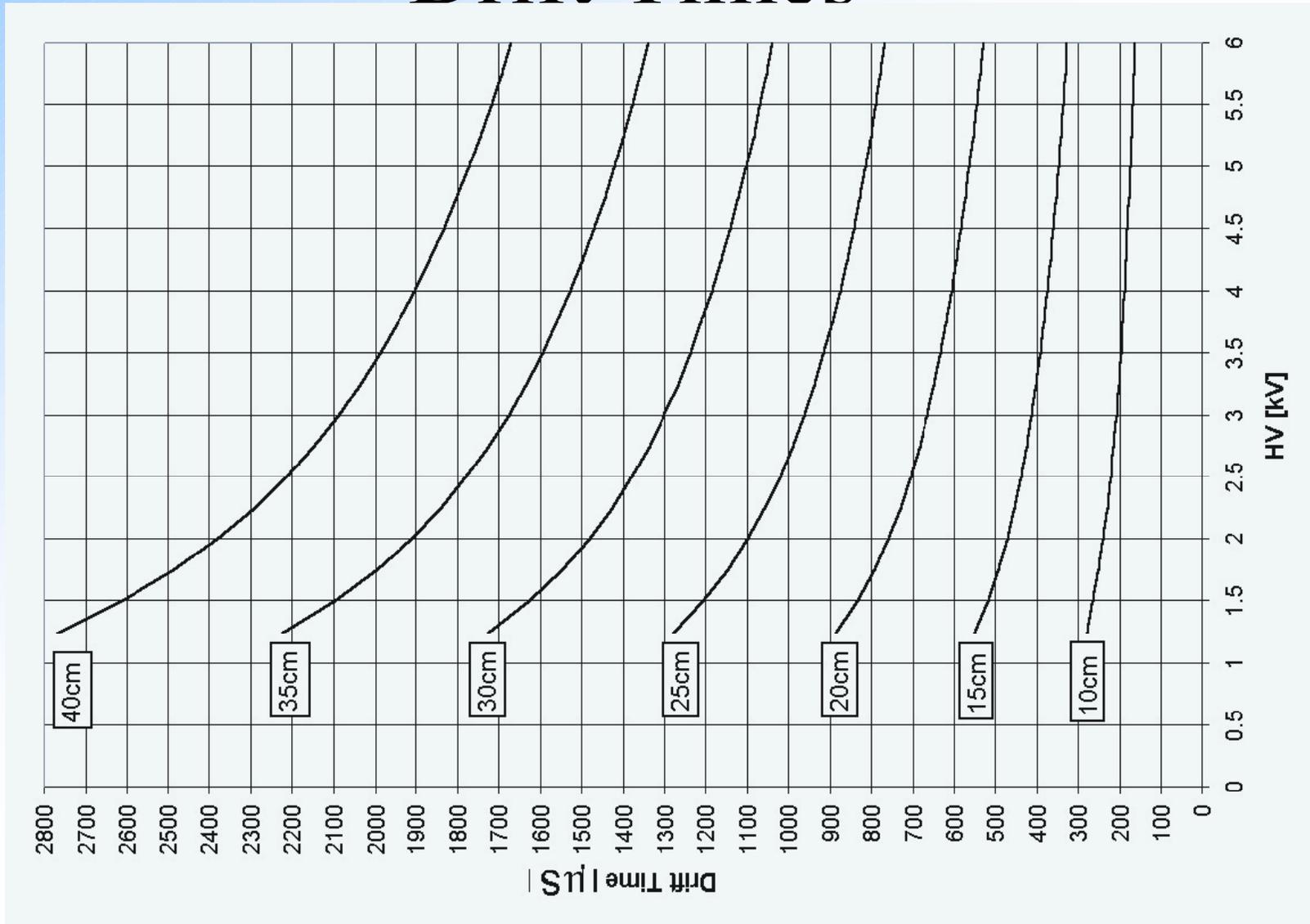
$$\alpha_{\max} = 31.4^{\circ} \cdot \frac{D_{\text{Chamber}}}{l}$$

Empirical
formula

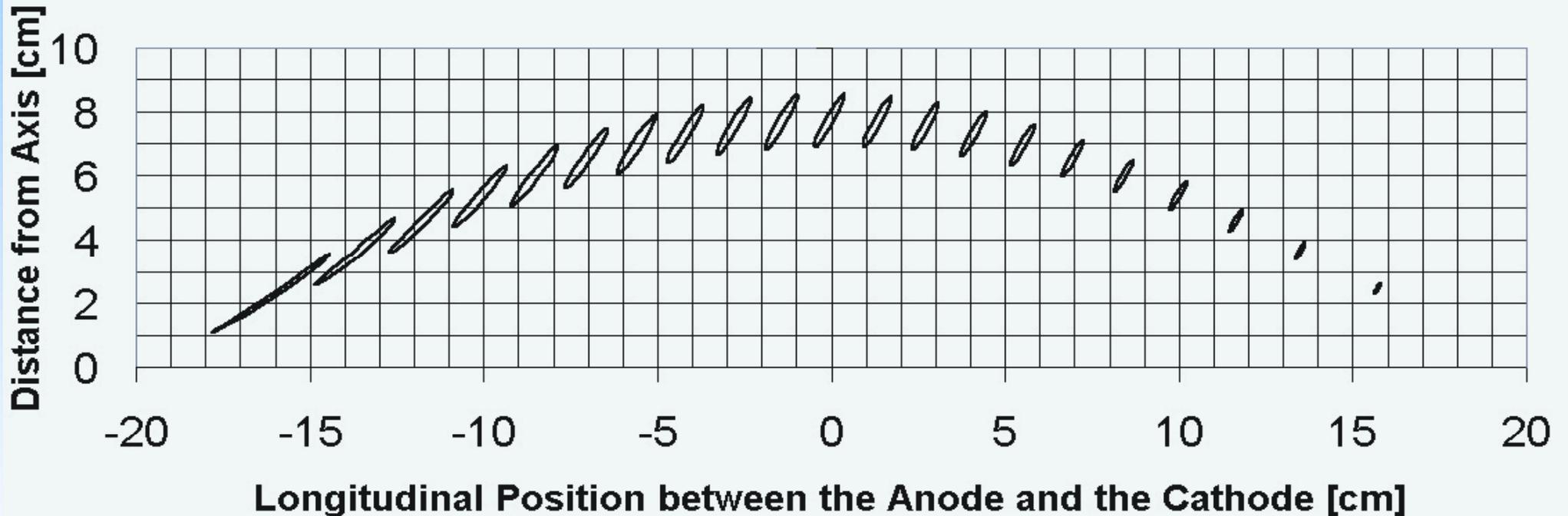
Calculation of Drift Velocity and Drift Times



Calculation of Drift Velocity and Drift Times

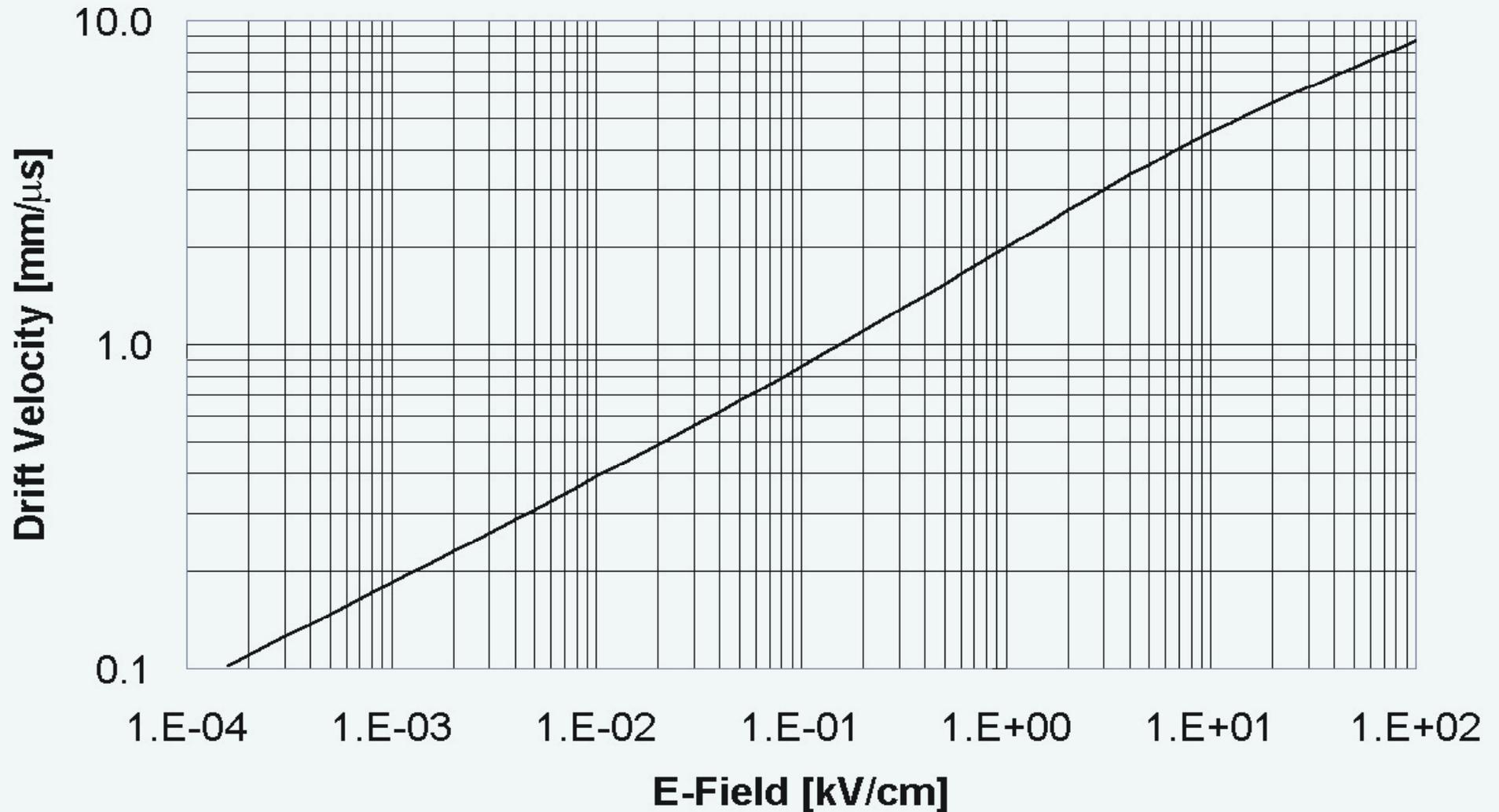


Calculation of Drift Velocity and Drift Times

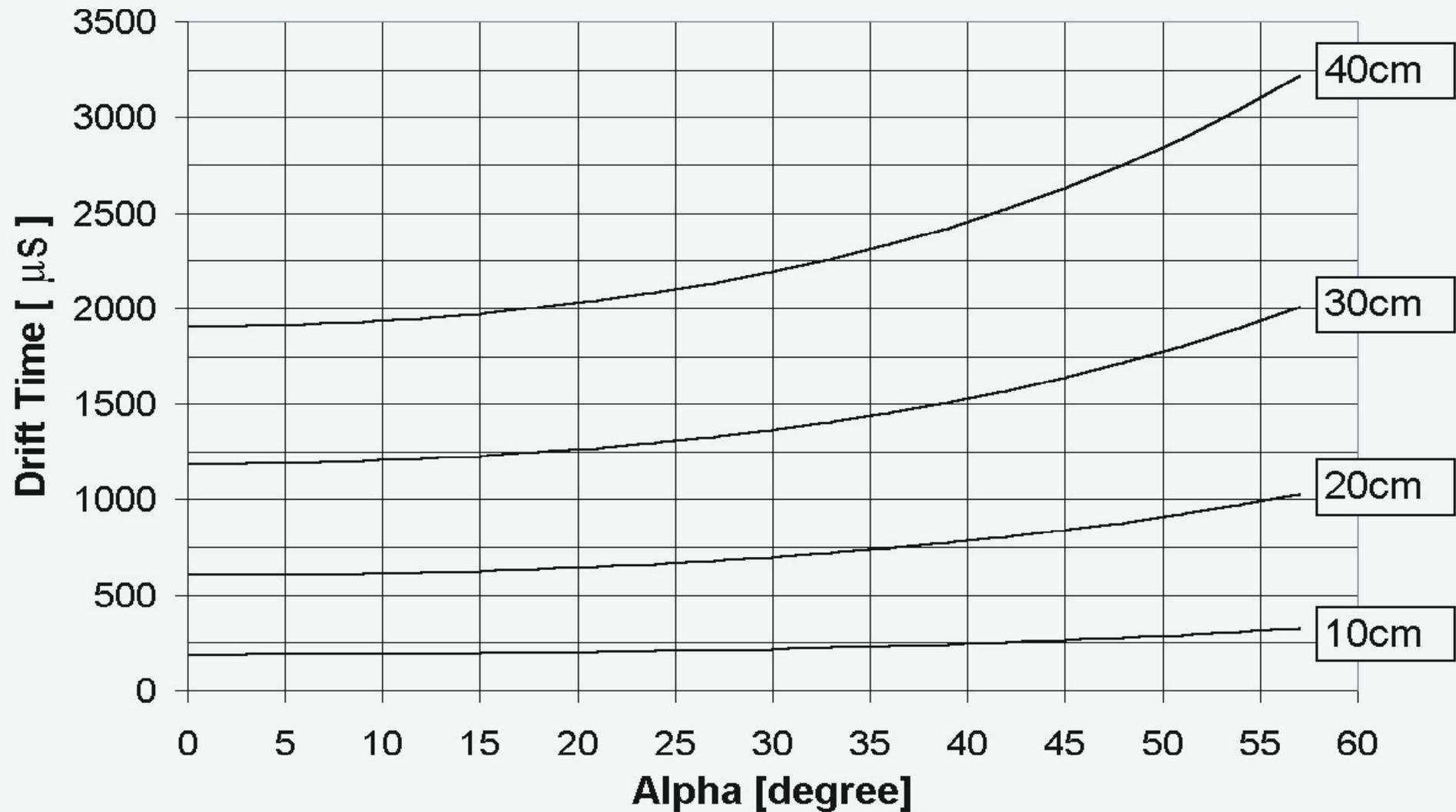


“Stroboscopic” view of an e^- cloud moving in a dipole field. The contour is 1σ of the (almost) Gaussian e^- distribution. Time intervals of $100 \mu\text{s}$ are shown.

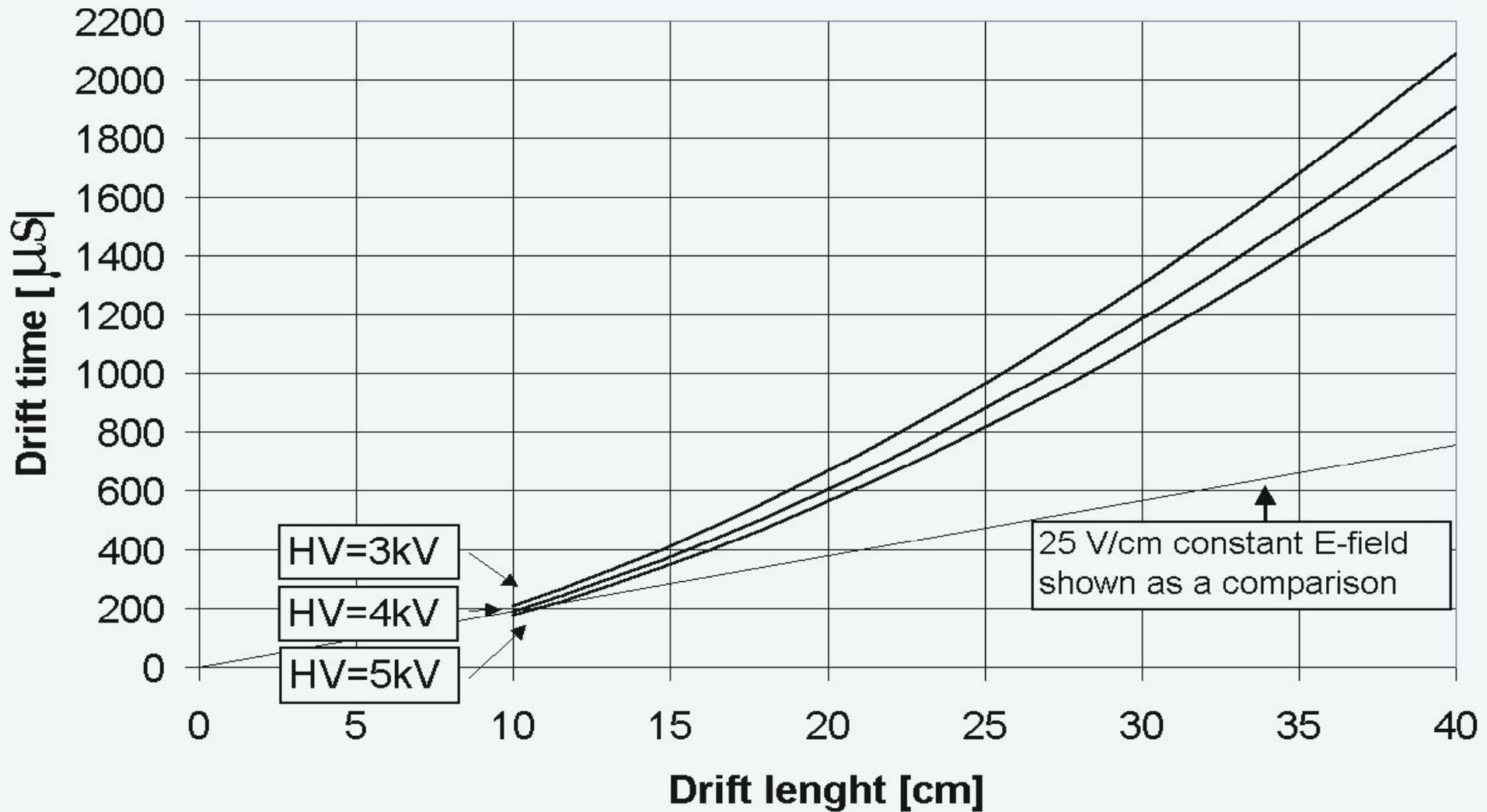
Calculation of Drift Velocity and Drift Times



Calculation of Drift Velocity and Drift Times



Calculation of Drift Velocity and Drift Times



Experimental Setup

Purity Monitor
Mechanics

Materials for high purity, vacuum
and cryogenic temperature

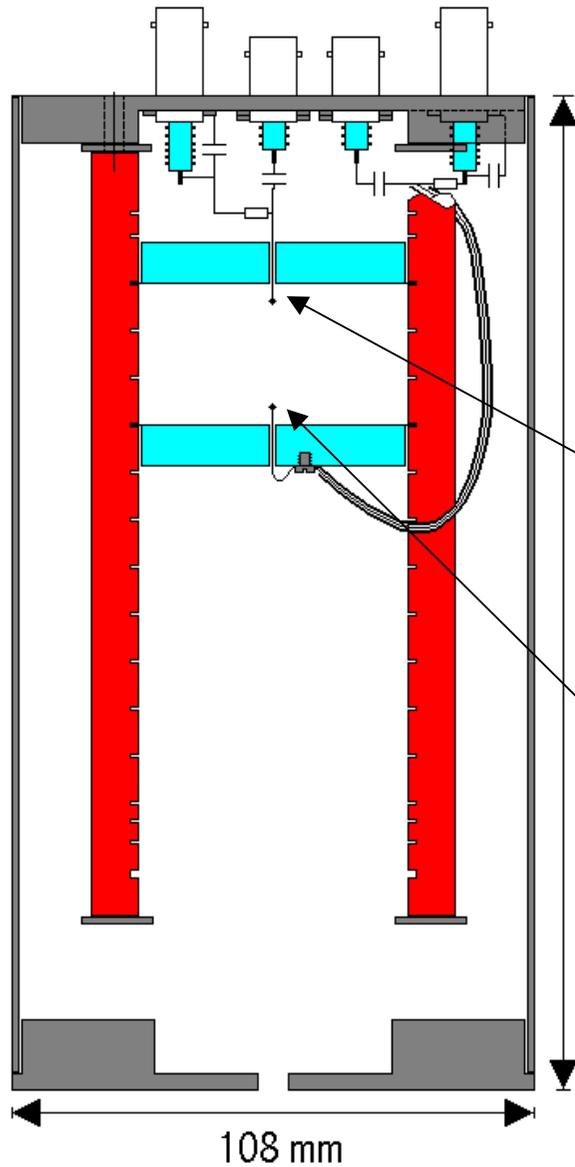
Electronics
and DAQ

10k e⁻ preamp and filters to have
best S/N. DAQ card @20Ms/s

Cryogenic
Vessel

High purity LAr (CuO filter).
Materials for purity, vacuum
and cryogenic temperature

Purity Monitor Mechanics

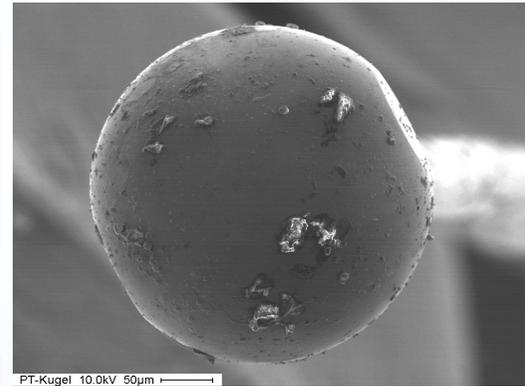


- Stainless Steel
- Polyethylene
- Macor

Materials for high purity, vacuum and cryogenic temperature

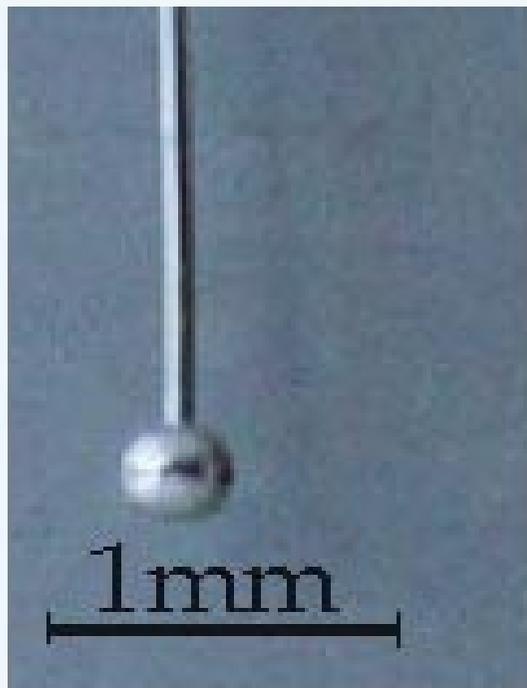
Spherical anode

Pb^{210} spherical source



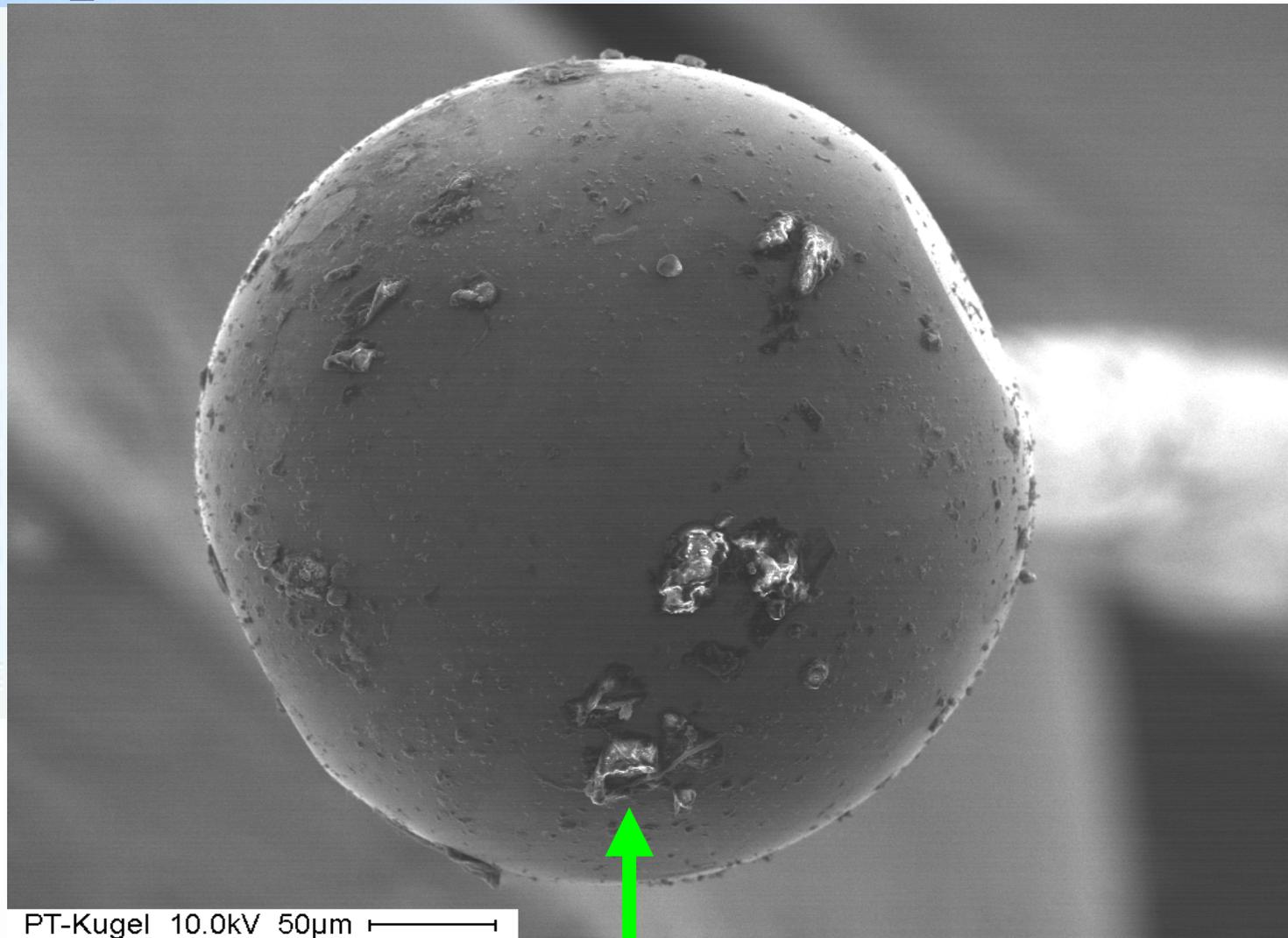
Built by our group

Spherical Anode



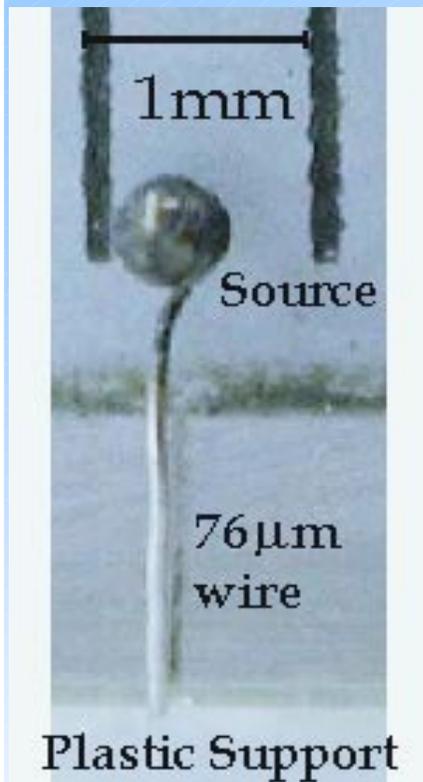
A Pt-wire is melted
in a flame

⇒spherical by
superficial tension



Plastic particle coming from a plastic envelope

Spherical Pb^{210} α -Source



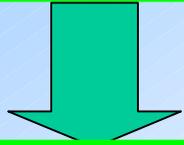
A Pt-wire is melted in a flame



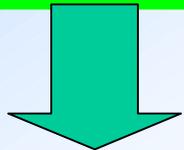
\Rightarrow The source is deposited by an electro-chemical procedure

Spherical Pb^{210} α -Source

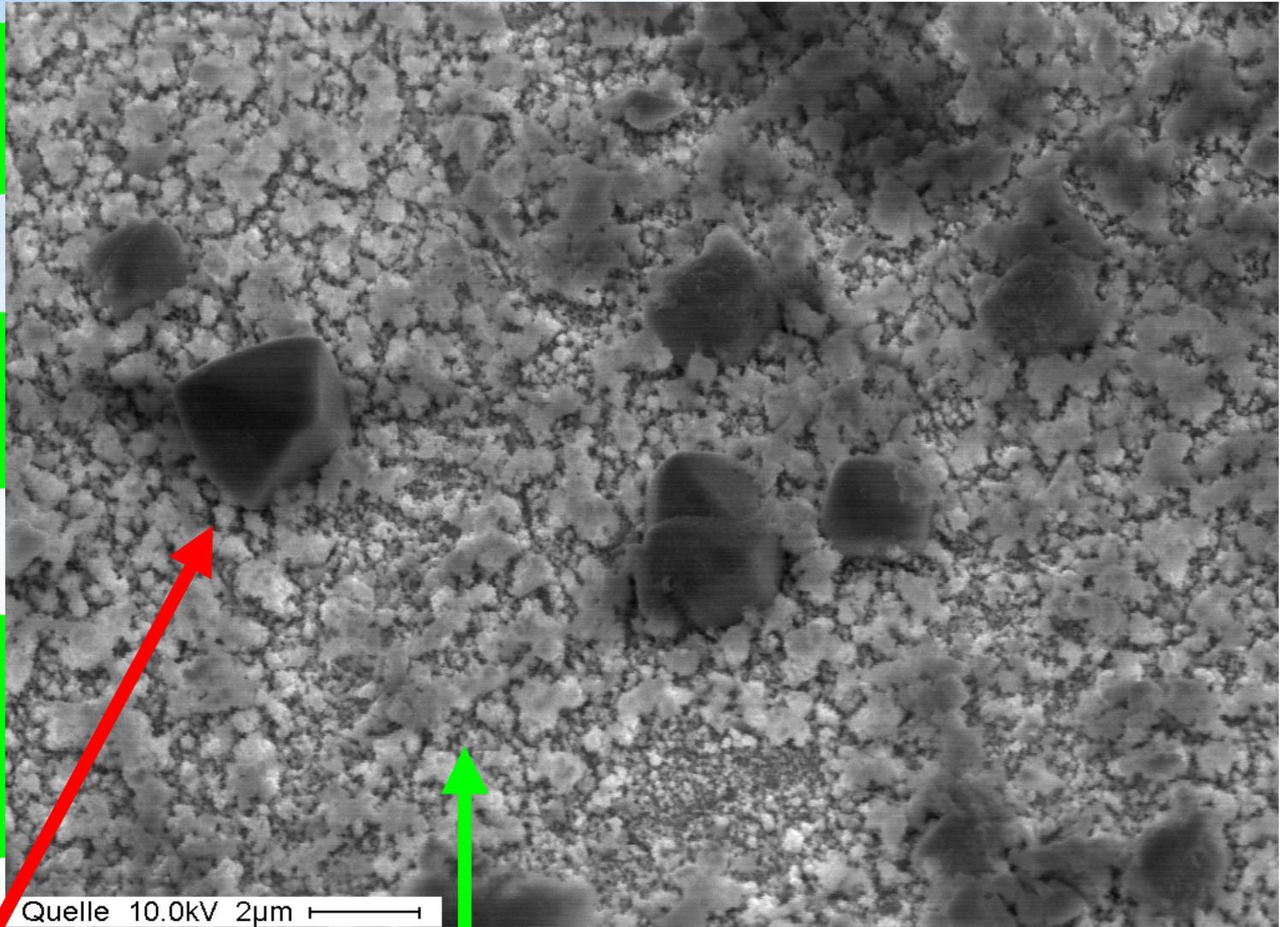
1.2Mol HNO_3 Pb^{210}
solution $\sim 20\text{kBq}$



The solution is dried
out



Dissolved in 1Mol
Acetic acid and
1Mol NH_4 -acetate



Unknown crystals

Pb crystals grown at the Pt surface

Spherical Pb²¹⁰ α-Source

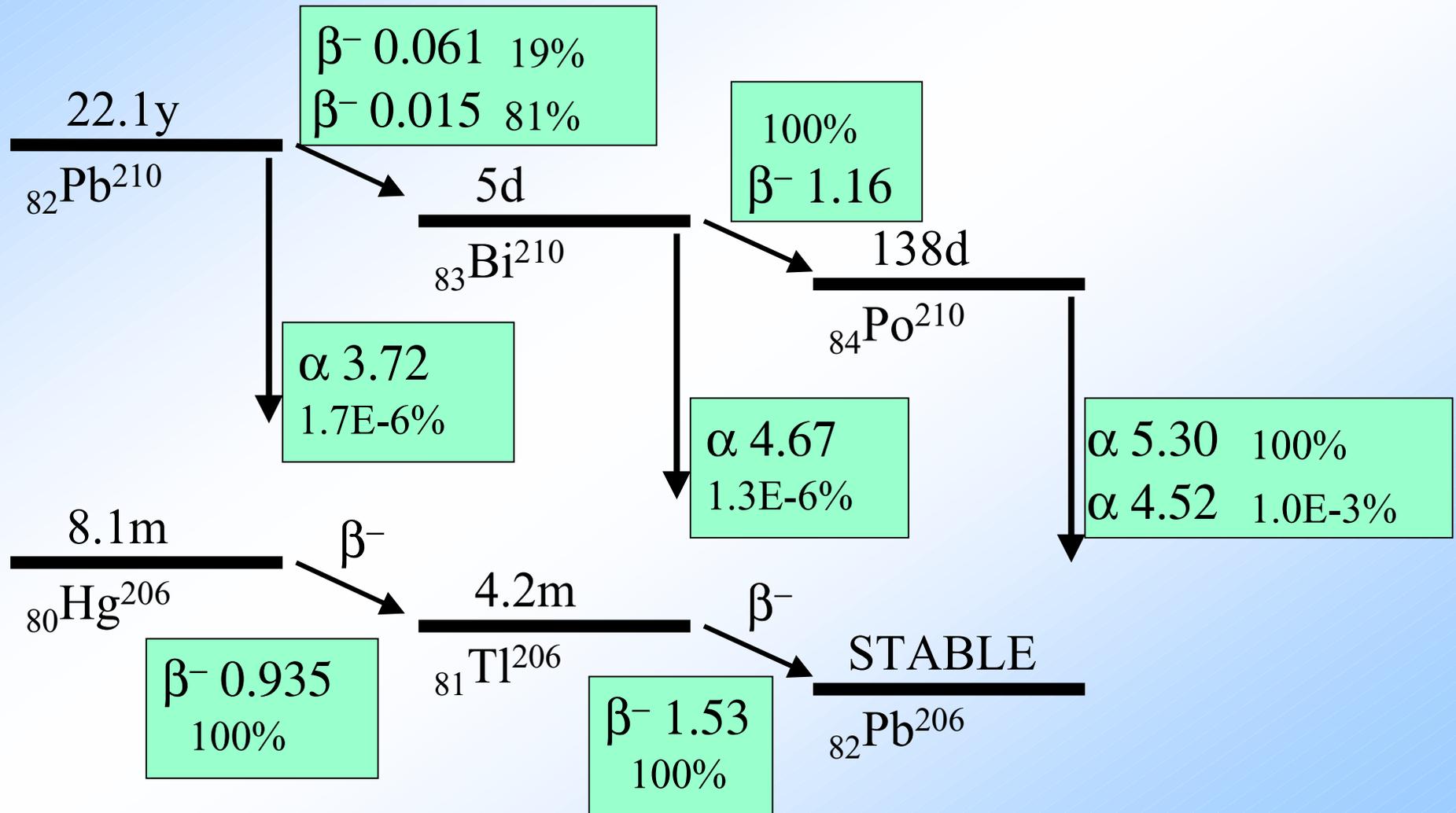
Isotope	Life-time	Modes of Decay	Particle Energy [Mev]	Particle Intensities
⁸² Pb ²¹⁰	21y	β	0.015	0.81
			0.061	0.19
		α	3.72	1.7E-8
		γ	~0.047	With β
Isotope	Life-time	Modes of Decay	Particle Energy [Mev]	Particle Intensities
⁸⁰ Hg ²⁰⁶	8m	β ⁻	0.935	1.00
		γ	0.305	With β ⁻
Isotope	Life-time	Modes of Decay	Particle Energy [Mev]	Particle Intensities
⁸¹ Tl ²⁰⁶	4m	β ⁻	~1.530	1.00
		γ	1.163	6.2E-4

Spherical Pb²¹⁰ α-Source

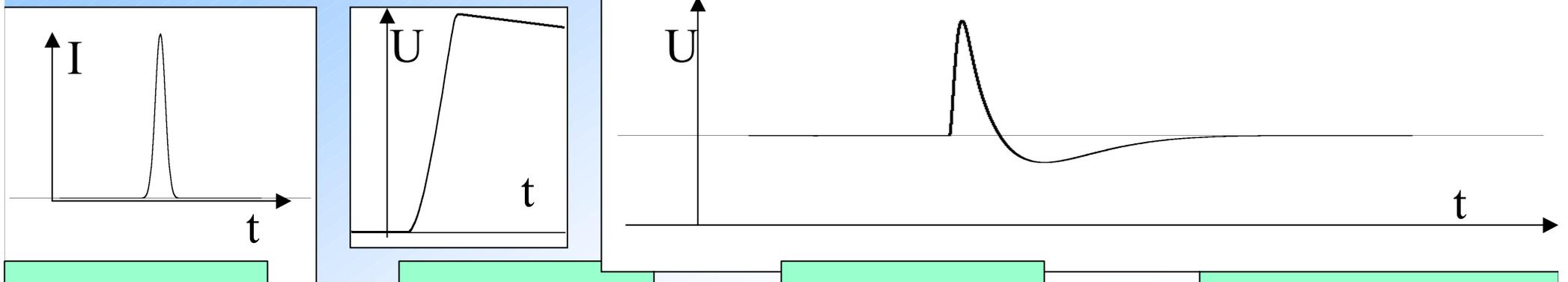
Isotope	Life-time	Modes of Decay	Particle Energy [Mev]	Particle Intensities
⁸³ Bi ²¹⁰	5d	β ⁻	1.16	1.00
		α	~4.67	1.3E-6
		γ	~0.3	With α

Isotope	Life-time	Modes of Decay	Particle Energy [Mev]	Particle Intensities
⁸⁴ Po ²¹⁰	138d	α	5.30	1.00
		α	4.52	1E-5
		γ	0.80	1.6E-6

Spherical Pb²¹⁰ α-Source



Electronics and DAQ



Charge to
Voltage
Preamp

Filter to
optimize
S/N

Offset and
Amplifier

DAQ Computer
Card

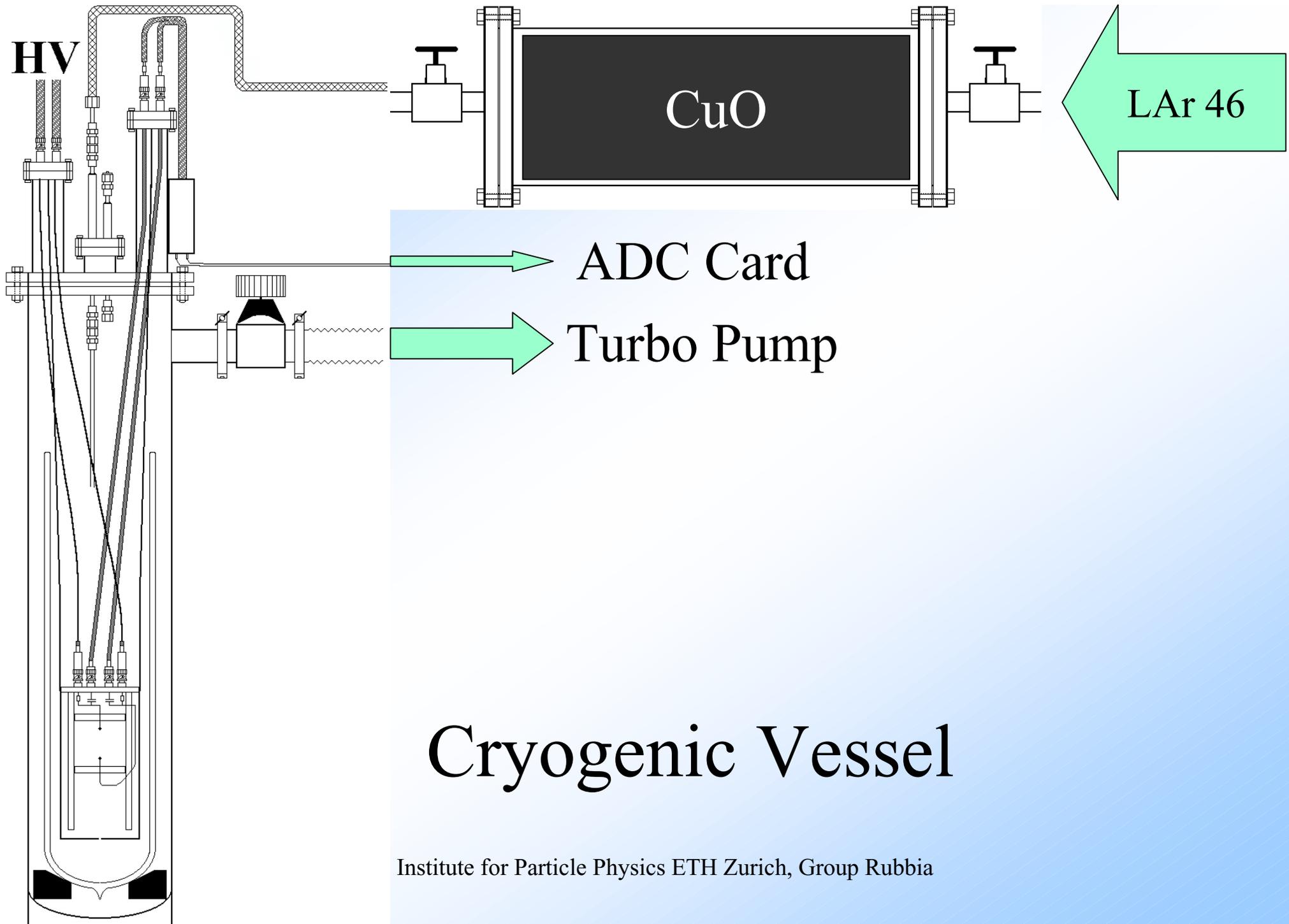
Preamp developed
for ICARUS

→ Sensitivity
of $0.95 \text{ mV}/10\text{ke}^-$

Double band-pass
filter. Center
frequency $\sim 15\text{kHz}$

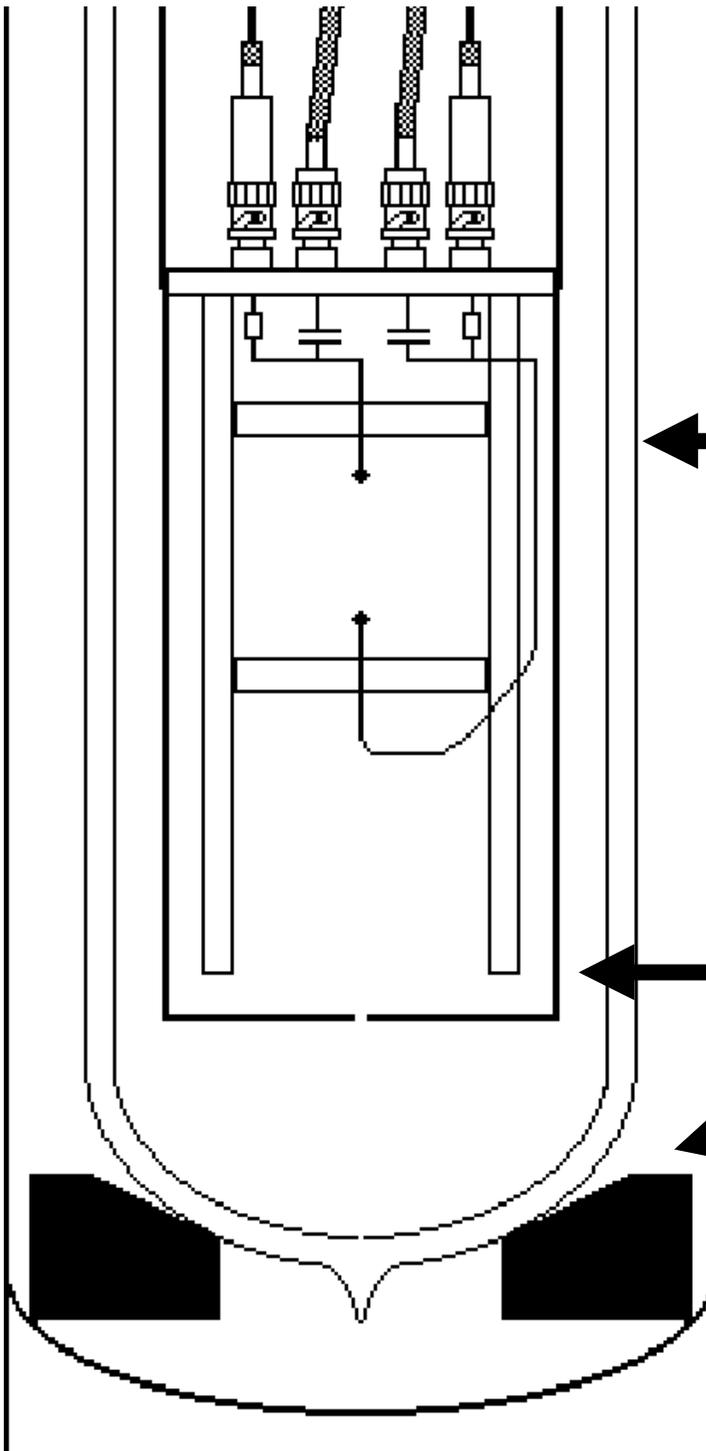
Low noise Op-Amp

20Ms/s 12bit DAC



Institute for Particle Physics ETH Zurich, Group Rubbia

Cryogenic Vessel

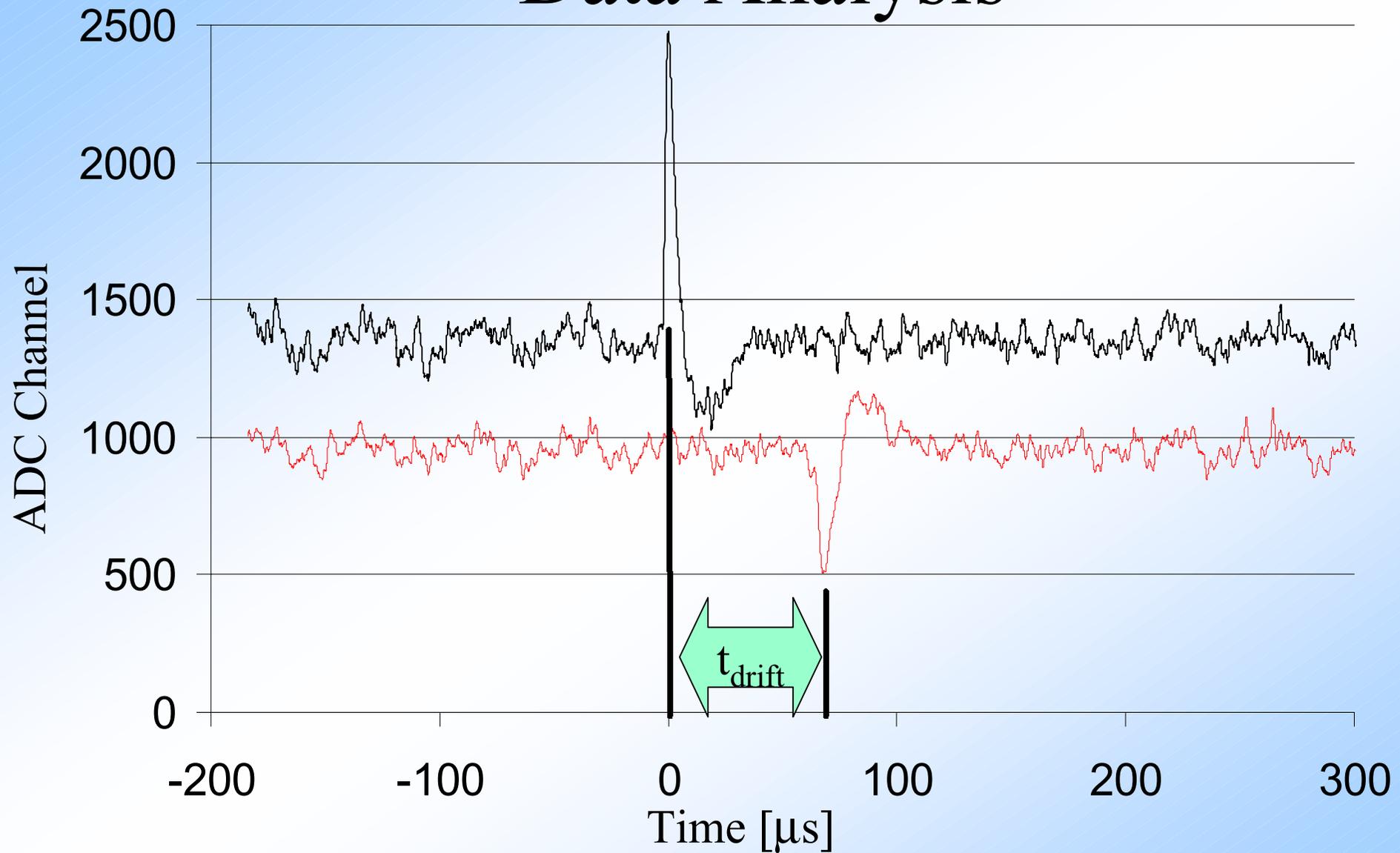


Glass Dewar

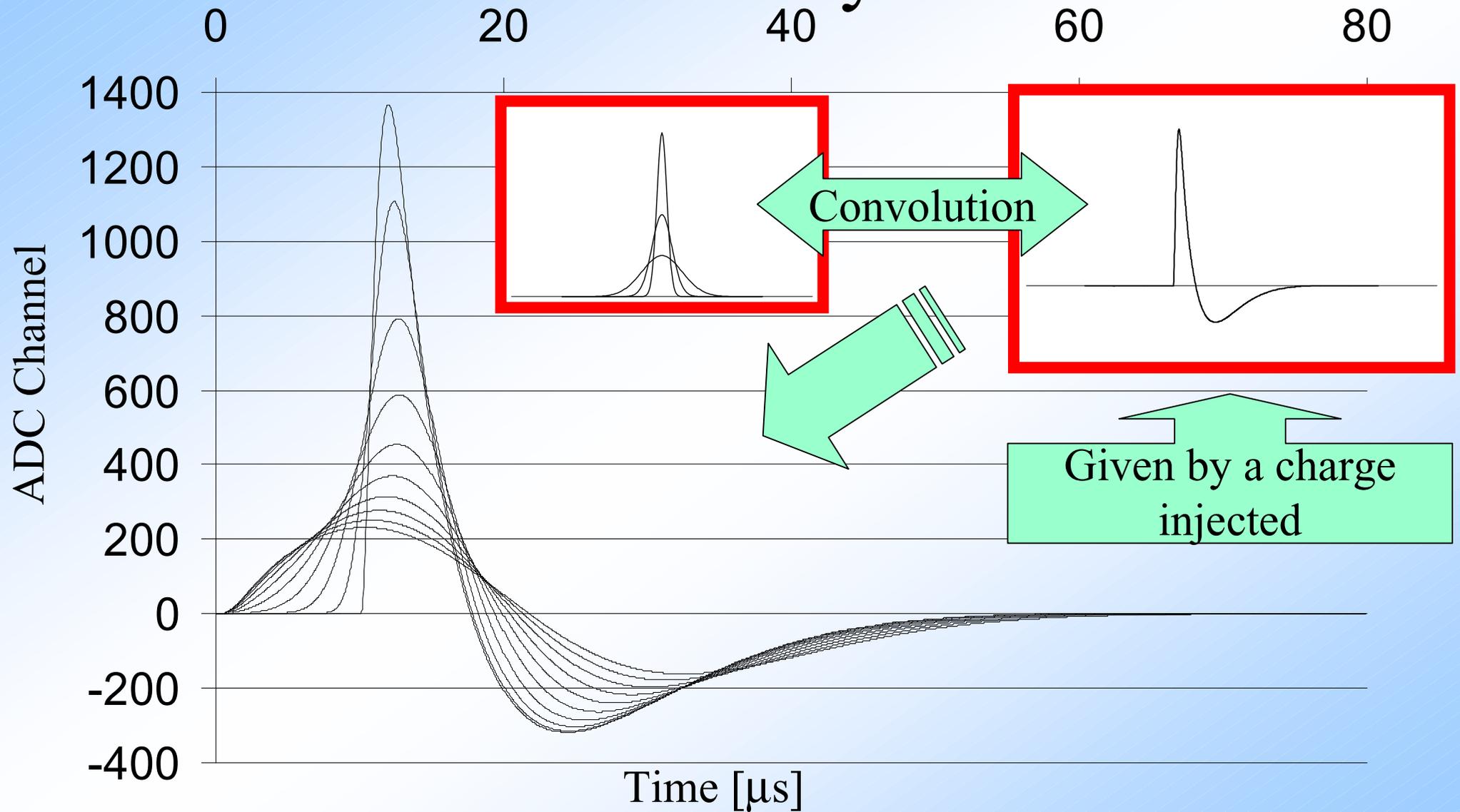
Purity Monitor

Polyethylene

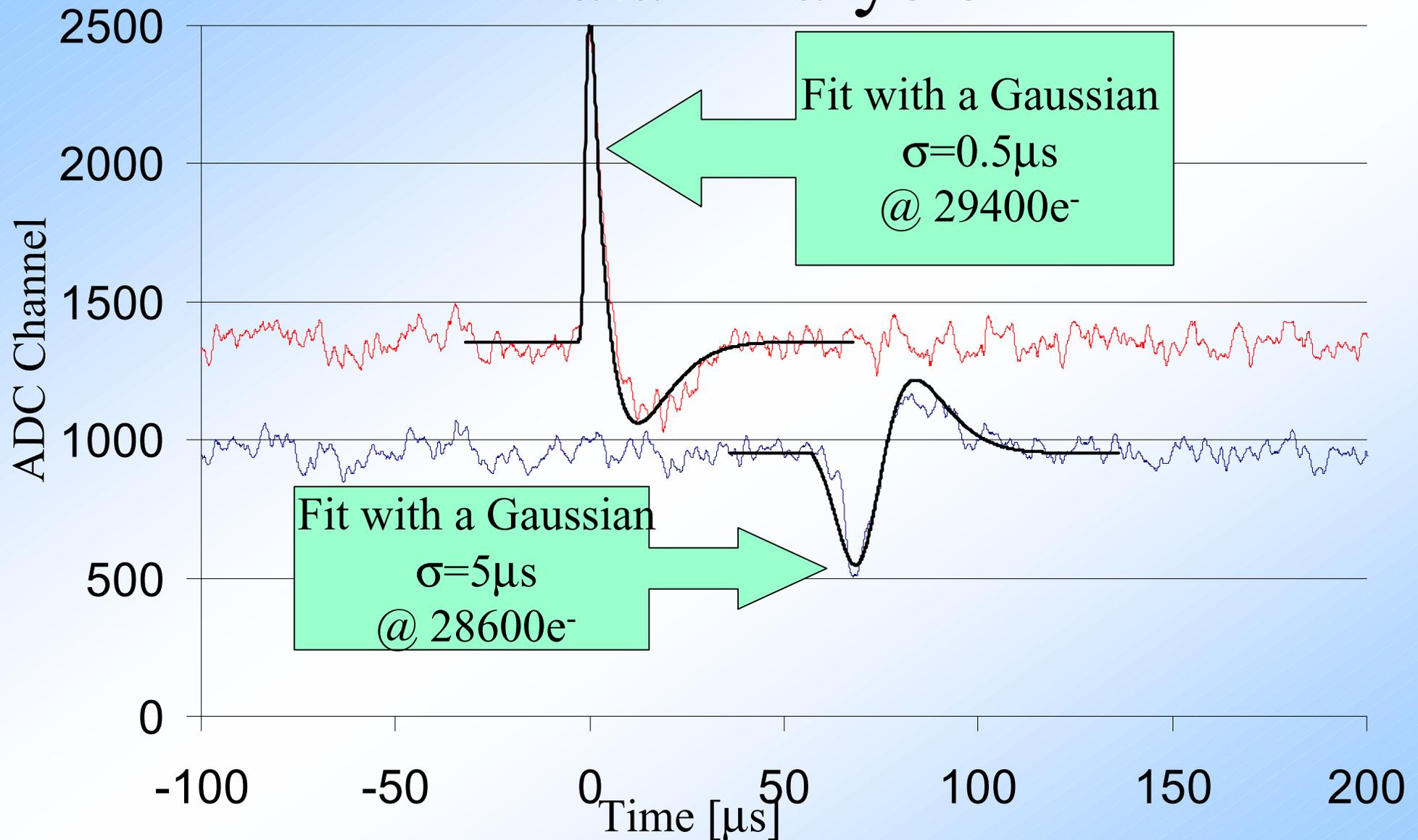
Data Analysis



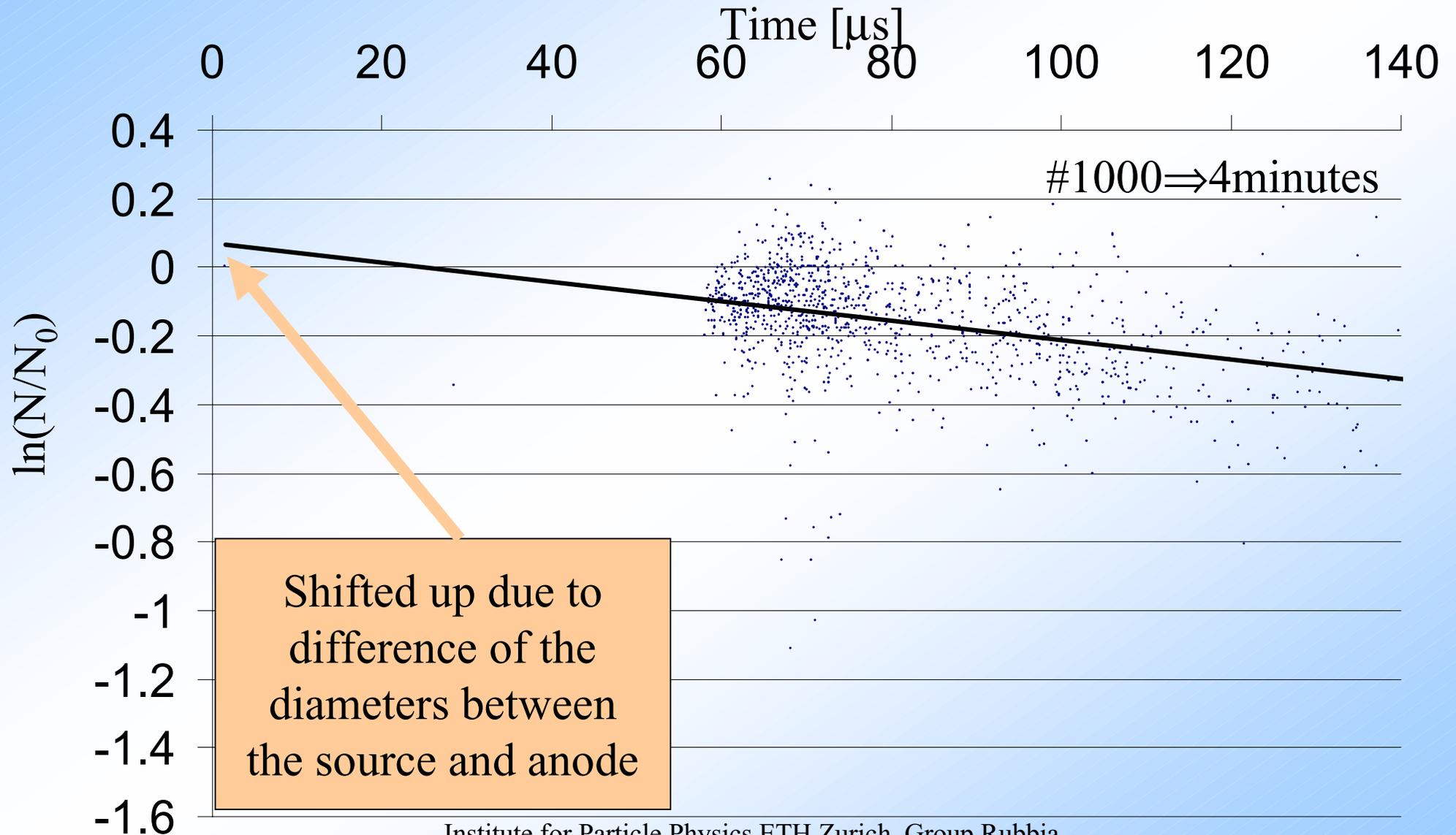
Data Analysis



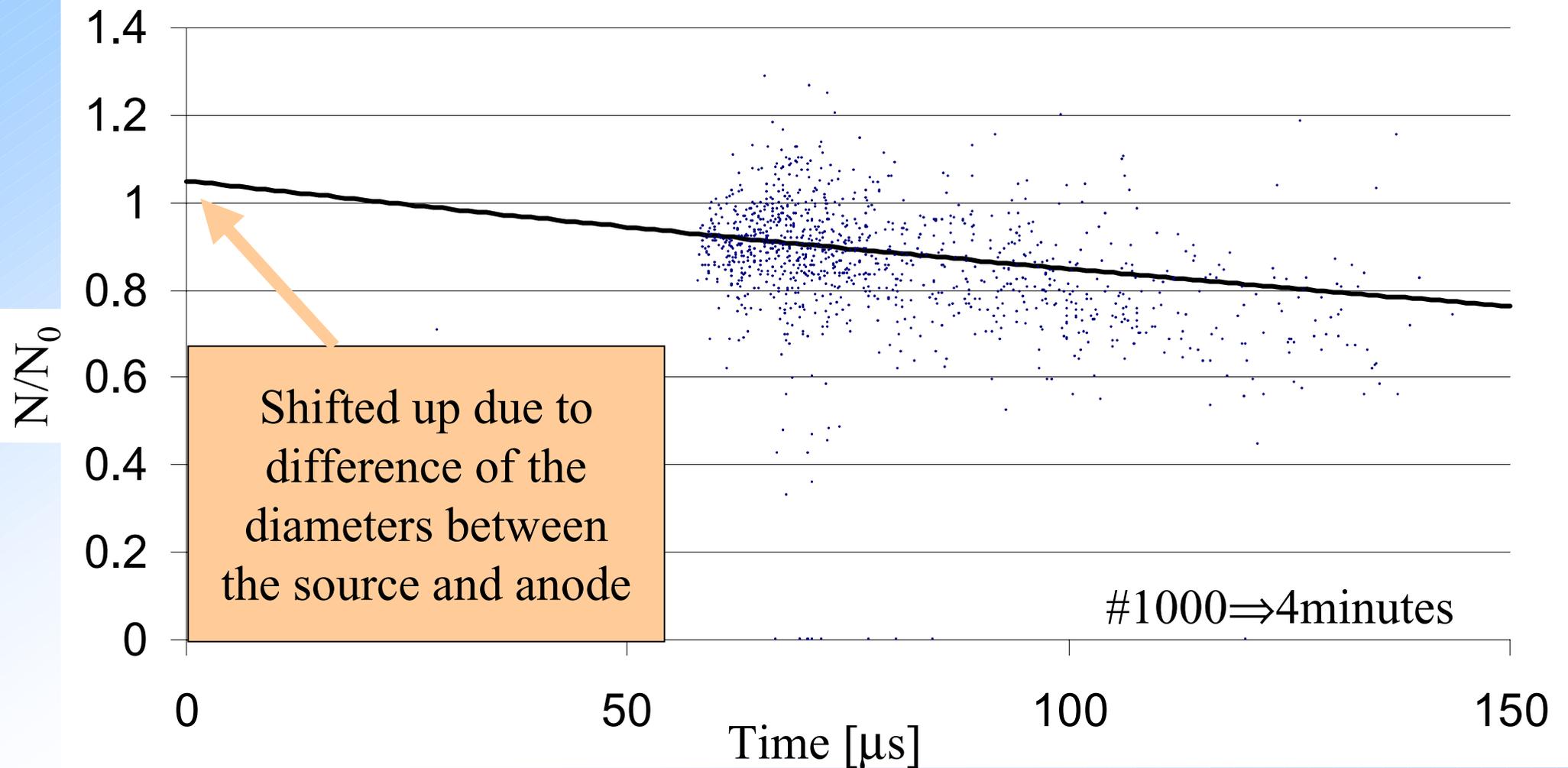
Data Analysis



Data Analysis

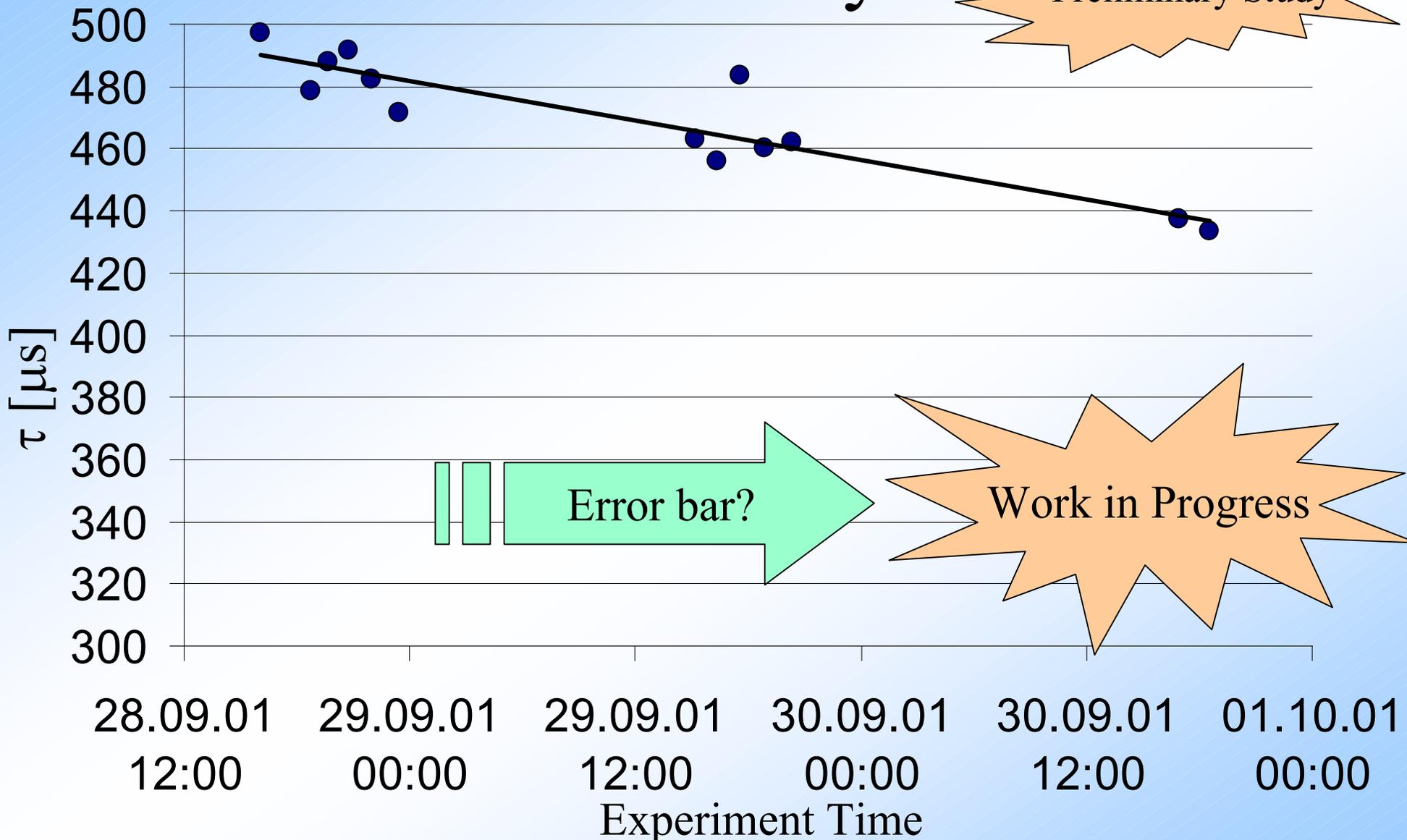


Data Analysis



Data Analysis

Preliminary Study



Need for a Purity Monitor

- Long Drift Length Needs Long Drift Time
 - The LAr should be so pure that almost all electrons produced by the particle arrive at the grid.
 - Filling and on-line operations will need the information of the LAr purity.
- Signal Correction in the Analysis
 - The LAr is a fine grain homogeneous calorimeter, it means that the number of free e^- is proportional to the deposited energy.
 - Off-line analysis will need precise information of the LAr purity

Purity Monitor with Spherical Source

After 1½ Years of Work!

