
Physics potential at the Neutrino Factory: Can we benefit from more than just detecting muons?

A.Bueno, M.Campanelli, A.Rubbia
ETH Zurich

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Neutrino Oscillations at the NuFactory

In a Neutrino Factory we can study in principle 12 independent processes:

Oscillation probabilities:

$\mu^- \rightarrow e^- \bar{\nu}_e$	ν_μ	
	$\nu_e \rightarrow e^-$	appearance
	$\nu_\mu \rightarrow \mu^-$	disappearance
	$\nu_\tau \rightarrow \tau^-$	appearance
$\bar{\nu}_e \rightarrow e^+$		appearance
$\bar{\nu}_\mu \rightarrow \mu^+$		disappearance
$\bar{\nu}_\tau \rightarrow \tau^+$		appearance

Plus their charge conjugates with μ^+ beam

$$P(\nu_e \rightarrow \nu_e) = 1 - \sin^2 2\theta_{13} \Delta_{32}^2$$

$$P(\nu_e \rightarrow \nu_\mu) = \sin^2 2\theta_{13} \sin^2 \theta_{23} \Delta_{32}^2$$

$$P(\nu_e \rightarrow \nu_\tau) = \sin^2 2\theta_{13} \cos^2 \theta_{23} \Delta_{32}^2$$

$$P(\nu_\mu \rightarrow \nu_\mu) = 1 - 4 \cos^2 \theta_{13} \sin^2 \theta_{23} (1 - \cos^2 \theta_{13} \sin^2 \theta_{23}) \Delta_{32}^2$$

$$P(\nu_\mu \rightarrow \nu_\tau) = \cos^4 \theta_{13} \sin^2 2\theta_{23} \Delta_{32}^2$$

$$P(\nu_\tau \rightarrow \nu_\tau) = 1 - 4 \cos^2 \theta_{13} \cos^2 \theta_{23} (1 - \cos^2 \theta_{13} \cos^2 \theta_{23}) \Delta_{23}^2$$

$$\Delta_{32}^2 = \sin^2(1.27 \Delta m_{32}^2 L / E)$$

How to get most of the potential?

Experimentally, it is not possible to disentangle exactly all 12 oscillation processes.

However, we believe that once such a powerful machine is built, one should try to extract most of the information, with a detector as versatile as possible

LAr TPC at the Neutrino Factory

ICANOE is one of the two large detectors proposed for the CERN-Gran Sasso beam.

- ★ Liquid Argon target for fine-grained event imaging
- ★ Calorimeter modules for tail-catching and muon charge+momentum determination

→Good technology for a Neutrino Factory

Event classes in ICANOE-like detector

Detector able to identify γ , e, μ and hadrons,
charge is measured only for muons.

Events can be classified into four classes, according
to the leading particle:

- ★ Electron of any charge
- ★ Muons of same sign as those circulating in ring
- ★ Muons of opposite sign (oscillation, or bg)
- ★ No leading leptons

Possible baselines

With the high fluxes foreseen at the Neutrino Factory we can think of very long baselines:

Ring location	Distance to GS	Mean density
CERN	732 km	2.8 g/cm ³
Canary	2900 km	3.2 g/cm ³
FNAL	7400 km	3.7 g/cm ³
KEK	8815 km	4.0 g/cm ³



Event rates for a 10 kton detector

		Rates		
		L=732 km	L=2900 km	L=7400 km
μ^- 10^{20} decays	ν_μ CC	226000	14400	2270
	ν_μ NC	67300	4120	680
	$\bar{\nu}_e$ CC	87100	5530	875
	$\bar{\nu}_e$ NC	30200	1990	300
μ^+ 10^{20} decays	$\bar{\nu}_\mu$ CC	101000	6380	1000
	$\bar{\nu}_\mu$ NC	35300	2240	350
	ν_e CC	197000	12900	1980
	ν_e NC	57900	3670	580

No oscillations

$E_\mu = 30 \text{ GeV}$

No polarization

No beam divergence

Detector simulation

ICANOE fully simulated for CNGS studies.

For this study, events fully simulated and passed through ICANOE fast simulation.

$$\frac{\sigma(E)_{e.m.}}{E} = \frac{3\%}{\sqrt{E(GeV)}} \quad \frac{\sigma(E)_{had}}{E} = \frac{20\%}{\sqrt{E(GeV)}} \quad \frac{\sigma(P_\mu)}{P_\mu} = 20\%$$

$$\frac{\sigma(\theta)}{\theta} = 130 \text{ mrad} / \sqrt{p(GeV)}$$

Proper neutrino cross section used

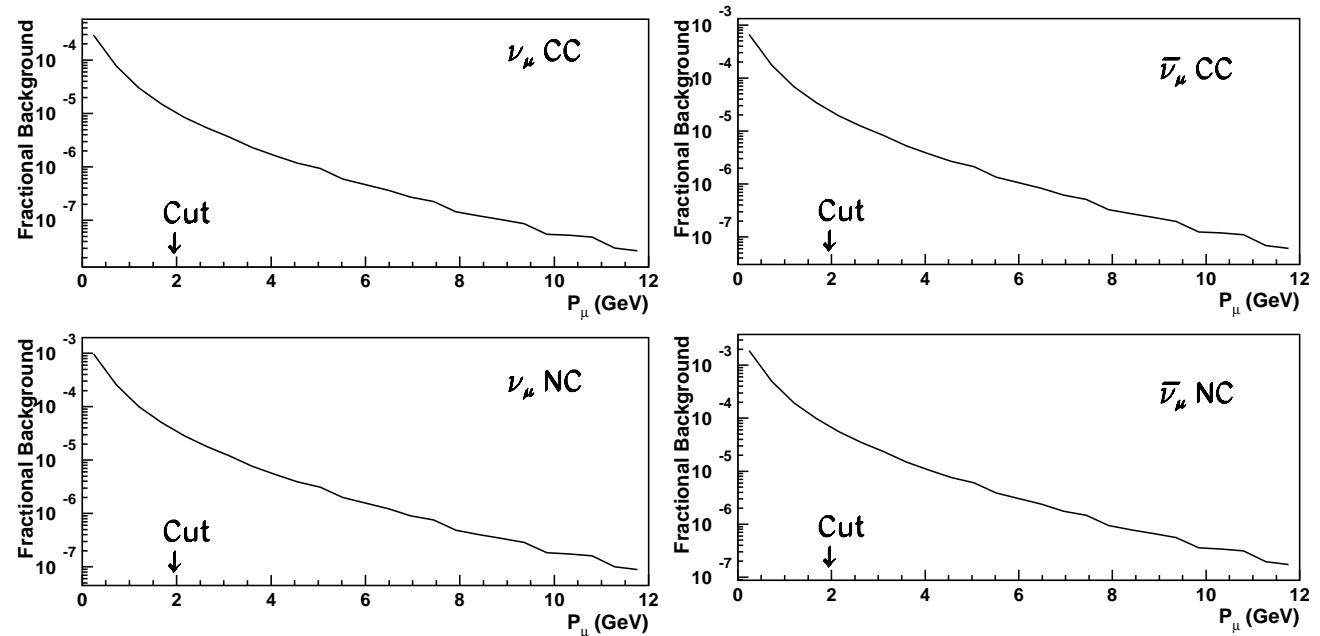
Charged π^\pm, K^\pm decay into μ^\pm for BG treatment

Background treatment

Hadron decays can produce “fake” wrong-sign muon events.

We are not aiming at large background rejection factors.

Simple momentum cut $P_\mu > 2$ GeV applied



Parameters used

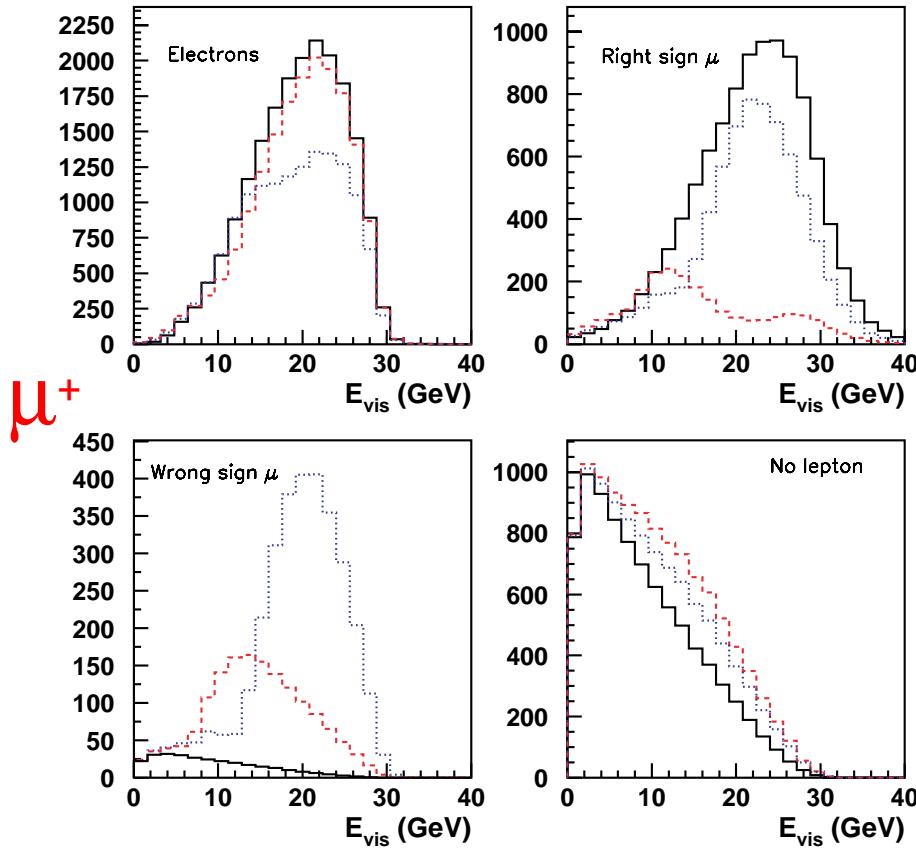
For this study, we consider as our default:

- ★ 2×10^{20} decays of 30 GeV ($\mu^+ + \mu^-$)
- ★ 3-family mixing with:
 - $\Delta m_{23}^2 = (3.5, 5, 7) \times 10^{-3} \text{ eV}^2$
 - $\sin^2 \theta_{23} = 0.5$
 - $\sin^2 2\theta_{13} = 0.05$
- ★ 10 kton ICANOE-like detector

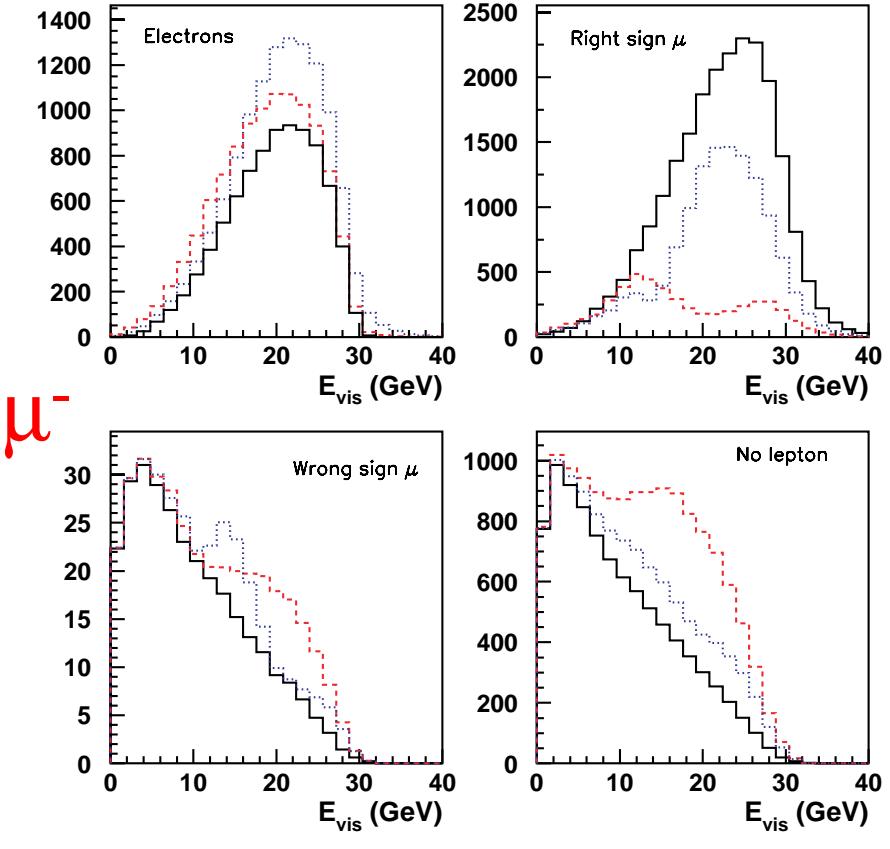
Observed Spectra

Neutrino oscillations visible in the spectra:

$L=7400 \text{ km}, 10^{21} \mu^+ \text{ decays}$



$L=7400 \text{ km}, 10^{21} \mu^- \text{ decays}$



ICANOE fast simulation

Antonio Bueno & Mario Campanelli & André Rubbia, ETH/Zurich May 2000

No oscillation

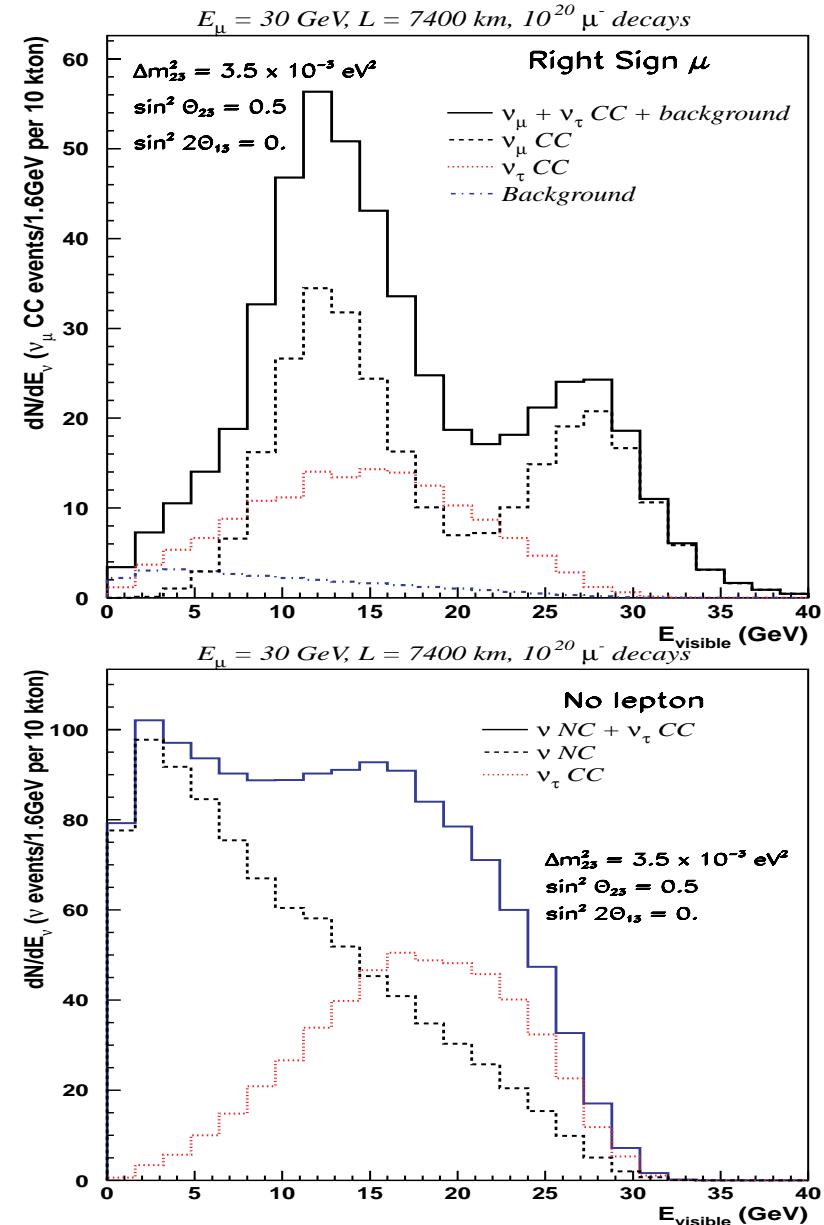
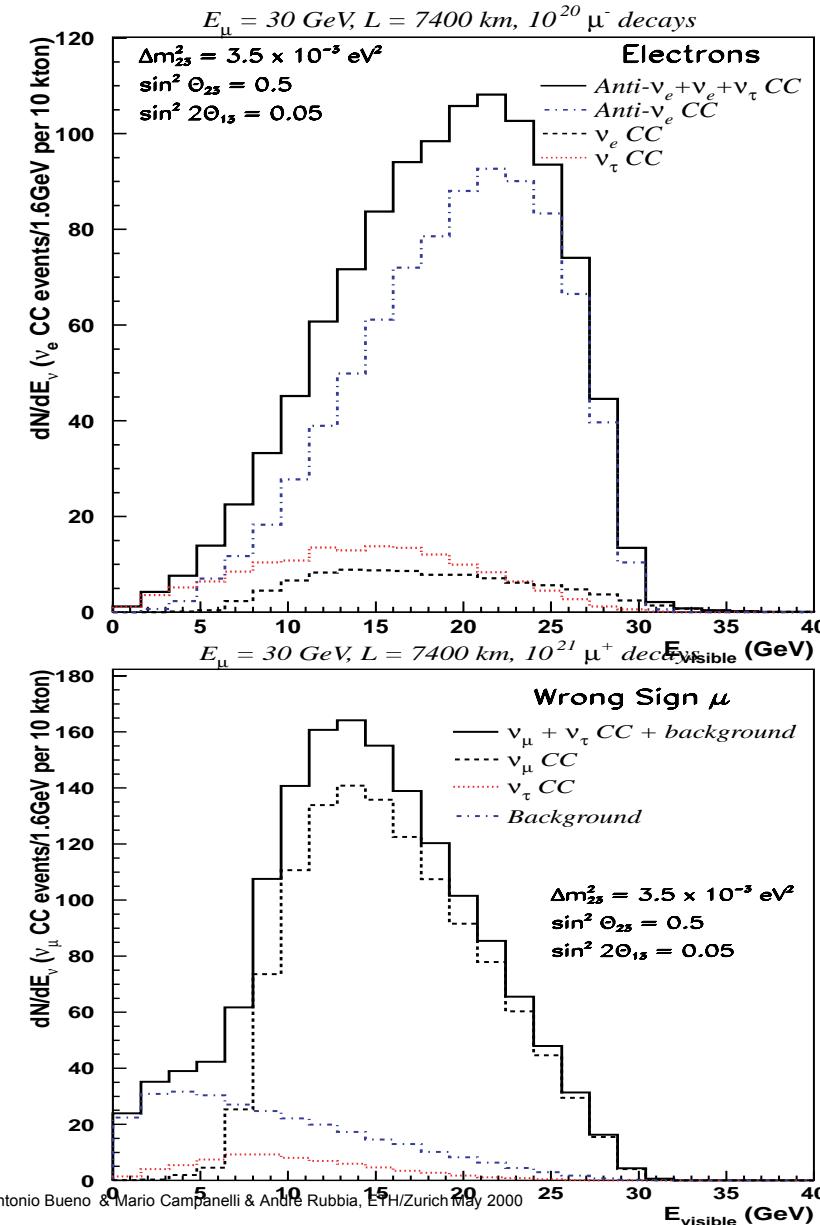
$\Delta m^2 = 3.5 \cdot 10^{-3} \text{ eV}^2$

$\Delta m^2 = 7 \cdot 10^{-3} \text{ eV}^2$

$\theta_{23} = \pi/4$

$\sin^2 \theta_{13} = 0.05$

A closer look to different event classes



ν_τ appearance

ICANOE has been designed to perform ν_τ appearance searches at the CNGS.

A similar detector at the Neutrino factory would benefit from the better signal/BG ratio given by the longest baselines:

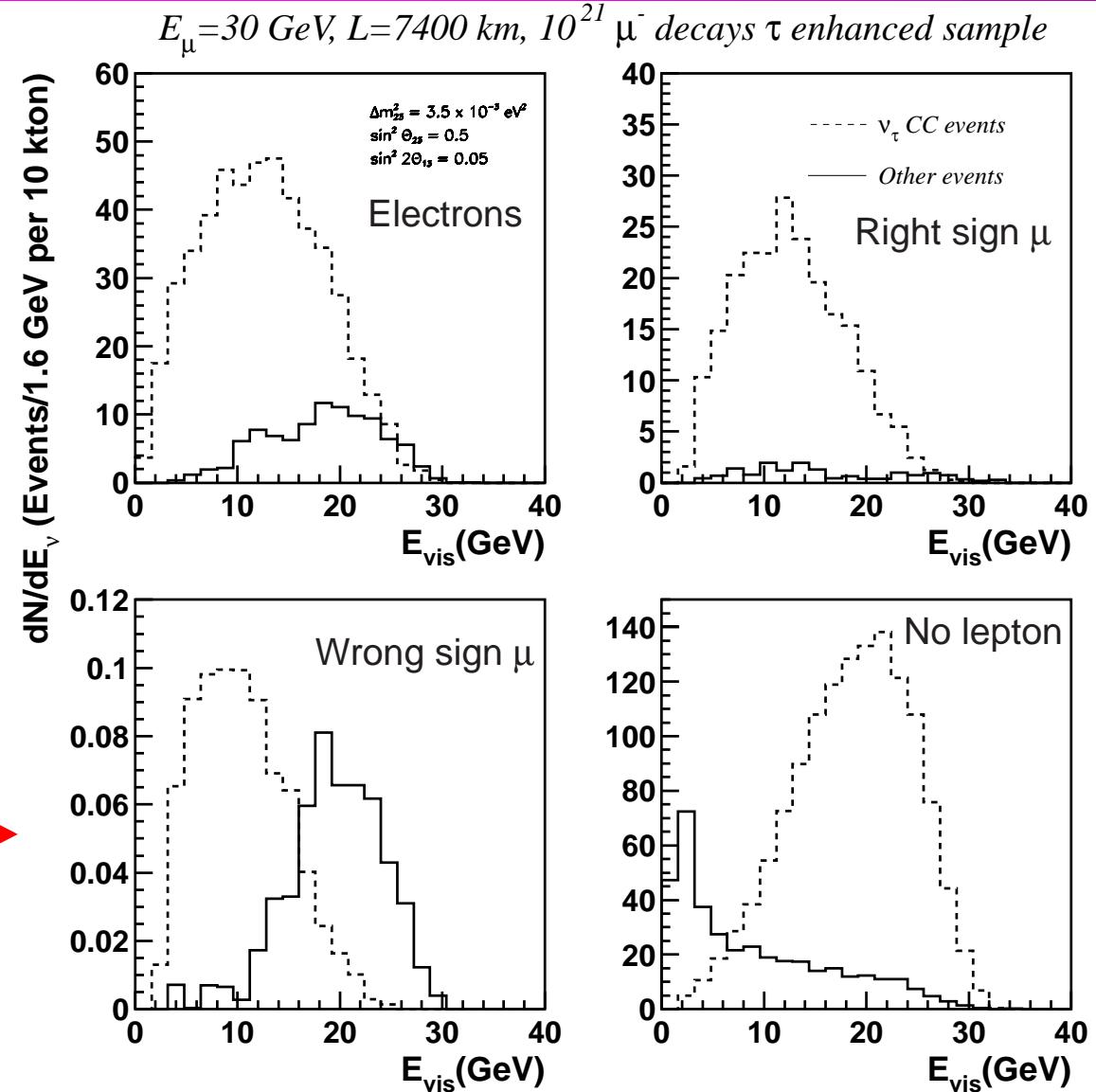
$\nu_\mu \rightarrow \nu_\tau$ appearance search					
Cuts	$\tau \rightarrow l$	CC background	Cuts	$\tau \rightarrow h$	NC background
Initial	100%	100%	Initial	100%	100%
Loose cuts					
$P_T^l < 0.5\text{GeV}$	50%	14%	$P_T^{miss} < 1\text{GeV}$	72%	40%
$P_T^{miss} > 0.6\text{GeV}$	40%	0.5%	$Q_T > 0.5\text{GeV}$	30%	2%
Tight cuts					
$P_T^l < 0.5\text{GeV}$	50%	14%	$P_T^{miss} < 1\text{GeV}$	72%	40%
$P_T^{miss} > 1\text{GeV}$	20%	0.08%	$Q_T > 1\text{GeV}$	6%	0.07%

1 event
background

τ -enhanced sample

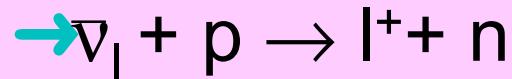
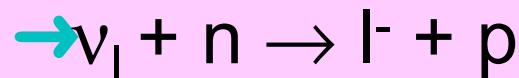
Tau contribution
enhanced for all event
classes

$\nu_e \rightarrow \nu_\tau$
oscillations!!



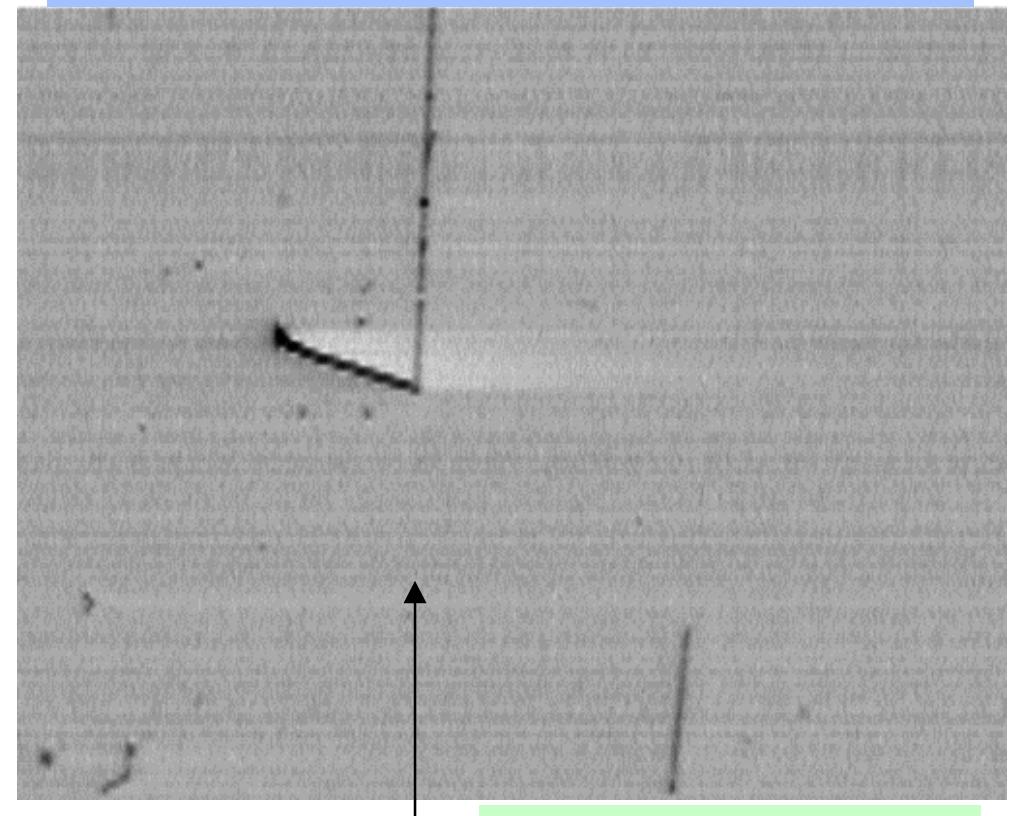
Quasi-elastic events

Provide ν - $\bar{\nu}$ separation for all flavors:



Only way to determine helicity of electron neutrino without explicit electron charge measurement

Real ν event
in the 50 liter LAr TPC in
CERN WANF beam



CERN ν -beam



Goals of Experiments at NUFACt

For second generation long baseline experiments,
the main goals will be:

- ★ Precise determination of Δm^2_{23} and Θ_{23}
- ★ Measurement of Θ_{13}
- ★ Study of matter effects
- ★ Study of CP violation

Fitting procedure

Parameters are determined by fit of visible energy from the different classes.

We use:

- ★ χ^2 for >40 events in bin
- ★ -log **L** for <40 events in bin

Beam systematics: 2% uncorrelated (25 bins)

Background added in fit

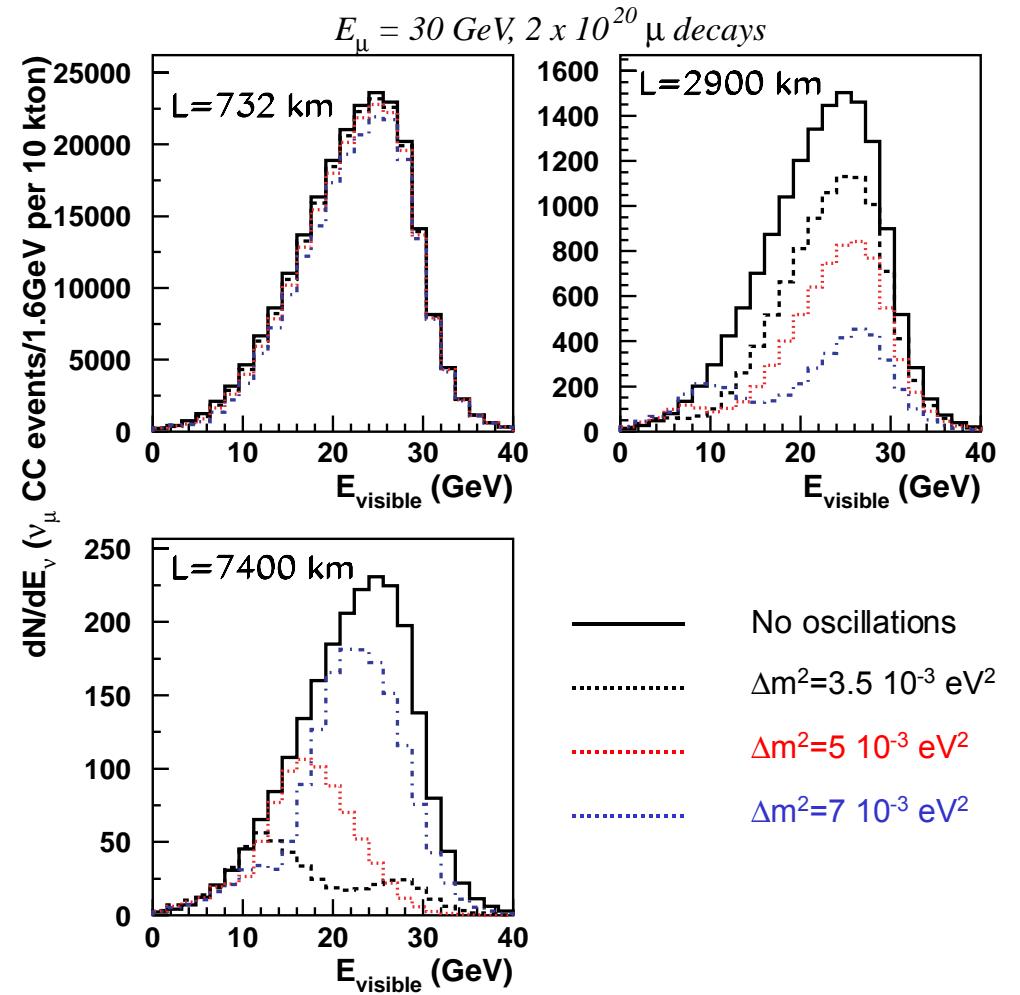
Earth density and oscillation parameters can vary in the fit or be fixed to reference value

Precise determination of $\Delta m^2_{23}, \Theta_{23}$

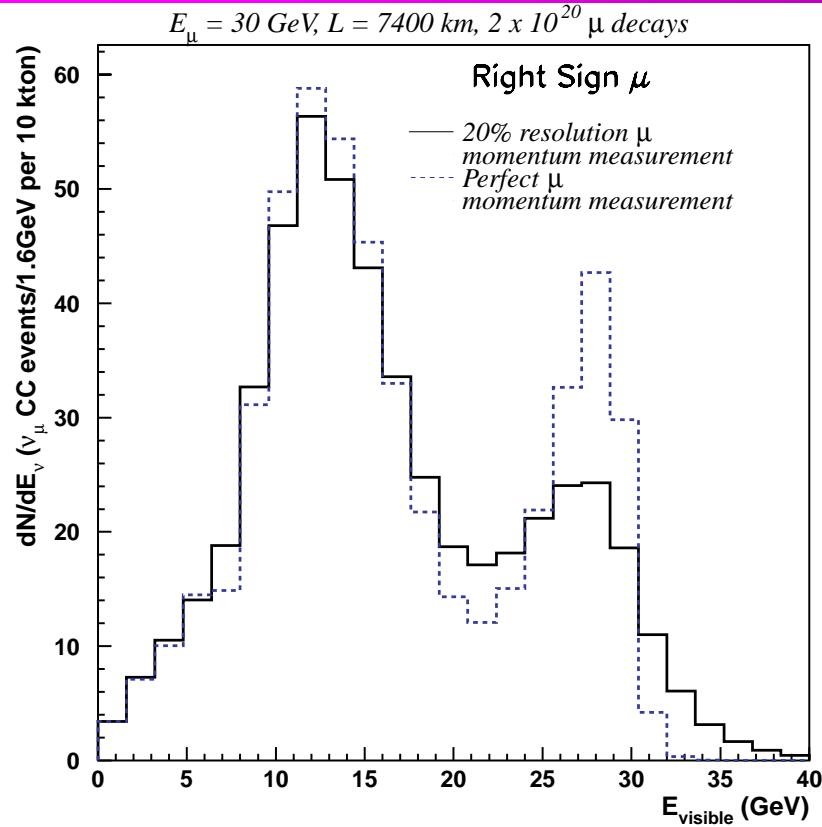
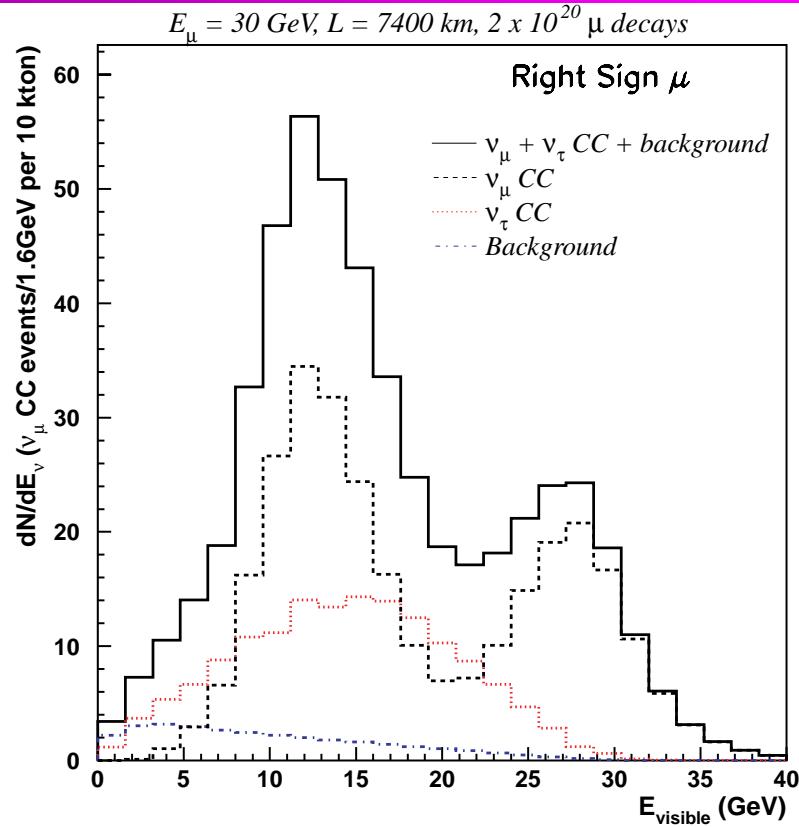
Assume $\Theta_{13} = 0 \Rightarrow$ 2-family $\nu_\mu \rightarrow \nu_\tau$ oscillations

Measurement dominated
by disappearance dip
at large distances
for right-sign μ :

- ★ Position: Δm^2_{23}
- ★ Height: Θ_{23}



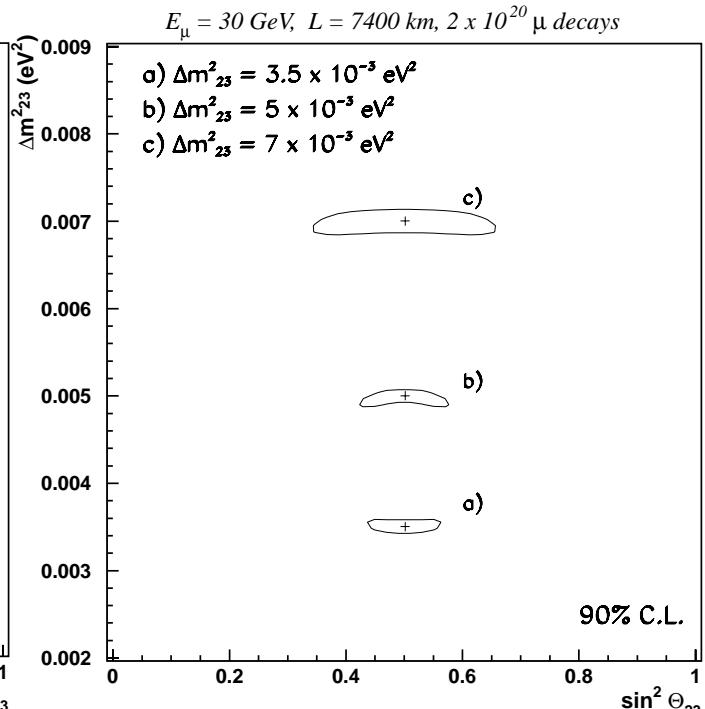
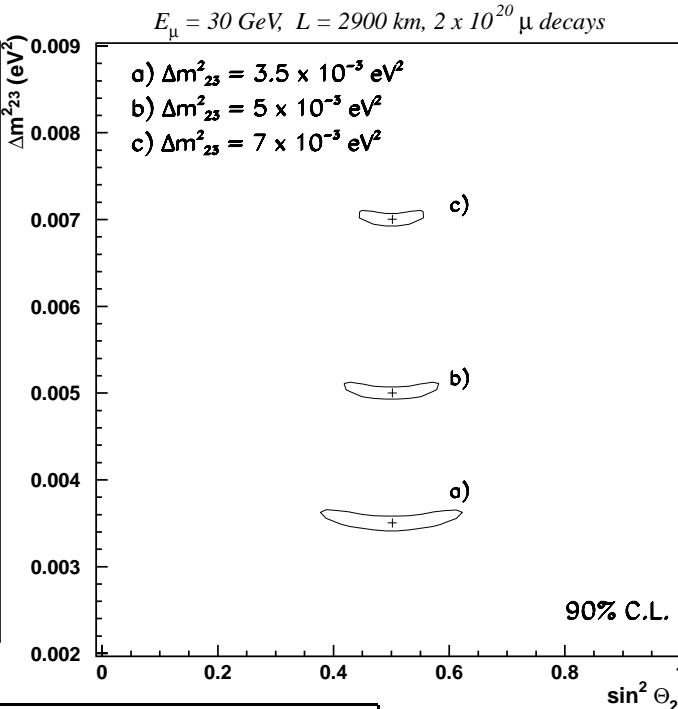
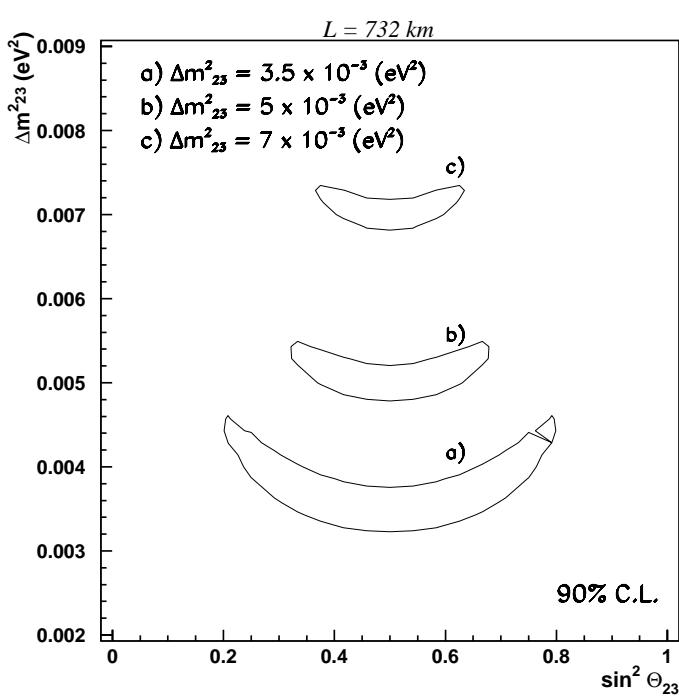
Right-sign muon disappearance



Contributions to events in the dip:

- ★ Resolution
- ★ $\nu_\mu \rightarrow \nu_\tau \rightarrow \tau \rightarrow \mu$ decays
- ★ background

Sensitivity for $\Delta m^2_{23}, \theta_{23}$ measurements



$\sin^2 \theta_{23}$ measurement		
Δm^2_{23} (eV 2)	L=2900 km	L=7400 km
7×10^{-3}	0.50 ± 0.11	0.50 ± 0.04
5×10^{-3}	0.50 ± 0.06	0.50 ± 0.06
3.5×10^{-3}	0.50 ± 0.05	0.50 ± 0.09

Error on $\Delta m^2_{23} = 1\%$

Event simulation includes:

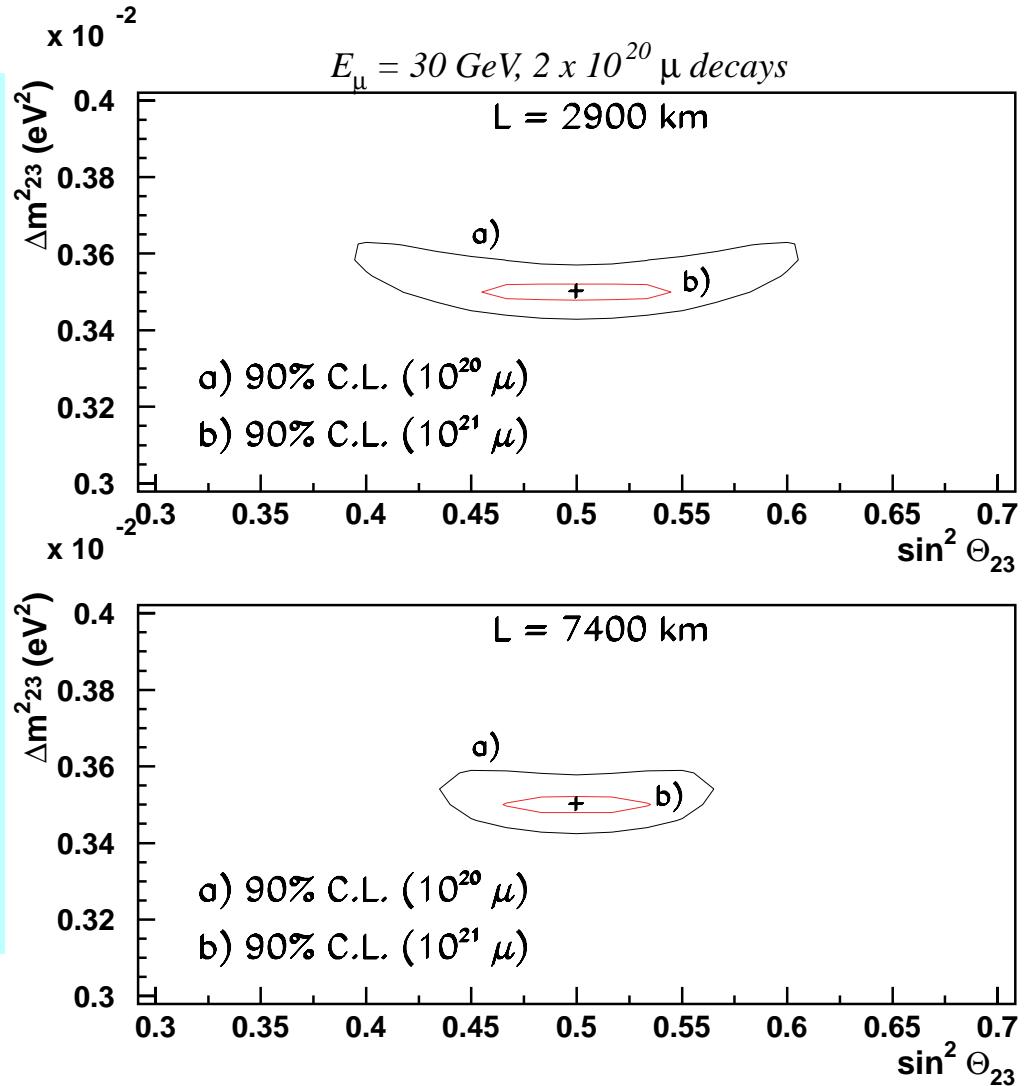
- Background
- Exclusive τ decays
- Resolution
- 2% Beam systematics

Consistent with Barger et al. hep-ph/9911524

Statistical improvements

A factor 10 more statistics can still improve the measurement at very long distances

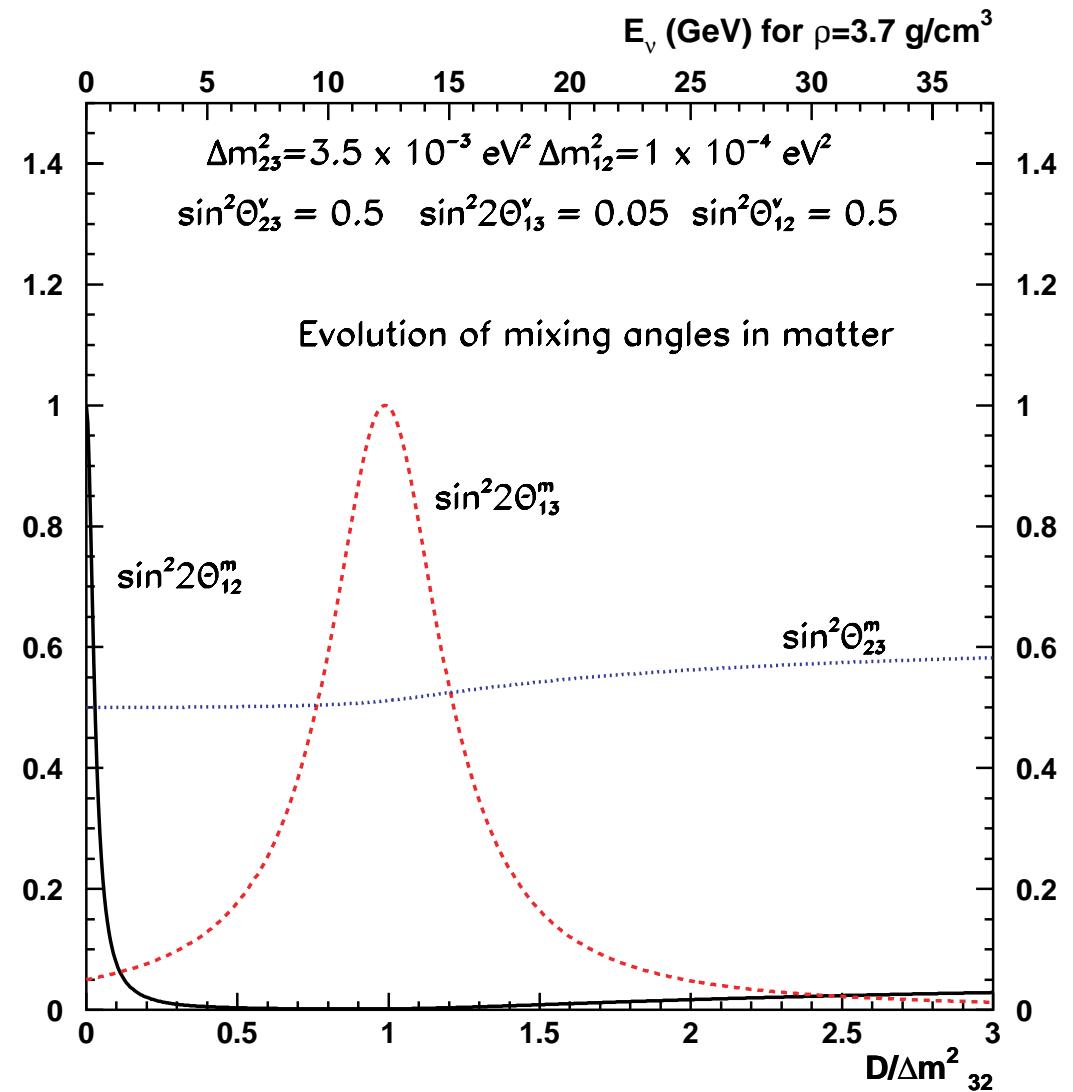
(2% systematics)



3-family mixing

With $\Theta_{13} \neq 0$, all flavors mix.

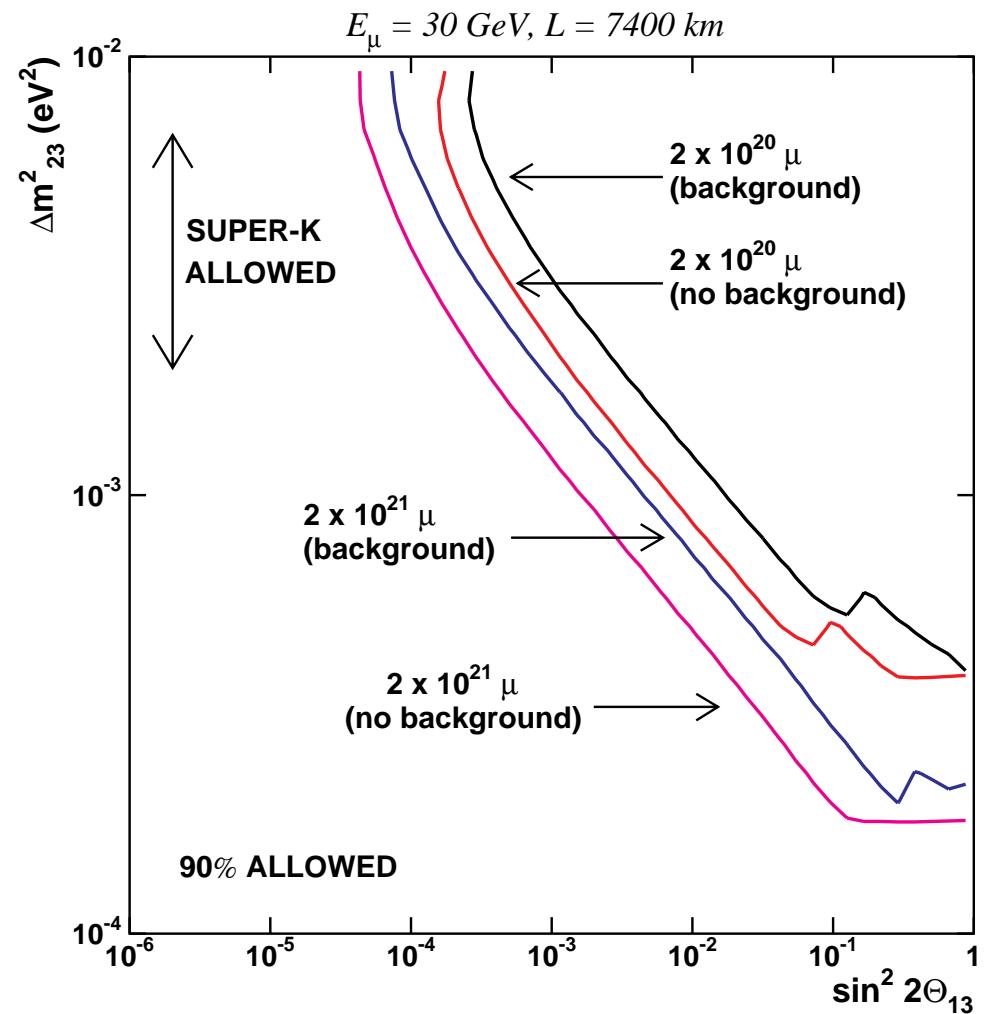
Assuming $\Delta m^2_{23} > 0$, oscillations involving ν_e ($\bar{\nu}_e$) are enhanced (suppressed) by MSW interactions with matter



Sensitivity to θ_{13}

Sensitivity to Θ_{13}
strongly depends on
background level
assumed for wrong-
sign muons

Negligible contribution
from other classes



2 orders of magnitude better than ICANOE at CNGS

Quasi-elastics

Quasi-elastic events can confirm discovery of $\nu_\mu \rightarrow \nu_e$ oscillations. For $\sin^2 2\theta_{13} = 0.05$:

ν_e appearance search with quasi-elastic Electron Class: Events for $10^{21} \mu^-$ decays						
Baseline	$\bar{\nu}_e$ CC		$\nu_\mu \rightarrow \nu_e$ CC		$\nu_\mu \rightarrow \nu_\tau$ CC, $\tau \rightarrow e$	
	Total	Elastic	Total	Elastic	Total	Elastic
$L = 732$ km	860000	43000	2090	84	3990	110
$L = 2900$ km	54300	2700	1720	70	3300	90
$L = 7400$ km	8300	410	960	40	1450	40

Measurement of Θ_{13}

Three-family mixing $\sin^2 \theta_{23} = 0.5, \sin^2 2\theta_{13} = 0.05$				
	All classes χ^2_{all}		Only muons $\chi^2_{rs\mu} + \chi^2_{ws\mu}$	
	L=2900 km	L=7400 km	L=2900 km	L=7400 km
$\Delta m_{32}^2 = 3.5 \times 10^{-3} \text{ eV}^2,$				
$\delta(\Delta m_{32}^2)$	1.4%	0.9%	1.4%	0.9%
$\delta(\sin^2 \theta_{23})$	14%	8%	16%	9%
$\delta(\sin^2 2\theta_{13})$	15%	10%	17%	15%
$\Delta m_{32}^2 = 5 \times 10^{-3} \text{ eV}^2$				
$\delta(\Delta m_{32}^2)$	0.4%	0.8%	0.4%	0.8%
$\delta(\sin^2 \theta_{23})$	11%	8%	10%	12%
$\delta(\sin^2 2\theta_{13})$	11%	9%	14%	16%
$\Delta m_{32}^2 = 7 \times 10^{-3} \text{ eV}^2$				
$\delta(\Delta m_{32}^2)$	0.4%	0.6%	0.4%	0.6%
$\delta(\sin^2 \theta_{23})$	7%	8%	8%	18%
$\delta(\sin^2 2\theta_{13})$	8%	6%	9%	20%

Θ_{13} mainly determined by wrong-sign muons.

Including electrons, NC and kinematics for τ can improve precision by more than 30%

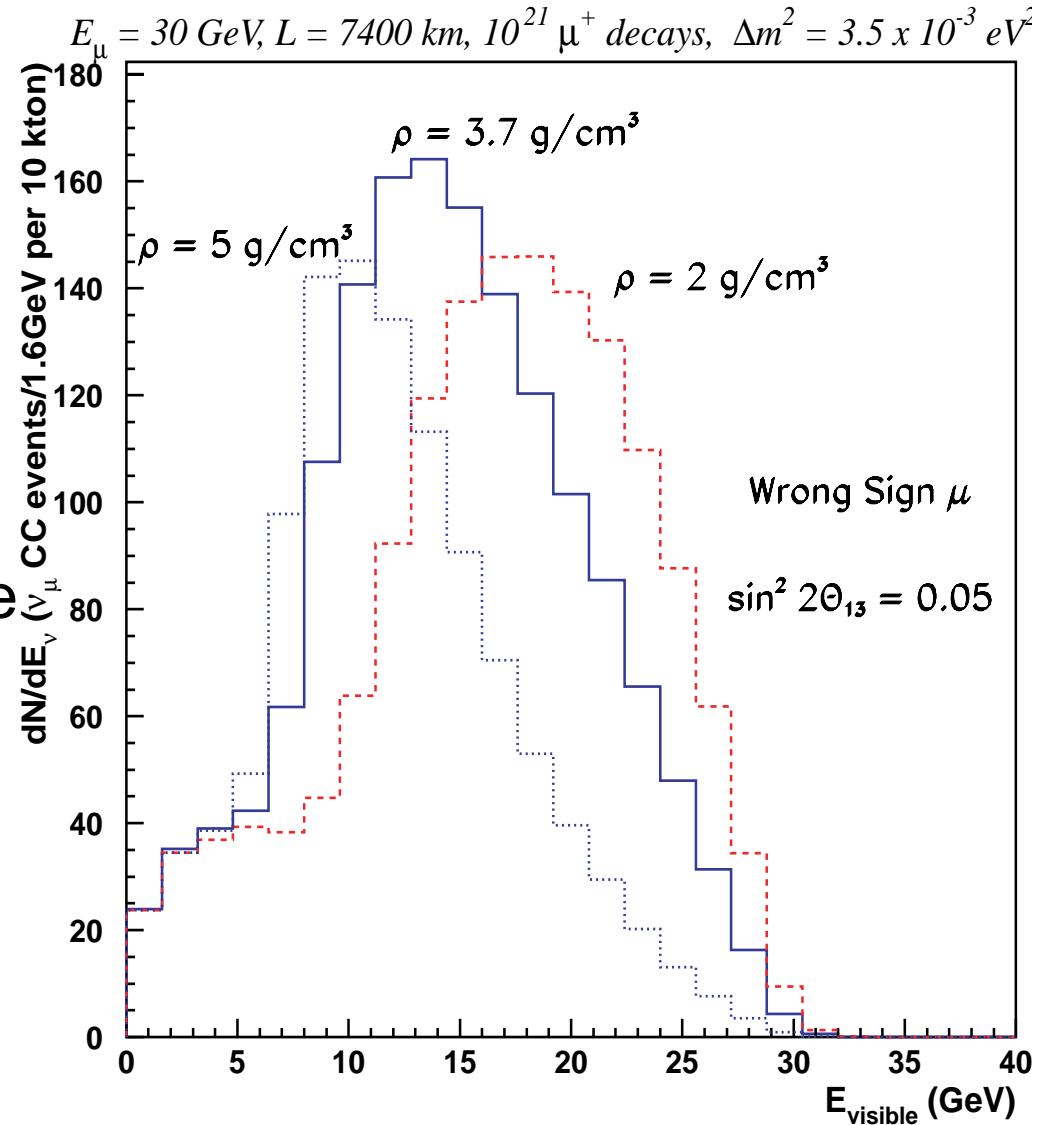
Earth density

Resonance position depends
on $\Theta_{13}\Delta m^2_{13}, \rho$

$$E_\nu^{res} \approx \frac{1.32 \times 10^4 \cos 2\theta_{13} \Delta m_{23}^2 (eV^2)}{\rho (g/cm^3)}$$

- For small θ_{13} , $\cos 2\theta_{13} \approx 1$
- Δm^2_{23} measured independently by right-sign muon disappearance

→ The resonance position, visible in wrong-sign muons, gives a measurement of the mean density



Influence of density

Density fixed to
true value:

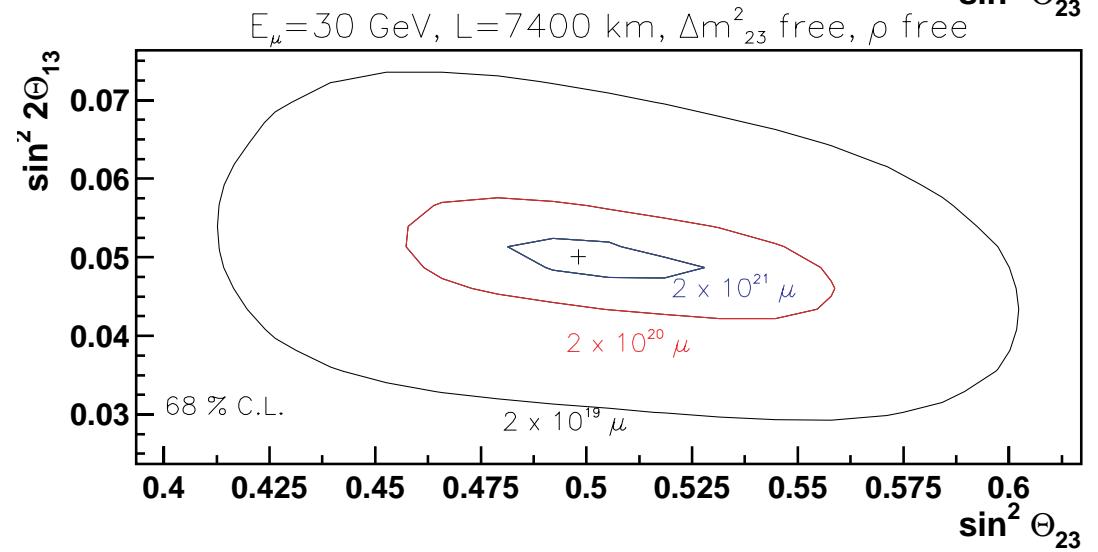
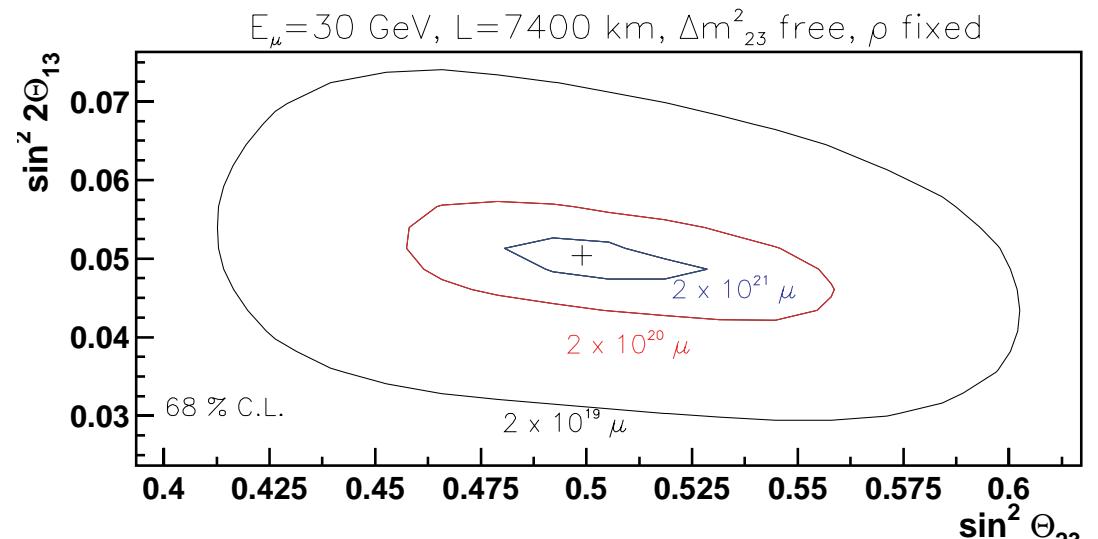
$$\sigma(\sin^2 2\theta_{13}) = 0.0071$$

$$\sigma(\sin^2 \theta_{23}) = 0.044$$

Density left free
in the fit:

$$\sigma(\sin^2 2\theta_{13}) = 0.0074$$

$$\sigma(\sin^2 \theta_{23}) = 0.050$$



Over-constraining the oscillation

For 3 active neutrinos:

$$\sum_{x=e, \mu, \tau} P(\nu_x \rightarrow \nu_y) = 1$$

Assuming new phenomena in oscillations to τ neutrinos, probabilities would change:

$$P(\nu_\mu \rightarrow \nu_\tau) \rightarrow \alpha P(\nu_\mu \rightarrow \nu_\tau)$$

$$P(\nu_e \rightarrow \nu_\tau) \rightarrow \beta P(\nu_e \rightarrow \nu_\tau)$$

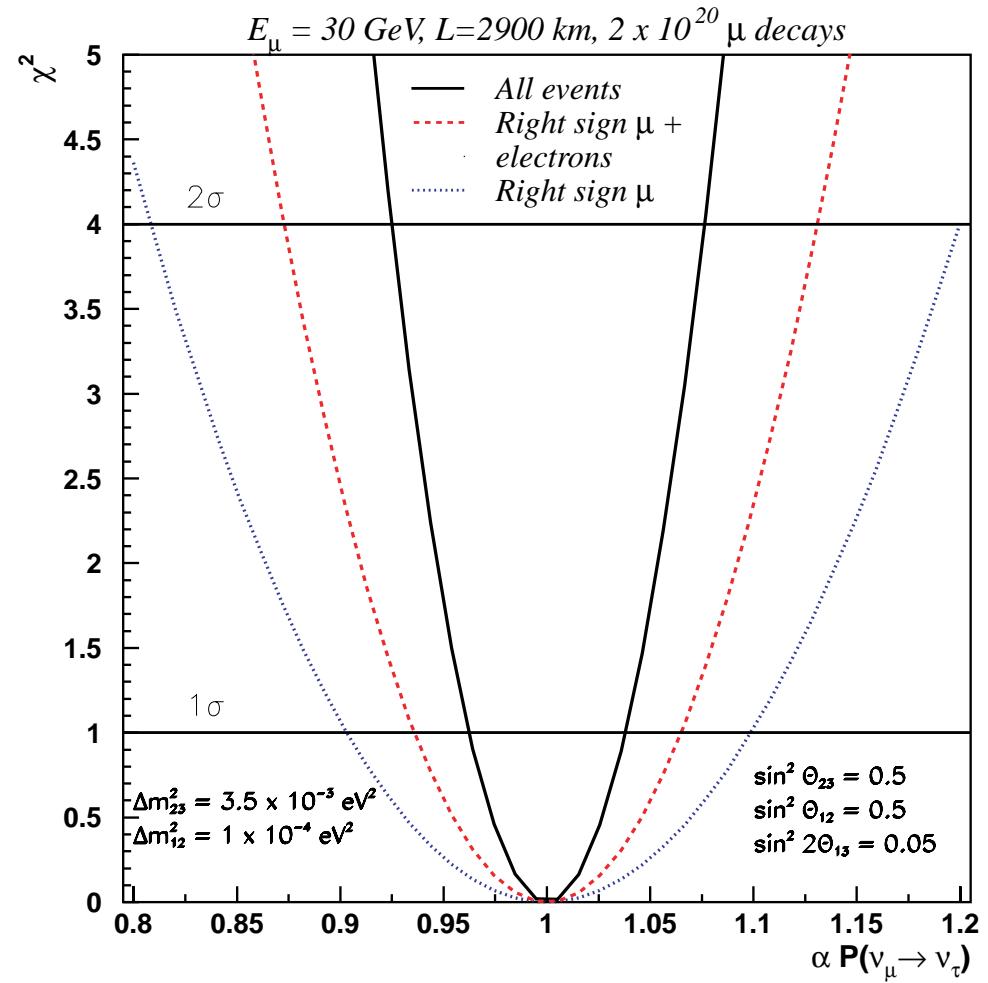
Precision on α : O(1%)

Precision on β : O(20%)

Appearance/disappearance test				
Baseline	$\Delta m_{32}^2 (\times 10^{-3} \text{ eV}^2)$	$10^{20} \mu^\pm$	$10^{21} \mu^\pm$	$10^{22} \mu^\pm$
Precision on $\alpha \Rightarrow \alpha \times P(\nu_\mu \rightarrow \nu_\tau)$				
7400 km	3.5	5.5%	2%	0.6%
	5	6%	2%	0.6%
	7	11%	3%	1%
2900 km	3.5	4%	2%	0.6%
	5	3%	1%	0.4%
	7	2.5%	1%	0.4%
Precision on $\beta \Rightarrow \beta \times P(\nu_e \rightarrow \nu_\tau)$				
7400 km	3.5	60%	20%	7%
	5	35%	10%	5%
	7	25%	7%	2%
2900 km	3.5	75%	25%	9%
	5	25%	15%	5%
	7	30%	10%	4%

Over-constraining the oscillation

As expected, all event classes contribute to over-constrain the oscillation to ν_τ

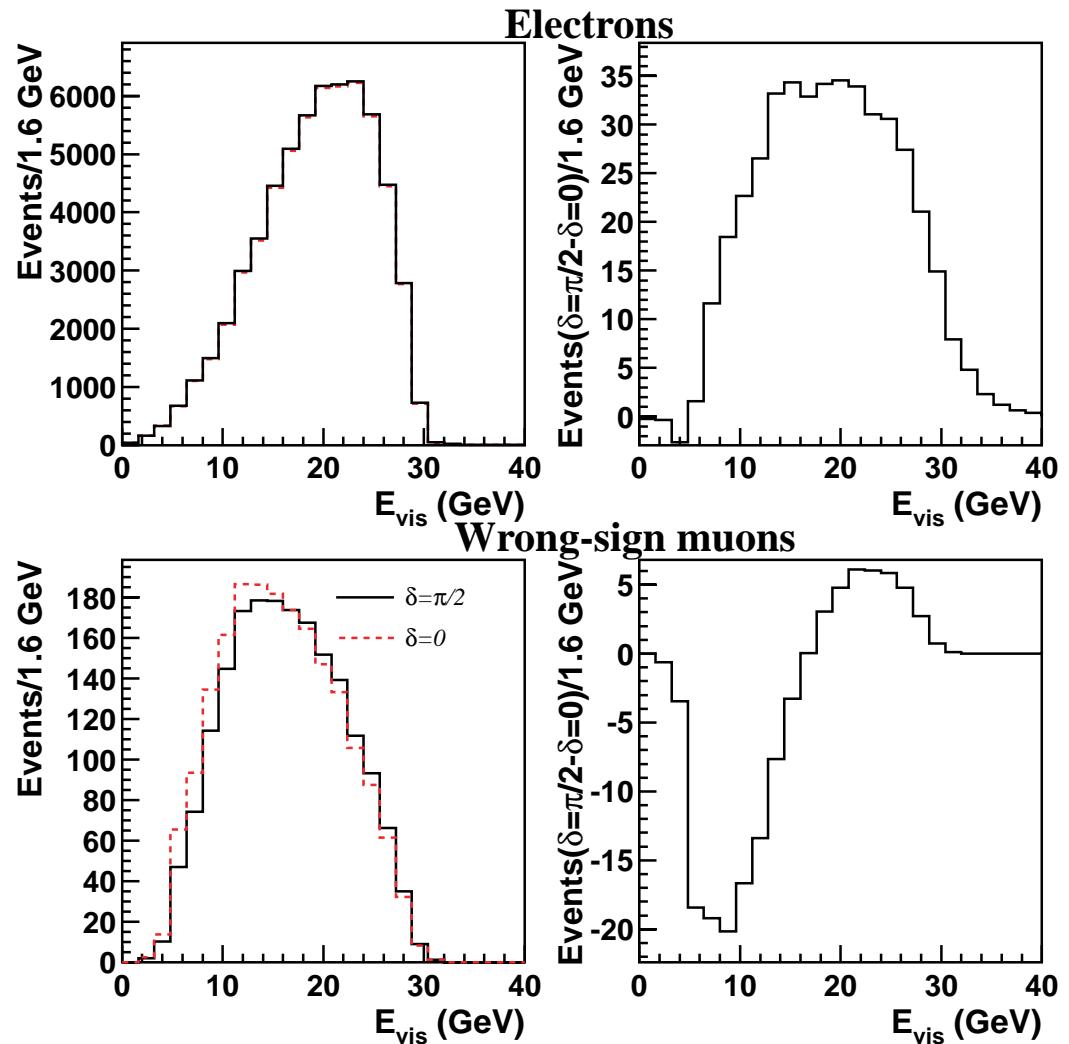


Effects of CP Violation

$L=2900 \text{ km}$, $\Delta m^2_{12}=10^{-4} \text{ eV}^2$, $\Delta m^2_{23}=3.5*10^{-3} \text{ eV}^2$, $\sin^2\theta_{23}=\sin^2\theta_{12}=0.5$, $\sin^2 2\theta_{13}=0.05$

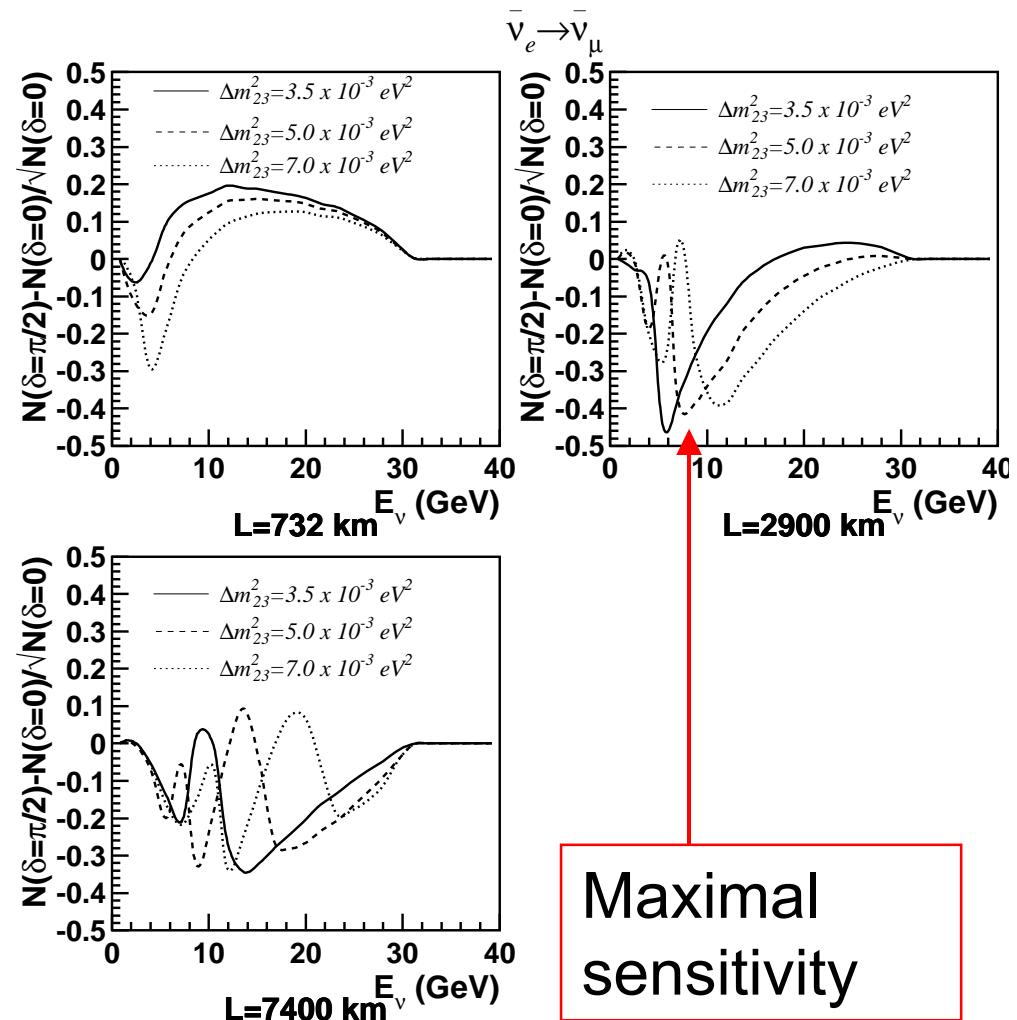
Increase number of electrons
(large, but drawn in the BG)

Change shape in wrong-sign muons
(smaller but much cleaner)



Sensitivity to CP violation

Sensitivity to CP for wrong-sign muons is maximal at L=2900 km, for neutrino energies around 10 GeV

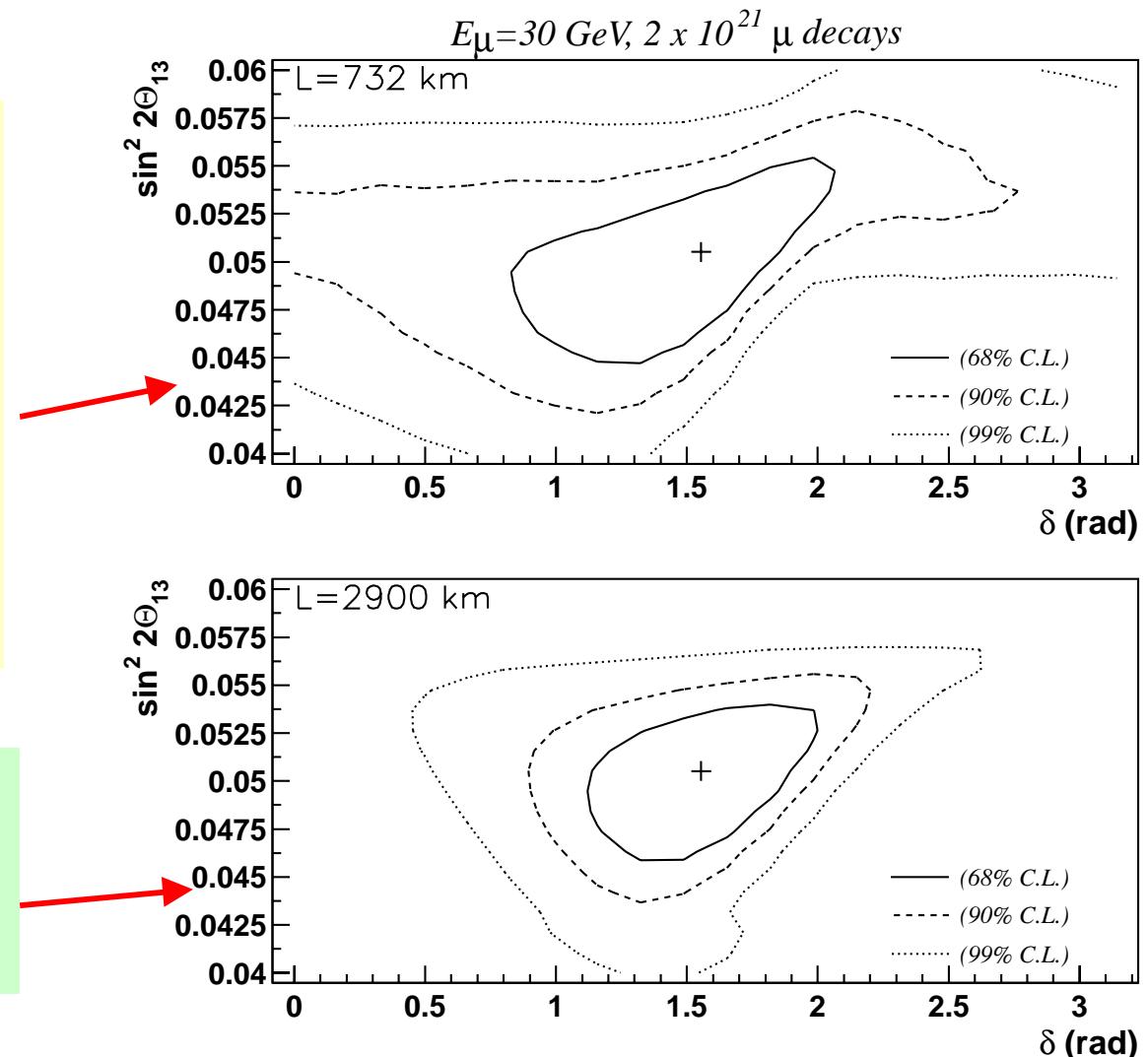


Maximal sensitivity

δ vs θ_{13}

When CP violation only leads to a normalization factor, the effect is very correlated to a variation in θ_{13}

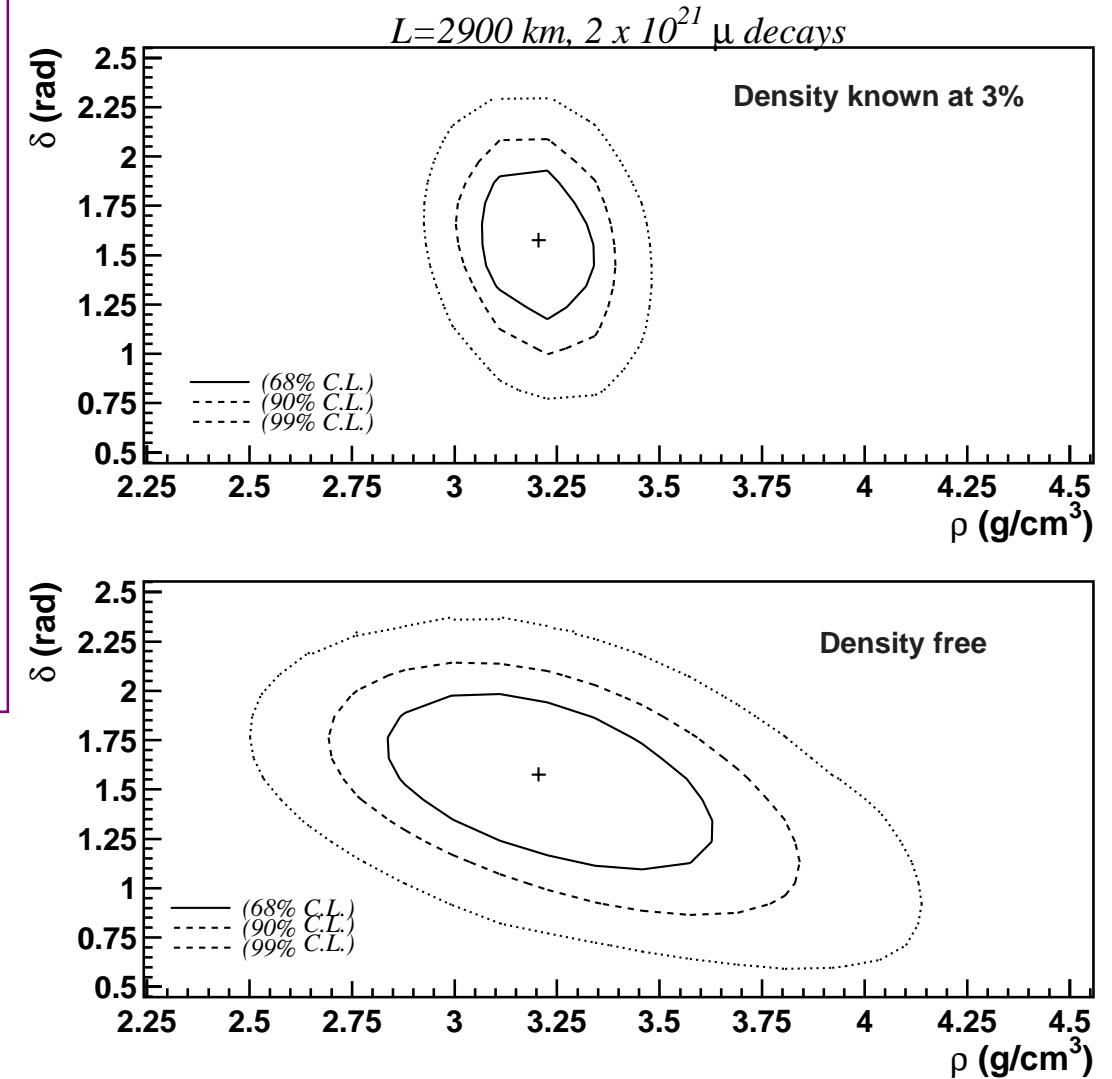
Correlation is much smaller when also a shape change occurs



δ vs ρ

In a global fit, even a complete ignorance of the matter density **does not** spoil the determination of CP

No need for more baselines!



Use of quasi-elastic events

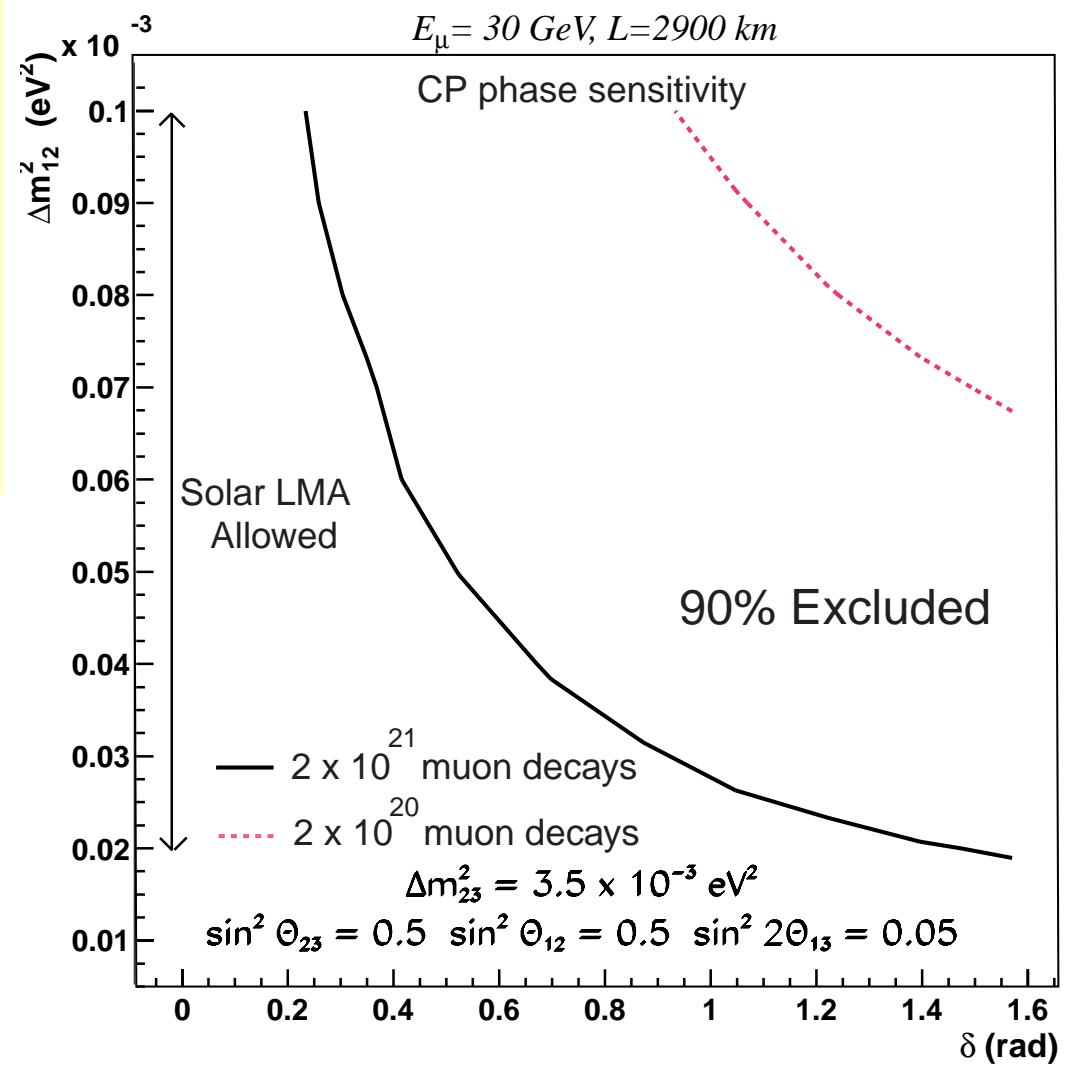
The CP violation effect is large for the electrons, but is hard to see due to ν_e background from the beam. Clean signal in quasi-elastic events

CP-violation with quasi-elastic events				
L=2900 km		N_{ele} ($\delta = 0$)	N_{ele} ($\delta = \pi/2$)	Stat. significance
$\Delta m_{32}^2 = 3.5 \times 10^{-3} \text{ eV}^2$	$10^{21} \mu^\pm$	35	26	1.5σ
$\sin^2 \theta_{23} = 0.5, \sin^2 2\theta_{23} = 0.05$	$5 \times 10^{21} \mu^\pm$	175	130	3.4σ
$\Delta m_{12}^2 = 10^{-4} \text{ eV}^2, \sin^2 \theta_{12} = 0.5$	$10^{22} \mu^\pm$	350	260	4.8σ
$\Delta m_{32}^2 = 7 \times 10^{-3} \text{ eV}^2$	$10^{21} \mu^\pm$	96	85	1.1σ
$\sin^2 \theta_{23} = 0.5, \sin^2 2\theta_{23} = 0.05$	$5 \times 10^{21} \mu^\pm$	480	425	2.5σ
$\Delta m_{12}^2 = 10^{-4} \text{ eV}^2, \sin^2 \theta_{12} = 0.5$	$10^{22} \mu^\pm$	960	850	3.6σ

Sensitivity to δ

Amplitude of CP-violating effects strongly depends on all oscillation parameters, in particular Δm_{12}^2 .

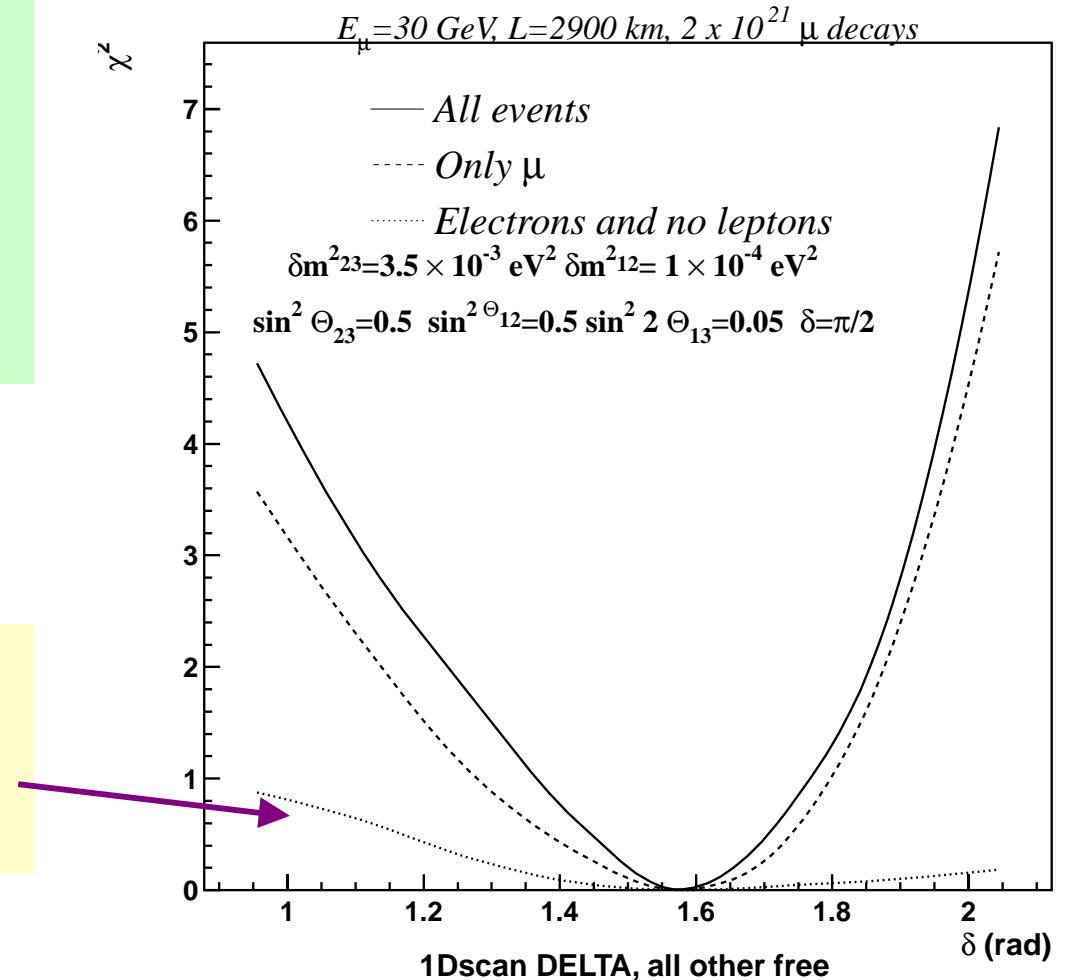
High fluxes needed, apart from a little corner of the parameter space



Measurement of CP violation

In the most favorable case, the phase δ can be measured with a 20% accuracy at $L=2900$ km.

Electron and NC-like events also contribute, albeit marginally, to the fit



Conclusions

- ★ Neutrino factory allows simultaneous study of many oscillation phenomena
 - A general-purpose detector is needed for:
 - Electrons
 - Tau identification
 - Quasi-elastic events
 - Neutral current-like events.
- ★ These events contribute to the main measurements:
 - θ_{13} , $\nu_e \rightarrow \nu_\tau$, CP violation and provide essential cross-checks
- ★ To look for new phenomena, we need to over-constrain the oscillation pattern